



## $O_2(b^1\Sigma_g^+, v = 0, 1)$ Relative Yield in $O(^1D) + O_2$ Energy Transfer

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Energy transfer from excited  $O(^1D)$  atoms to ground-state  $O_2(X^3\Sigma_g^-)$  leads to production of  $O_2$  in the first two vibrational levels of the  $O_2(b^1\Sigma_g^+)$  state:  $O(^1D) + O_2 \rightarrow O(^3P) + O_2(b^1\Sigma_g^+, v = 0, 1)$ . Subsequent radiative decay of  $O_2(b^1\Sigma_g^+, v = 0, 1)$  to the ground state results in the Atmospheric Band emission, a prominent feature of the terrestrial airglow. The relative yield for production of  $O_2(b^1\Sigma_g^+, v = 0, 1)$  in the above process,  $k_1/k_0$ , is an important parameter in modeling of the observed  $O_2$  Atmospheric Band emission intensities.

In the laboratory experiments, the output of a pulsed fluorine laser at 157 nm is used to photodissociate molecular oxygen in an  $O_2/N_2$  mixture flowing through a heated gas cell. Photodissociation of  $O_2$  produces a ground-state  $O(^3P)$  atom and an excited  $O(^1D)$  atom.  $O(^1D)$  rapidly transfers energy to the remaining  $O_2$  to produce  $O_2(b^1\Sigma_g^+, v = 0, 1)$ . The populations of  $O_2(b^1\Sigma_g^+, v = 0, 1)$  are monitored by observing emissions in the  $O_2(b-X)$  0–0 and 1–0 bands at 762 and 688 nm, respectively. The value of  $k_1/k_0$  is extracted from the time-dependent  $O_2(b^1\Sigma_g^+, v = 0, 1)$  fluorescence signals using computer simulations. We find that production of  $v = 1$  is substantially larger than that of  $v = 0$ .

We will present measurements on  $k_1/k_0$  and its temperature dependence, and discuss the significance of these and other relevant laboratory measurements on the interpretation of the  $O_2$  Atmospheric Band emission.

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