

Soundproof simulations of stratospheric gravity waves on unstructured meshes

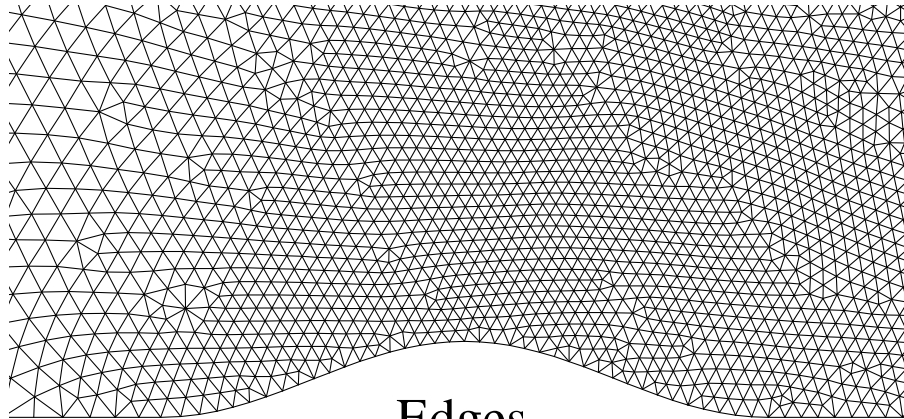
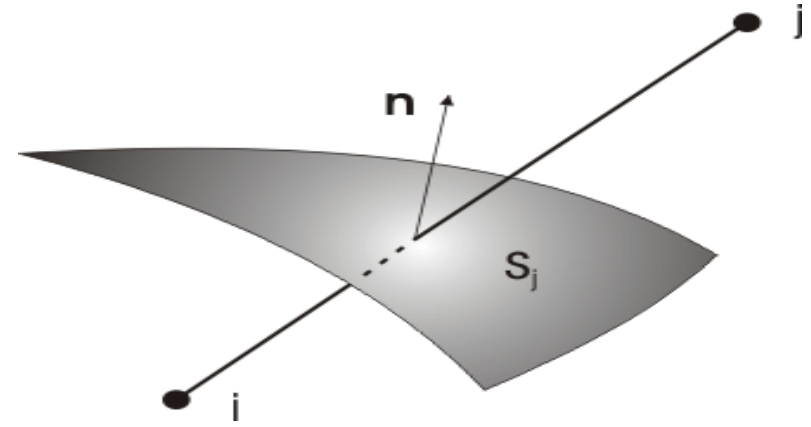
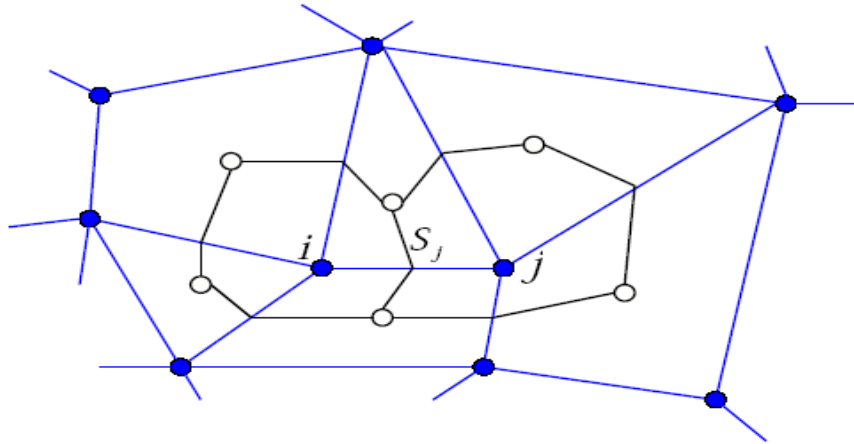
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Aims:

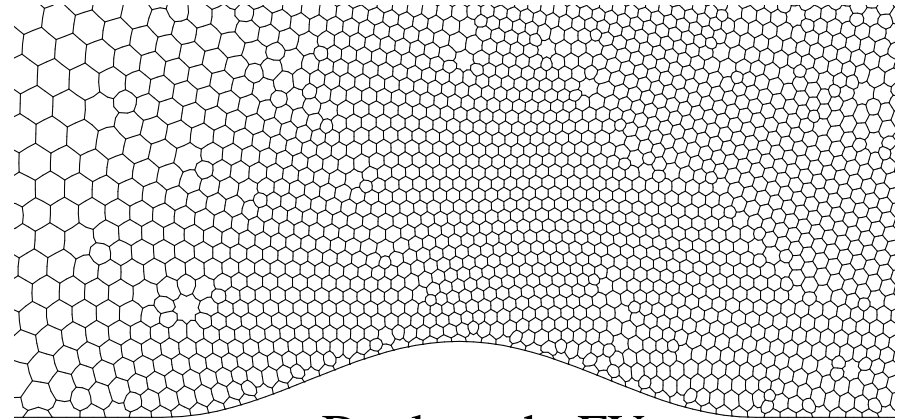
- a) highlight the progress with the development of a nonhydrostatic, soundproof, unstructured-mesh model for atmospheric flows;
- b) assess the accuracy of unstructured-mesh discretization relative to equivalent structured-grid methods for wave dynamics;
- c) assess relative merits of anelastic and pseudo-incompressible PDEs versus fully compressible PDEs, debated in the literature

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Nonoscillatory forward-in-time semi-implicit numerics (of EULAG) with flexible edge-based finite-volume (FV) spatial discretisation



Edges



Dual mesh, FVs

See Szmelter & Smolarkiewicz, *Comp. Fluids* (2011), for our FV solutions to canonical problems of hydrostatic/nonhydrostatic, strongly/weakly stratified, and rapidly/slowly rotating mountain flows.

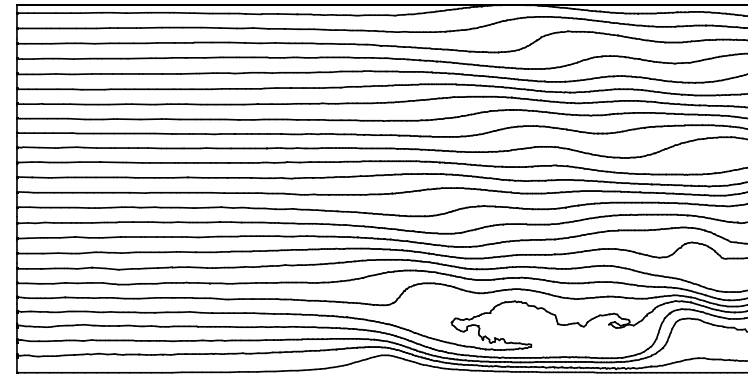
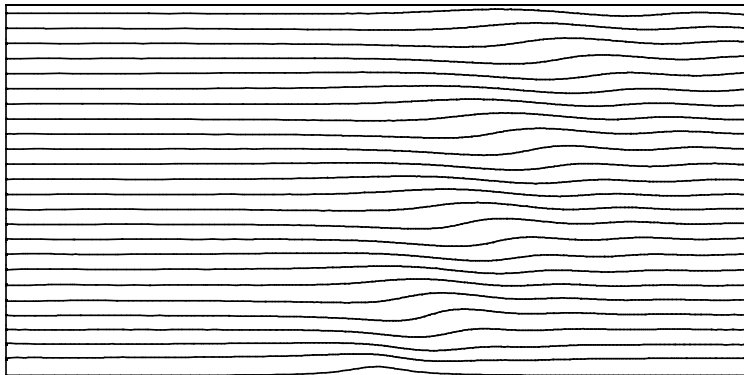
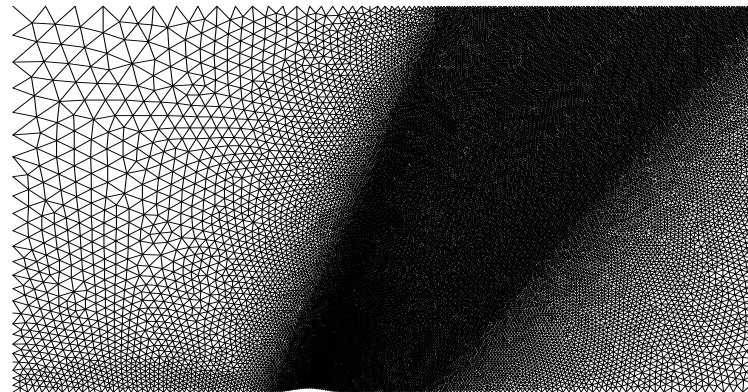
Nonhydrostatic Boussinesq mountain wave

Szmelter & Smolarkiewicz , *Comp. Fluids*, 2011

$$\nabla \bullet (\mathbf{V} \rho_o) = 0 ,$$

$$\frac{\partial \rho_o V^I}{\partial t} + \nabla \bullet (\mathbf{V} \rho_o V^I) = -\rho_o \frac{\partial \tilde{p}}{\partial x^I} + g \rho_o \frac{\theta'}{\theta_o} \delta_{I2}$$

$$\frac{\partial \rho_o \theta}{\partial t} + \nabla \bullet (\mathbf{V} \rho_o \theta) = 0 .$$



$$Fr \lesssim 2$$

$$NL/U_o = 2.4$$

$$Fr \lesssim 1,$$

Comparison with the EULAG's results and the linear theories (Smith 1979, Durran 2003):
3% in wavelength; 8% in propagation angle; wave amplitude loss 7% over 7 wavelengths

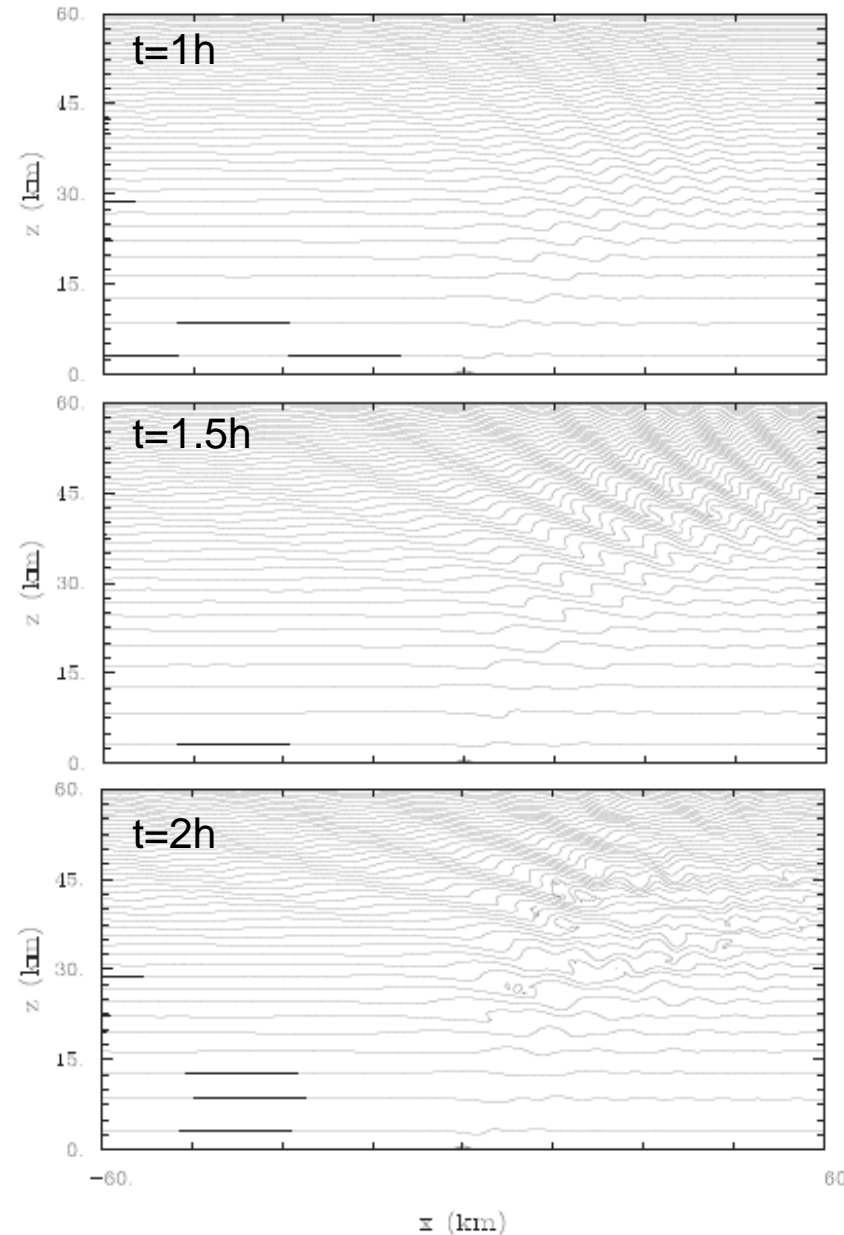
Non-Boussinesq amplification and breaking of deep stratospheric gravity wave

isothermal reference
profiles ; $H_\theta = 3.5 H_\rho$

$$NL/U_o \approx 1, Fr \approx 1.6;$$
$$\lambda_o = 2\pi \text{ km} \lesssim H_\rho \Rightarrow$$
$$A(H/2) = 10h_o = \lambda_o$$

EULAG “reference” solution
using terrain-following →
coordinates

Prusa et al., *JAS* 1996;
Smolarkiewicz & Margolin, *Atmos. Ocean*, 1997;
Klein, *Ann. Rev. Fluid Dyn.*, 2010



Soundproof generalizations

Smolarkiewicz & Szmelter, *Acta Geophysica* (2011)

$$\nabla \cdot (\rho^* \mathbf{v}) = 0, \quad \frac{D\theta'}{Dt} = -\mathbf{v} \cdot \nabla \theta_e, \quad \frac{D\mathbf{v}}{Dt} = -\Theta \nabla \phi' - \mathbf{g} \Upsilon \frac{\theta'}{\bar{\theta}}$$

For [anelastic, pseudo-incompressible]:

$$\rho^* = [\bar{\rho}, \bar{\rho} \bar{\theta} / \theta_o]; \quad \Theta = [1, \theta / \theta_o]; \quad \text{and} \quad \Upsilon = [1, \bar{\theta} / \theta_e]$$

$$\theta' = \theta - \theta_e \quad \phi' = [(p - p_e) / \bar{\rho}, c_p (\pi - \pi_e) \theta_o]$$

$$0 = -\nabla \frac{p_e - \bar{p}}{\bar{\rho}} - \mathbf{g} \frac{\theta_e - \bar{\theta}}{\bar{\theta}} \quad 0 = -c_p \theta_e \nabla (\pi_e - \bar{\pi}) - \mathbf{g} \frac{\theta_e - \bar{\theta}}{\bar{\theta}}$$

$$\frac{\partial \rho^* \psi}{\partial t} + \nabla \cdot (\rho^* \mathbf{v} \psi) = \rho^* R \quad \rightarrow \quad \psi_i^{n+1} = \mathcal{A}_i(\tilde{\psi}, \mathbf{v}^{n+1/2}, \rho^*) + 0.5 \delta t R_i^{n+1} \equiv \widehat{\psi}_i + 0.5 \delta t R_i^{n+1}$$

$$u = \hat{u} - \delta_h t \Theta \partial_x \phi' - \delta_h t \alpha (u - u_e)$$

$$w = \hat{w} - \delta_h t \Theta \partial_z \phi' + \delta_h t \beta \theta' - \delta_h t \alpha (w - w_e)$$

$$\theta' = \hat{\theta}' - \delta_h t w \partial_z \theta_e - \delta_h t \alpha' \theta' ,$$

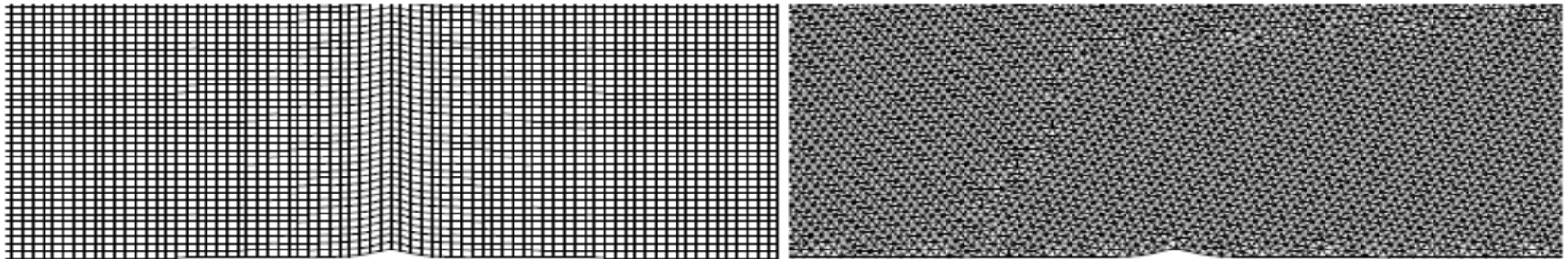
$$\begin{aligned} u &= \hat{u} - C^{xx} \partial_x \phi' \\ w &= \hat{w} - C^{zz} \partial_z \phi' \end{aligned} \quad \theta' = \frac{\hat{\theta}' - \delta_h t w \partial_z \theta_e}{1 + \delta_h t \alpha'}$$

$$\hat{u} = \frac{\hat{u} + \delta_h t \alpha u_e}{1 + \delta_h t \alpha} , \quad \hat{w} = \frac{\hat{w} + \delta_h t \alpha w_e + \delta_h t \beta \hat{\theta}' (1 + \delta_h t \alpha')^{-1}}{(1 + \delta_h t \alpha)(1 + \delta_h t \alpha') + (\delta_h t)^2 \beta \partial_z \theta_e (1 + \delta_h t \alpha')^{-1}}$$

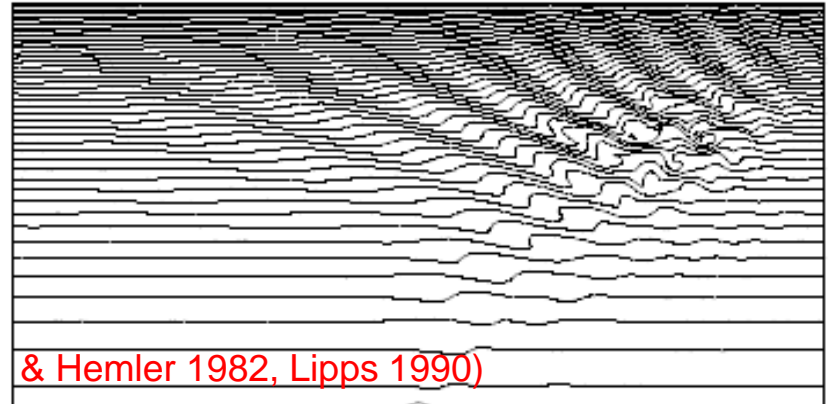
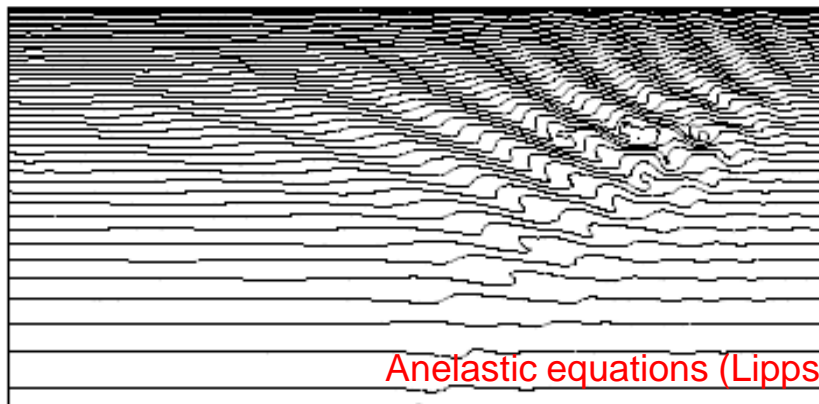
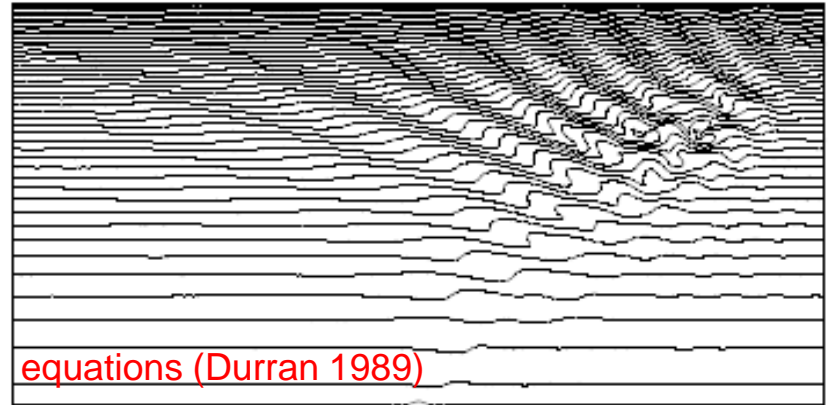
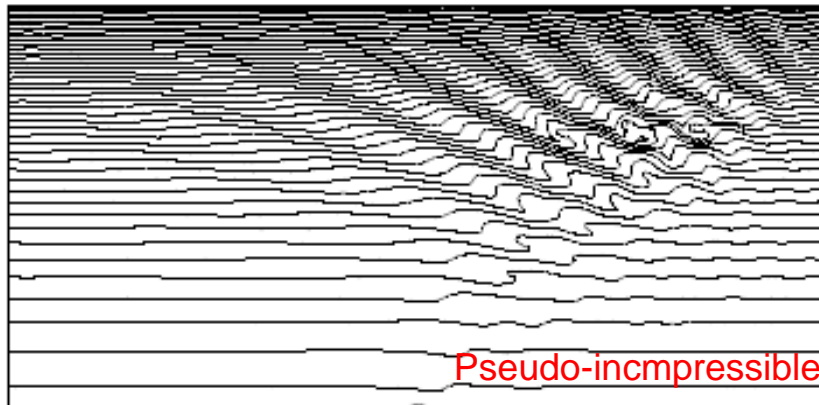
$$C^{xx} = \frac{\delta_h t \check{\Theta}}{(1 + \delta_h t \alpha)} , \quad C^{zz} = \frac{\delta_h t \check{\Theta}}{(1 + \delta_h t \alpha)(1 + \delta_h t \alpha') + (\delta_h t)^2 \beta \partial_z \theta_e (1 + \delta_h t \alpha')^{-1}}$$

$$\frac{1}{\rho^*} \left[\partial_x \rho^* (\hat{u} - C^{xx} \partial_x \phi') + \partial_z \rho^* (\hat{w} - C^{zz} \partial_z \phi') \right] \equiv -(\mathcal{L} \phi' - \mathcal{R}) = 0$$

unstructured-mesh “EULAG” solution, using mesh (left) mimicking terrain-following coordinates, and (right) a fully unstructured mesh



$t=1.5h$



$$H_{\theta} \lesssim H_{\rho} \quad (\text{RE: Achatz et al., } JFM, 2010)$$

t=0.75h



NCAR

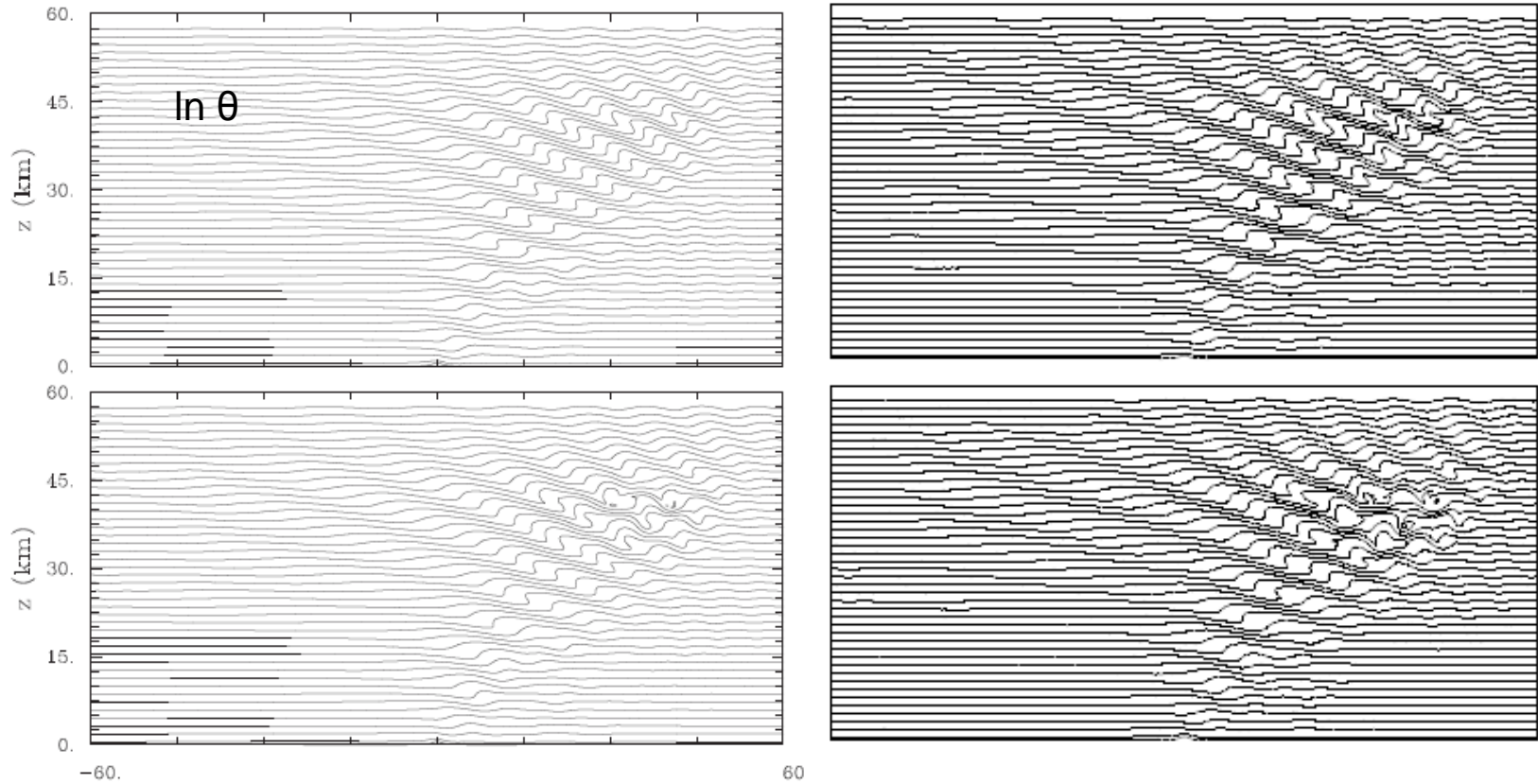


Table 1: Normalized vorticity: maximum, minimum, mean and standard deviation.

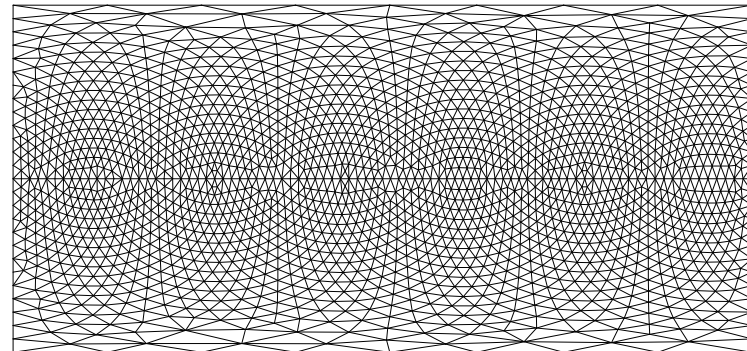
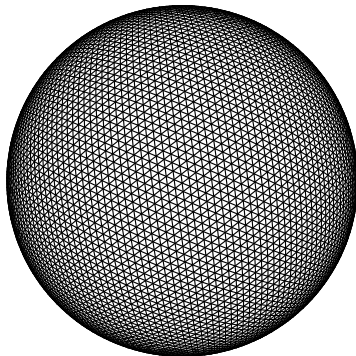
<i>eqs.</i>	<i>numerics</i>	$\max(\omega)$	$\min(\omega)$	$\bar{\omega}$	$\overline{(\omega - \bar{\omega})^2}^{1/2}$
PSI	CV/grid	0.17	-0.21	$1.6 \cdot 10^{-4}$	$3.3 \cdot 10^{-2}$
ANL	CV/grid	0.27	-0.41	$6.4 \cdot 10^{-5}$	$3.5 \cdot 10^{-2}$
PSI	CV/mesh	0.28	-0.24	$2.0 \cdot 10^{-4}$	$3.7 \cdot 10^{-2}$
ANL	CV/mesh	0.24	-0.36	$9.5 \cdot 10^{-5}$	$3.6 \cdot 10^{-2}$
PSI	SL/grid	0.28	-0.30	$2.1 \cdot 10^{-4}$	$3.1 \cdot 10^{-2}$
ANL	SL/grid	0.18	-0.24	$7.2 \cdot 10^{-5}$	$3.0 \cdot 10^{-2}$

Unstructured-mesh framework for atmospheric flows

Smolarkiewicz & Szmelter, pubs in *JCP*, *IJNMF*, *Comp. Fluids*, 2005-2011

- Differential manifolds formulation $\frac{\partial G\Phi}{\partial t} + \nabla \cdot (V\Phi) = G\mathcal{R}$, $V(x, t) := G\dot{x}$
- Finite-volume NFT numerics with a fully unstructured spatial discretization, heritage of EULAG and its predecessors ($A = MPDATA$)

$$\Phi_i^{n+1} = \mathcal{A}_i(\Phi^n + 0.5\delta t \mathcal{R}^n, V^{n+1/2}, G) + 0.5\delta t \mathcal{R}_i^{n+1}$$

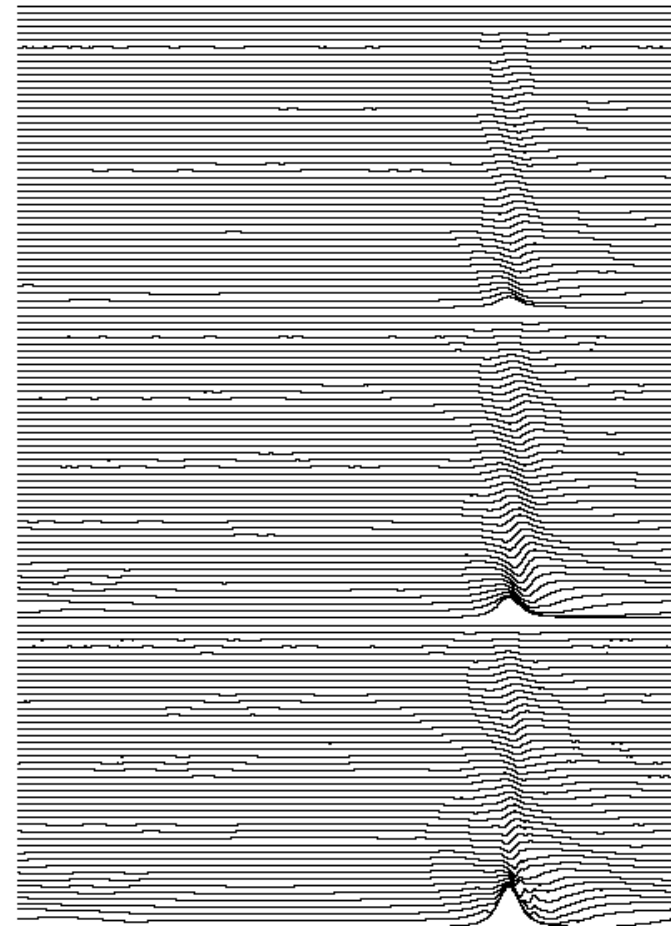
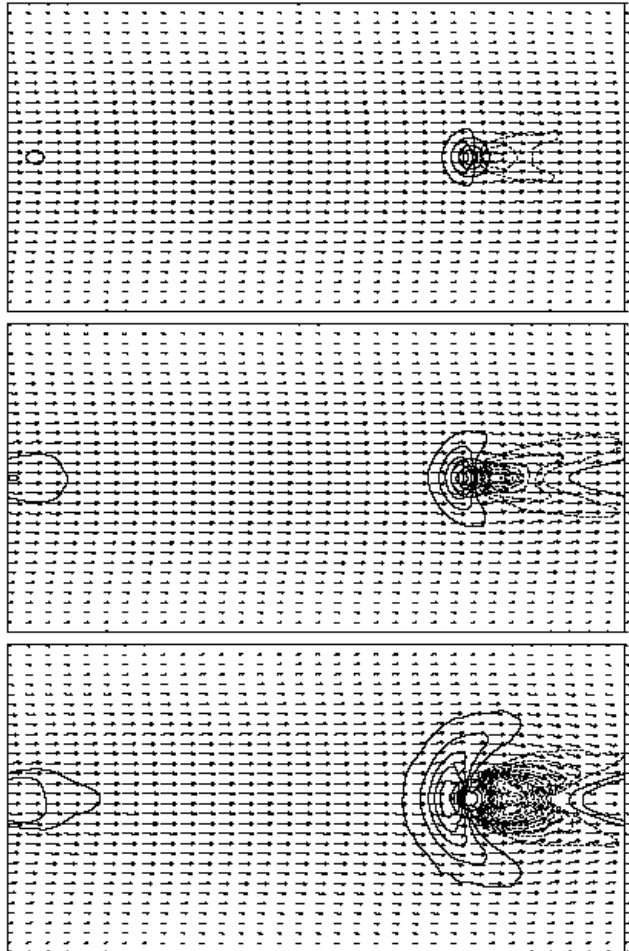


- Focus (so far) on wave phenomena across a range of scales and Mach, Froude & Rossby numbers

Stratified (mesoscale) flow past an isolated hill on a reduced planet

4 hours

$$Fr = U_0/Nh$$



Fr=2

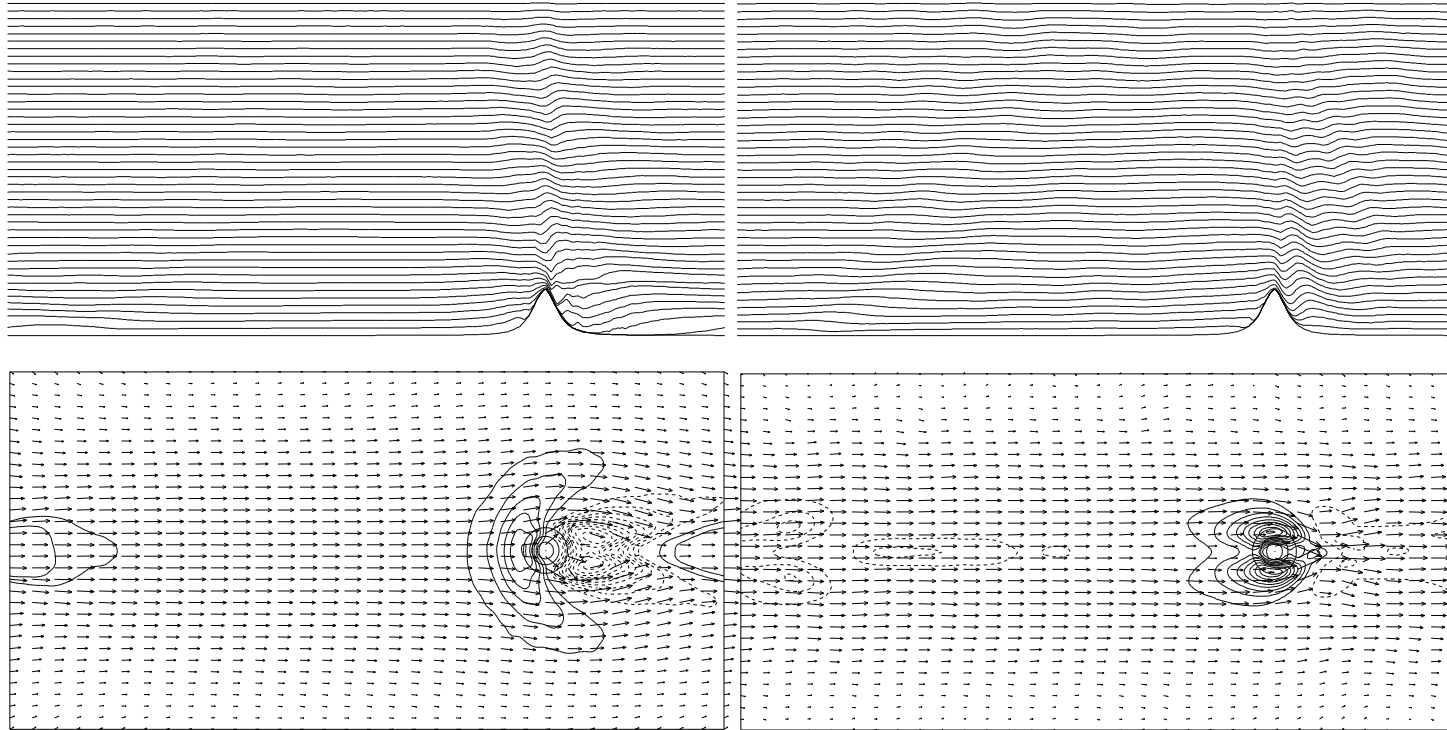
Fr=1

Fr=0.5

$$Fr=0.5$$

$$Ro \gg 1$$

$$Ro \gtrsim 1$$



Conclusions:

Unstructured-mesh discretization sustains the accuracy of structured-grid discretization and offers full flexibility in spatial resolution.

Soundproof models appear effective for a range of atmospheric flows much broader than thought the decade ago, while having considerable advantages over fully compressible models.