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# **Computing streamfunction and velocity potential** in a limited domain of arbitrary shape

**Jie CAO** IAP, Chinese Academy of Sciences, Beijing, CHINA; CIMMS, University of Oklahoma, USA Qin XU NOAA/National Severe Storms Laboratory, Norman, USA

# 1. Introduction

Streamfunction and velocity potential,  $\psi$  and  $\chi$ , are useful scalars of flow fields and powerful tools for studies of fluid dynamics. They provide useful tools for weather and climate map analyses, for deriving velocity error covariance representations in large-scale data assimilation. **However** the direct usage of  $\psi$  and  $\chi$  in the basic dynamics of the rotational and divergent components of the global or regional atmospheric circulation has not been well studied. Main reason lies on the lack of ways to compute  $\psi$  and  $\chi$  precisely and conveniently.

## 3. Numerical methods & Exps.

 $\succ$  Discretization schemes for  $(\chi_i, \psi_i)$ : to deal with the singular point P.



**S1**: the solution and source fields are of the same resolution; but staggered by  $\frac{1}{2} \triangle d$ .

**S2**: the same grids; but the grid is nested locally around P into P<sub>1,2,3,4</sub>.

**S3**: similar to S2, except that the nested discrete integrations around P are replaced by piecewise continuous integration.

**Problems to be solved:**  $\Delta \chi = \alpha, \Delta \psi = \zeta$ , inside the domain;  $\partial_n \chi - \partial_s \psi = v_n$ ,  $\partial_s \chi + \partial_n \psi = v_s$ , along the boundary.

to explore the utilities of classic integral formulae for **Objective:** computing  $\psi$  and  $\chi$ , and to design efficient numerical methods to accurately compute  $\psi$  and  $\chi$  in both regular and irregular domains either in the presence or absence of data hole.

### 2. Theory & Integral formulae

**Method:** to minimize the difference between the domainintegrated kinetic energy of the original horizontal velocity and that of the reconstructed one, i.e.,  $J = \int_D d\mathbf{x} |\nabla \boldsymbol{\chi} + \mathbf{k} \times \nabla \boldsymbol{\psi} - \mathbf{v}|^2$ .

 $(\chi, \psi) = (\chi_i, \psi_i) + (\chi_e, \psi_e) + (\chi_d, \psi_d)$ 

> The internally induced solution

 $\chi_{i} = (2\pi)^{-1} \int_{D} dx' \alpha(x') ln |x-x'|; \psi_{i} = (2\pi)^{-1} \int_{D} dx' \zeta(x') ln |x-x'|.$ (1) is solved in unbounded domain, therefore, no B.C. (2) satisfies:  $\Delta \chi_i = \alpha$ , and  $\Delta \psi_i = \zeta$ .



 $\succ$  Discretization schemes for  $(\chi_e, \psi_e)$ : to reduce the error near corners when solve for the density function equation.

**x**, (or **x**,,)



**SD1**: Dirichlet, uniform-boundary grid. **SD2**: Dirichlet, local-nesting. **SD3**: Dirichlet, piecewise continuous integration. **SN2**: Neumann, local-nesting.

**Exp 2:**  $\psi_e^{\text{true}} = -8xy(x^2 - y^2) - 3(3x^2 - y^2)y - 8xy - 5y$ ... (1)  $\psi_{e}^{\text{true}} = \sum_{m} [a_{m} \sin(m\pi x/L) \sin(m\pi y/L) + b_{m} \sin(m\pi x/L) \sin(m\pi y/L)] \dots (2)$ 

$\psi_{ m e}^{ m true}$	Scheme	SCC	RRD
(1)	SD1	1.000	0.485%
	SD2	1.000	0.015%
	SD3	1.000	0.11%
	SN2	0.981	20.7%
(2)	SD1	1.000	0.822%
	SD2	1.000	0.060%
	SD3	1.000	0.25%
	SN2	0.728	70.5%

### The data-hole induced solution

 $\chi_{d} = \sum_{1} (2\pi)^{-1} ln |\mathbf{x} - \mathbf{x}_{k}| \int_{Sk} ds(v_{n} - v_{ni}); \ \psi_{d} = \sum_{1} (2\pi)^{-1} ln |\mathbf{x} - \mathbf{x}_{k}| \int_{Sk} ds(v_{si} - v_{si}).$ (1) is constructed by placing point sources of  $\alpha$  and  $\zeta$  inside each data hole to represent the net effects of the integrated flux and circulation around the data hole.

(2) should be subtracted from the total solution to modify B.C. and solvability conditions for  $(\chi_e, \psi_e)$ .

> The externally induced solution

**Approach 1**:  $(\chi_e, \psi_e) = (0, -Im\omega_e)$ , where the complex velocity potential  $\omega_{e} = \varphi_{e} - i \psi_{e} = (i2\pi)^{-1} \int_{S} dz' \omega_{e}(z')/(z'-z), \ \omega_{e}(s) = \omega_{e}(s_{0}) + \int_{0}^{s} ds' [(v_{s} - v_{s})]$ +  $i(v_n - v_{ni})$ ] on the boundary *S*, and  $\partial_n \psi_e = \partial_s \varphi_e$  along *S*.

#### Approach 2:

**2.1** With Dirichlet B.C. :

 $\psi_{e}(\mathbf{x}) = \int_{S} ds' \mu(\mathbf{x}') \mathbf{n}' \cdot \mathbf{r}/r^{2}$  for **x** in *D*, and the density function  $\mu(\mathbf{x}')$  satisfies:  $\psi_{\rho}(\mathbf{x}) = \pi \mu(\mathbf{x}) + \int_{S} ds' \mu(\mathbf{x}') \mathbf{n}' \cdot \mathbf{r}/r^2$  for  $\mathbf{x}$  on S.

2.2 With Neumann B.C. :

 $\psi_{\rho}(\mathbf{x}) = -\int_{S} ds' v(\mathbf{x}') \ln r$  for **x** over *D*, and the density function  $v(\mathbf{x}')$  satisfies:  $\partial_n \psi_e(\mathbf{x}) = \pi v(\mathbf{x}) + \int_S ds' v(\mathbf{x}') \mathbf{n} \cdot \mathbf{r}/r^2$  for  $\mathbf{x}$  on S.

### - Discretization schemes for $(\chi_d, \psi_d)$ : to format the B.C for integration.



**Exp 3**:







**2.3** With tangential-derivative B.C. :  $\psi_{e}(\mathbf{x}) = -\int_{S} ds' v(\mathbf{x}') \ln r$  for **x** over *D*, and  $v(\mathbf{x}')$  satisfies:  $-\int_{S} ds' v(\mathbf{x}') \mathbf{s} \cdot \mathbf{r}/r^2 = v_{ni} - v_n \text{ for } \mathbf{x} \text{ on } S.$ (1) This part contains the null space which leads to non-unique solution. (2)  $\chi_{\rm e}$  can be solved exactly the same as  $\psi_{\rm e}$ .

**Cons and Pros** when compare to other methods: ✓ The non-uniqueness of solution and compatibility between the coupled B.C. are well taken care of.

V This method can solve for limited domains of arbitrary shape with inner data holes.

**?** How to design numerical schemes to reduce discretization error? **?** How to design an automated algorithm to format the B.C. for the integrations along irregular external and internal boundaries?