

# A Reduced Model for Nonlinear Interactions of Gravity Waves with Deep Convective Clouds

EGU 2011, Vienna

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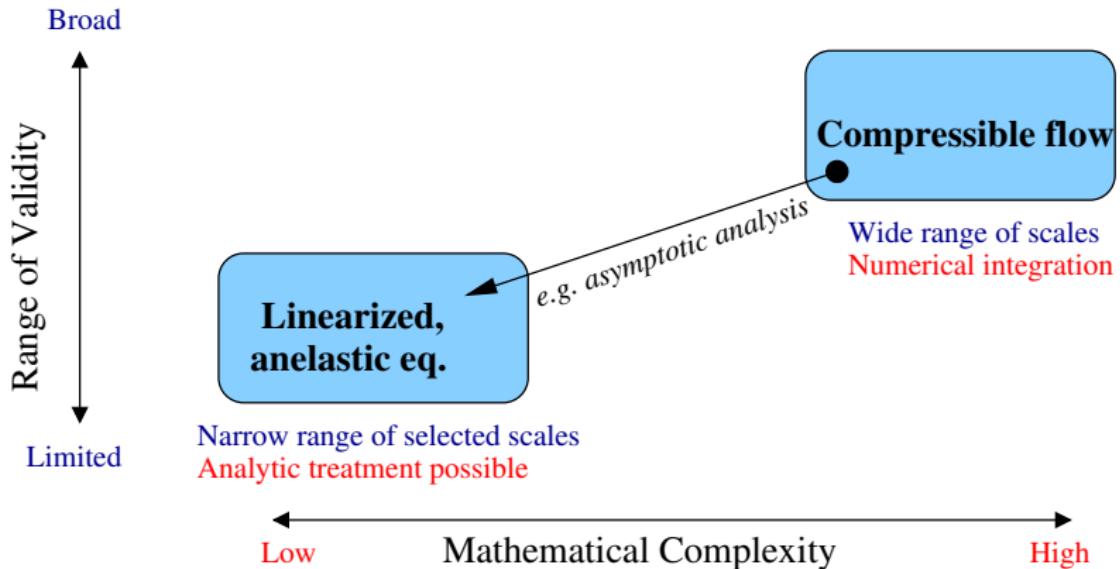
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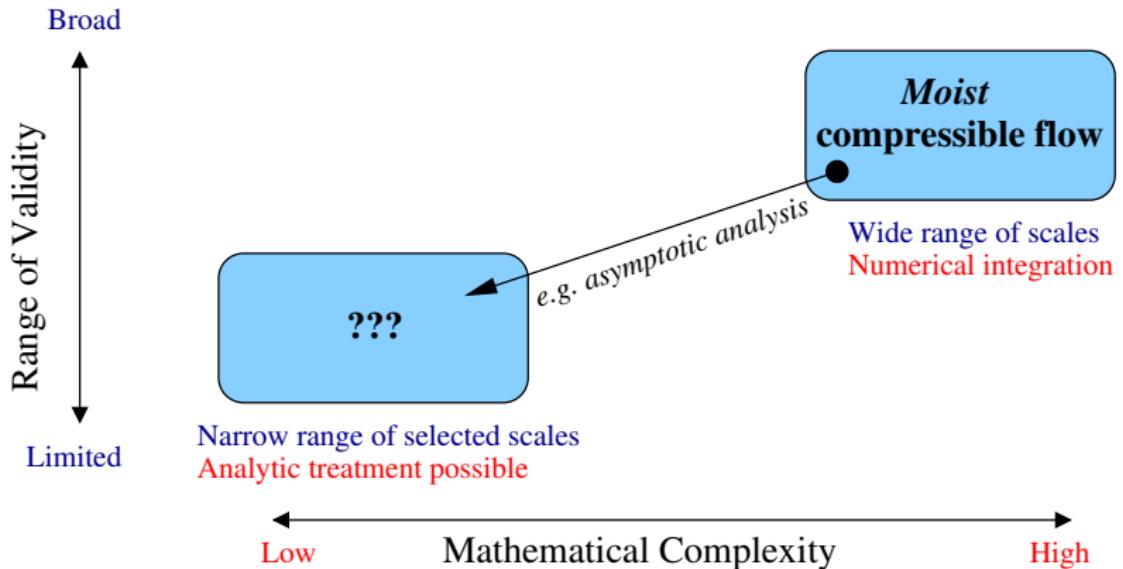
# Reduced Models



## Goal

Reduced model for gravity waves in atmosphere with deep convection.

# Reduced Models



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Reduced model for gravity waves in atmosphere with deep convection.

# Governing Equations

Conservation of mass, momentum, energy (dimensionless)

$$\begin{aligned} \rho_t + \nabla_{\parallel} \cdot (\rho \mathbf{u}) + (\rho w)_z &= 0 \\ \mathbf{u}_t + \mathbf{u} \cdot \nabla_{\parallel} \mathbf{u} + w \mathbf{u}_z + \frac{1}{Ro} (\boldsymbol{\Omega} \times \mathbf{v})_{\parallel} + \frac{1}{Ma^2} \frac{1}{\rho} \nabla_{\parallel} p &= 0 \\ w_t + \mathbf{u} \cdot \nabla_{\parallel} w + ww_z + \frac{1}{Ro} (\boldsymbol{\Omega} \times \mathbf{v})_{\perp} + \frac{1}{Ma^2} \frac{1}{\rho} p_z &= -\frac{1}{Fr^2} \\ \theta_t + \mathbf{u} \cdot \nabla_{\parallel} \theta + w \theta_z &= \underbrace{S_{\theta}}_{?} \end{aligned}$$

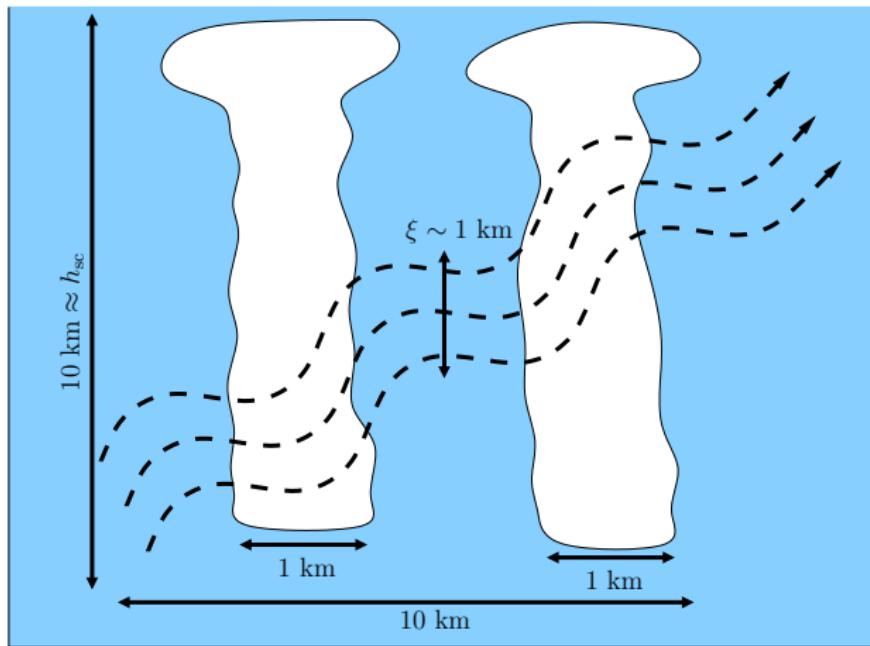
Bulk microphysical model

$$\begin{aligned} q_{v,t} + \mathbf{u} \cdot \nabla_{\parallel} q_v + w q_{v,z} &= C_{ev} - C_d \\ q_{c,t} + \mathbf{u} \cdot \nabla_{\parallel} q_c + w q_{c,z} &= C_d - C_{ac} - C_{cr} \\ q_{r,t} + \mathbf{u} \cdot \nabla_{\parallel} q_r + w q_{r,z} + \frac{1}{\rho} (\rho q_r V_T)_z &= C_{ac} + C_{cr} - C_{ev} \end{aligned}$$

# Scales

*Non-hydrostatic gravity waves modulated by deep convective towers*

## Coordinates:

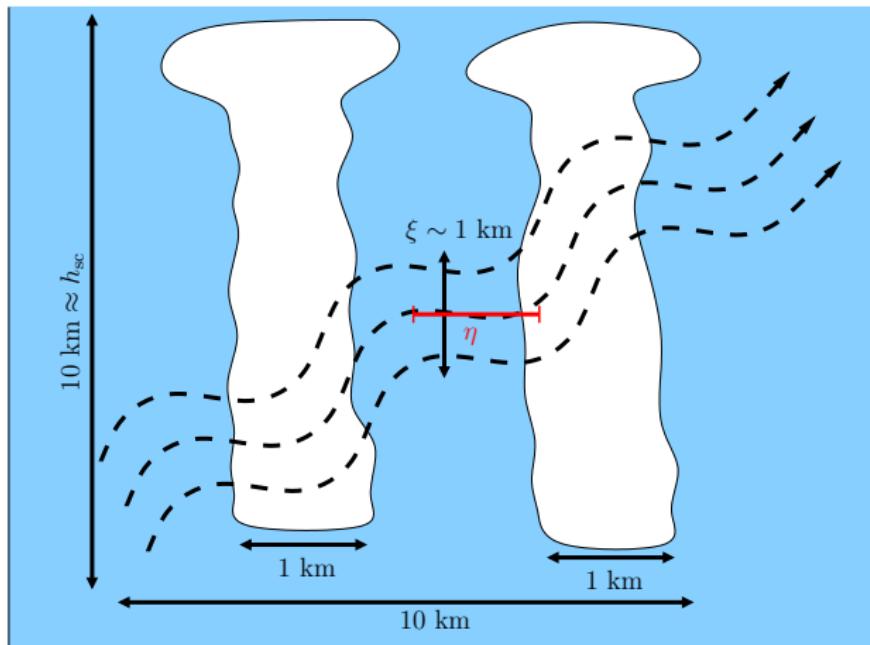


- Wave-scale:  
 $x \sim 10 \text{ km}$
  - Tower-scale:  
 $\eta \sim 1 \text{ km}$
  - Vert. scale:  
 $z \sim 10 \text{ km}$
  - Time scale:  
 $\tau \sim 100 \text{ s}$
- ⇒ Vert. Disp.:  
 $\xi \sim 1 \text{ km} = \mathcal{O}(\varepsilon)$
- ⇒ Condensate:  
 $q_{\text{released}} \sim \mathcal{O}(\varepsilon)$

# Scales

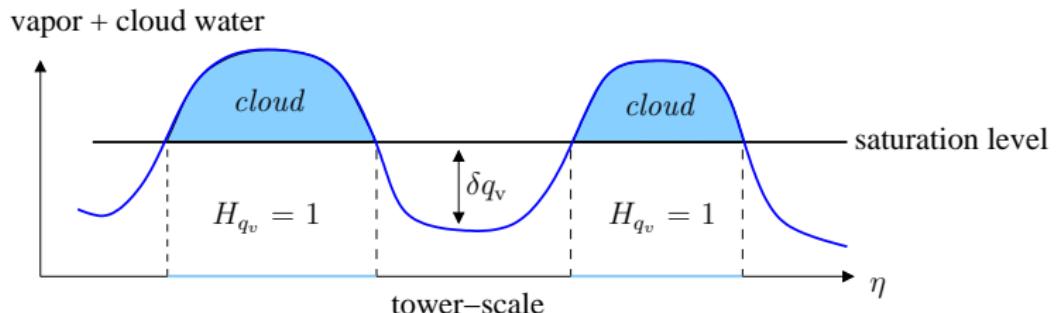
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# Leading Order Saturation Deficit



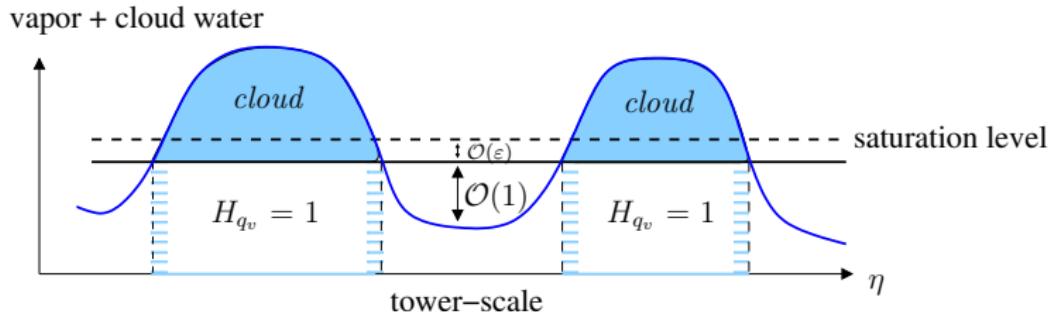
## Switching Function

$$H_{qv} := \begin{cases} 1 & : \text{Saturation at } \mathcal{O}(1) \quad (\text{Cloud}) \\ 0 & : \text{Under-saturation} \quad (\text{No cloud}) \end{cases}$$

## Saturated Area Fraction

$$\sigma := \overline{H_{qv}}^\eta, \quad \text{for } \delta q_v = \mathcal{O}(1) \text{ due to short time-scale: } D_\tau \sigma = 0.$$

# Leading Order Saturation Deficit



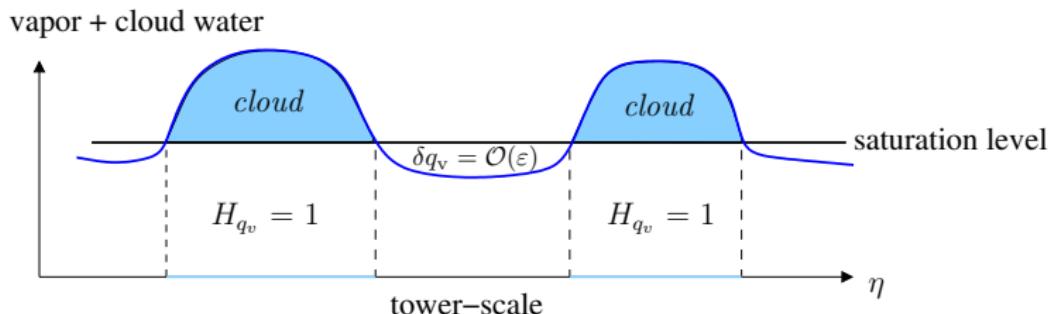
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# Systematically Small Saturation Deficit



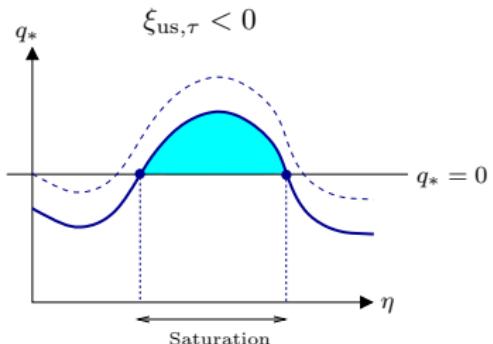
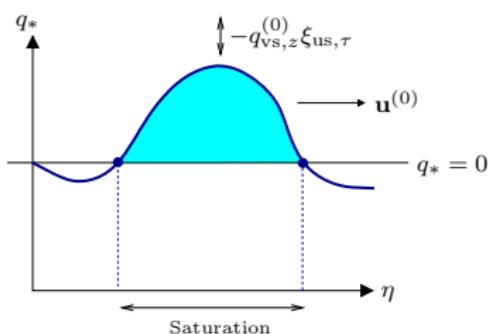
## Expansion of Vapor Mixing Ratio

$$q_v = \underbrace{q_{vs}^{(0)}}_{\text{saturation}} + \varepsilon q_v^{(1)} + \mathcal{O}(\varepsilon^2)$$

## Switching Function

$$H_{q_v} := \begin{cases} 1 & : \text{Saturation at } \mathcal{O}(\varepsilon) \\ 0 & : \text{Under-saturation} \end{cases}$$

# Tracking Saturated Areas



## Shape of Level Set Function

$$q_*(\eta, \tau) = q_{*,0}(\eta - \int_0^\tau \mathbf{u}^{(0)} d\tau') - q_{vs,z}^{(0)} \xi_{us}(\tau), \text{ Assume: } \xi_{us}(\tau) \approx \xi(\eta, \tau)$$

$$\text{Hence : } q_*(\eta, \tau) = 0 \Leftrightarrow q_{*,0}(\eta - \int_0^\tau \mathbf{u}^{(0)} d\tau') = q_{vs,z}^{(0)} \xi_{us}(\tau)$$

# Model Equations

Closed model equations:

Wave-scale dynamics:

$$\mathbf{u}_\tau^{(0)} + \mathbf{u}^\infty \cdot \nabla_x \mathbf{u}^{(0)} + \nabla_x \pi = 0$$

$$\bar{w}_\tau^{(0)} + \mathbf{u}^\infty \cdot \nabla_x \bar{w}^{(0)} + \pi_z = \bar{\theta}^{(3)}$$

$$\bar{\theta}_\tau^{(3)} + \mathbf{u}^\infty \cdot \nabla_x \bar{\theta}^{(3)} + (1 - \sigma) \Theta_z^{(2)} \bar{w}^{(0)} = \Theta_z^{(2)} w'$$

$$\nabla_x \cdot (\rho^{(0)} \mathbf{u}^{(0)}) + (\rho^{(0)} \bar{w}^{(0)})_z = 0$$

Effective tower-scale dynamics:

$$w'_\tau + \mathbf{u}^\infty \cdot \nabla_x w' + \frac{\sigma_\tau}{1 - \sigma} w' = \theta'$$

$$\theta'_\tau + \mathbf{u}^\infty \cdot \nabla_x \theta' + \sigma \Theta_z^{(2)} w' + \frac{\sigma_\tau}{1 - \sigma} \theta' = \sigma (1 - \sigma) \Theta_z^{(2)} \bar{w}$$

$$\sigma_\tau = \xi_{\text{us},\tau} \Psi(\xi_{\text{us}})$$

$$\xi_{\text{us},\tau} = \bar{w} - \frac{w'}{1 - \sigma}$$

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$$\bar{\theta}_\tau^{(3)} + \mathbf{u}^\infty \cdot \nabla_x \bar{\theta}^{(3)} + \Theta_z^{(2)} \bar{w}^{(0)} = 0$$

$$\nabla_x \cdot \left( \rho^{(0)} \mathbf{u}^{(0)} \right) + \left( \rho^{(0)} \bar{w}^{(0)} \right)_z = 0$$

Effective tower-scale dynamics:

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Effective tower-scale dynamics:

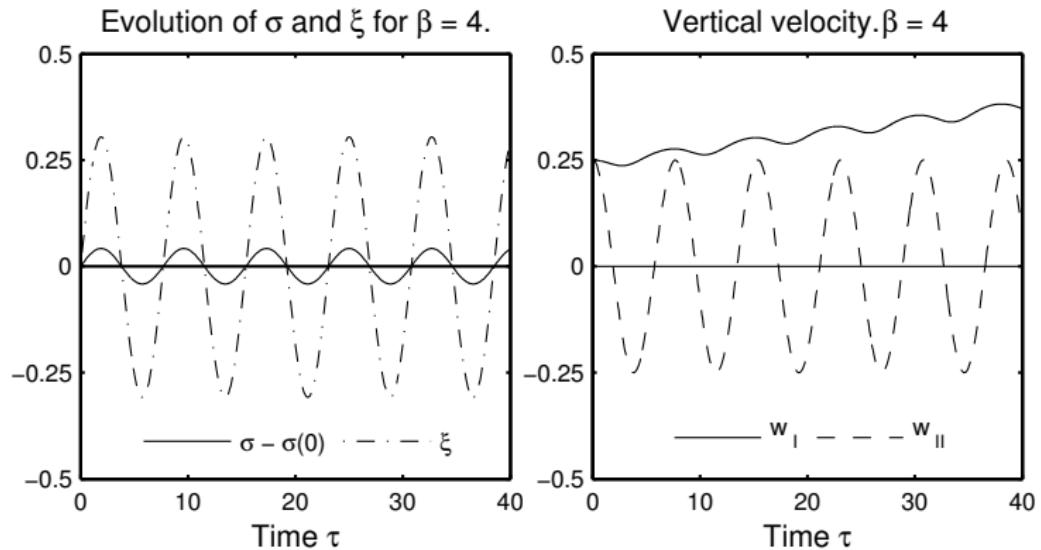
$$w'_\tau + \mathbf{u}^\infty \cdot \nabla_x w' + \frac{\sigma_\tau}{1 - \sigma} w' = \theta'$$

$$\theta'_\tau + \mathbf{u}^\infty \cdot \nabla_x \theta' + \sigma \Theta_z^{(2)} w' + \frac{\sigma_\tau}{1 - \sigma} \theta' = \sigma (1 - \sigma) \Theta_z^{(2)} \bar{w}$$

$$\sigma_\tau = \xi_{\text{us},\tau} \Psi(\xi_{\text{us}})$$

$$\xi_{\text{us},\tau} = \bar{w} - \frac{w'}{1 - \sigma}$$

# Plane Wave Amplitudes



Amplitudes of plane wave

$$\phi(x, z, \tau) = \hat{\phi}(\tau) \exp(i kx + i m z)$$

# Summary

- Derived **reduced model** for gravity waves in deep convecting atmosphere
- Assuming **small saturation deficit** leads to **nonlinear interactions** between waves and clouds
- Tracking of saturated spots by **level set function**: Simplifying assumptions allow to derive simple evolution equation
- Final model is extension of linearized anelastic equations

# Literature

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Theoretical & Computational Fluid Dynamics, 20 (2006), pp. 525-551.
-  D. Ruprecht, R. Klein, A. J. Majda.  
*Modulation of Internal Waves in a Multi-scale Model for Deep Convection on Mesoscales.*  
J. Atmos. Sci., 67 (2010), pp. 2504–2519.
-  D. Ruprecht, R. Klein.  
*A Model for Nonlinear Interactions of Internal Gravity Waves with Saturated Regions.*  
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