



Spatial and temporal compact equations for water waves

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A one-dimensional potential flow of an ideal incompressible fluid with a free surface in a gravity field is the Hamiltonian system with the Hamiltonian:

$$H = \frac{1}{2} \int dx \int_{-\infty}^{\eta} |\nabla \phi|^2 dz + \frac{g}{2} \int \eta^2 dx$$

$\phi(x, z, t)$ - is the potential of the fluid, g - gravity acceleration, $\eta(x, t)$ - surface profile

Hamiltonian can be expanded as infinite series of steepness:

$$\begin{aligned} H &= H_2 + H_3 + H_4 + \dots \\ H_2 &= \frac{1}{2} \int (g\eta^2 + \psi \hat{k} \psi) dx, \\ H_3 &= -\frac{1}{2} \int \{(\hat{k}\psi)^2 - (\psi_x)^2\} \eta dx, \\ H_4 &= \frac{1}{2} \int \{\psi_{xx} \eta^2 \hat{k} \psi + \psi \hat{k} (\eta \hat{k} (\eta \hat{k} \psi))\} dx. \end{aligned} \quad (1)$$

where \hat{k} corresponds to the multiplication by $|k|$ in Fourier space, $\psi(x, t) = \phi(x, \eta(x, t), t)$. This truncated Hamiltonian is enough for gravity waves of moderate amplitudes and can not be reduced.

We have derived self-consistent compact equations, both spatial and temporal, for unidirectional water waves.

Equations are written for normal complex variable $c(x, t)$, not for $\psi(x, t)$ and $\eta(x, t)$.

Hamiltonian for temporal compact equation can be written in x -space as following:

$$H = \int c^* \hat{V} c dx + \frac{1}{2} \int \left[\frac{i}{4} (c^2 \frac{\partial}{\partial x} c^{*2} - c^{*2} \frac{\partial}{\partial x} c^2) - |c|^2 \hat{K}(|c|^2) \right] dx \quad (2)$$

Here operator \hat{V} in K-space is so that $V_k = \frac{\omega_k}{k}$. If along with this to introduce Gardner-Zakharov-Faddeev bracket (for the analytic in the upper half-plane function)

$$\partial_x^+ \Leftrightarrow ik\theta_k \quad (3)$$

Hamiltonian for spatial compact equation is the following:

$$H = \frac{1}{g} \int \frac{1}{\omega} |c_\omega|^2 d\omega + \frac{1}{2g^3} \int |c|^2 (\ddot{c}^* c + \ddot{c} c^*) dt + \frac{i}{g^2} \int |c|^2 \hat{\omega} (\dot{c} c^* - c \dot{c}^*) dt. \quad (4)$$

equation of motion is:

$$\begin{aligned} & \frac{\partial}{\partial x} c + \frac{i}{g} \frac{\partial^2}{\partial t^2} c = \\ & = \frac{1}{2g^3} \frac{\partial^3}{\partial t^3} \left[\frac{\partial^2}{\partial t^2} (|c|^2 c) + 2|c|^2 \ddot{c} + \ddot{c}^* c^2 \right] + \\ & + \frac{i}{g^3} \frac{\partial^3}{\partial t^3} \left[\frac{\partial}{\partial t} (c \hat{\omega} |c|^2) + \dot{c} \hat{\omega} |c|^2 + c \hat{\omega} (\dot{c} c^* - c \dot{c}^*) \right]. \end{aligned} \quad (5)$$

It solves the spatial Cauchy problem for surface gravity wave on the deep water.

Main features of the equations are:

- Equations are written for complex normal variable $c(x, t)$ which is analytic function in the upper half-plane
- Hamiltonians both for temporal and spatial equations are very simple
- It can be easily implemented for numerical simulation

The equations can be generalized for "almost" 2-D waves like KdV is generalized to KP.

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