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Earth gravity field modeling and relativistic measurements with laser-ranged satellites and the LARASE research program



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Summary

- The LARASE experiment and its goals
- Relativistic effects to be measured
- Systematic errors from the background gravitational field
- Some recent measurements of relativistic effects
- Conclusions and future work



The LARASE experiment and its goals

The LARASE goals:

- The LAser RAnged Satellites Experiment (LARASE) main goal is to provide accurate measurements for the gravitational interaction in the weak-field and slow-motion limit of General Relativity by means of a very precise laser tracking of geodetic satellites orbiting around the Earth (the two LAGEOS and LARES)
- Beside the quality of the tracking observations, guaranteed by the powerful Satellite Laser Ranging (SLR) technique of the International Laser Ranging Service (ILRS), also the quality of the dynamical models implemented in the Precise Orbit Determination (POD) software plays a fundamental role in order to obtain precise and accurate measurements
- The models have to account for the perturbations due to both gravitational and non-gravitational forces in such a way to reduce as much as possible the difference between the *observed* range, from the tracking, and the *computed* one, from the models
- In particular, LARASE aims to improve the dynamical models of the current best laser-ranged satellites in order to perform a precise and accurate orbit determination, able to benefit also space geodesy and geophysics



The LARASE experiment and its goals

The LARASE activities:

- 1. Review of the literature, technical notes and all the documentation (NASA, ALENIA, ASI) related with the structure of the satellites and their physical characteristics
- 2. A reconstruction of the internal and external structure of the satellites with finite elements techniques
- 3. Review of the spin model of the two LAGEOS satellites and of their complex interaction with the Earth's magnetic field
- 4. Develop a spin model for LARES
- 5. Extension of the Yarkovsky–Schach thermal effect to the low spin-rate approximation
- 6. Impact of the neutral drag on the two **LAGEOS** satellites and on **LARES**
- 7. Solid and Ocean tides on the two **LAGEOS** satellites and on **LARES**
- 8. Precise Orbit Determination for the two LAGEOS satellites and for LARES
- 9. Finite Element Model for the Thermal effects





LAser **RE**lativity **S**atellite

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e

Ι

R

Μ

A/M

The LARASE experiment and its goals

LAGEOS, LAGEOS II and LARES

orbit, size, mass and materials

and LARES d materials	
LAGEOS II	LAser GEOdynamic Satellite LAGEOS II
12 163	LAGEOS (NASA 1976)
0.014	LAGEOS (NASA 1970) LAGEOS II (NASA/ASI 1992)

Par	ameter	LARES	LAGEOS	LAG	EOS II			LAGEOS II
	[km]	7 820	12 270	12 1	163		LAGEOS	NASA 1976)
		0.001	0.004	0.0	014		LAGEOS I	I (NASA/ASI 1992
	[deg]	69.5	109.8	52	2.7		LARES (A	SI 2012)
	[cm]	18.2	30	3	0			
	[kg]	386.8	406.9	40	5.4			
1	[m2/kg]	2.69.10-4	6.94·10 ⁻⁴	6.9	7.10-4		LARES	LAGEOS
	A	1 A			material		Tungsten	Al/Brass/Be/Cu
	$\frac{1}{M}$	$\cong \frac{1}{2.6} \frac{1}{M}$			CCR (suprasil	311)	92	422 + 4
		Lares	Lageos		bin		30 s	120 s

The LARASE experiment and its goals

- Despite the smaller A/M ratio, the nongravitational accelerations are not always smaller in magnitude for LARES with respect to LAGEOS II (or LAGEOS), due to the lower height (1450 vs. 5900 km) and the higher density of neutral atmosphere
- Being 50 times larger on LARES than on the two LAGEOS, the accurate modeling of neutral atmosphere drag needs special attention, because it might mask the presence of smaller and subtler effects

Effect	Estimate	LAGEOS II	LARES
Earth's monopole	$\frac{GM_{\oplus}}{r^2}$	2.69	6.51
Earth's oblateness	$3\frac{GM_{\oplus}}{r^2}\left(\frac{R_{\oplus}}{r}\right)^2 \bar{C}_{2,0}$	-1.1×10^{-3}	-6.4×10^{-3}
Low-order geopotential harmonics	$3\frac{GM_{\oplus}}{r^2}\left(\frac{R_{\oplus}}{r}\right)^2\bar{C}_{2,2}$	5.4×10^{-6}	3.2×10^{-5}
High-order geopotential harmonics	$19\frac{GM_{\oplus}}{r^2} \left(\frac{R_{\oplus}}{r}\right)^{18} \bar{C}_{18,18}$	1.4×10^{-12}	4.6×10^{-9}
Moon perturbation	$2\frac{GM}{r^3}r$	2.2×10^{-6}	1.4×10^{-6}
Sun perturbation	$2\frac{G\dot{M}_{\odot}}{r_{\odot}^3}r$	9.6×10^{-7}	6.2×10^{-7}
General relativistic correction	$\frac{GM_{\oplus}}{r^2} \frac{\overset{r_{\odot}}{G}M_{\oplus}}{\overset{1}{c^2}r}$	9.8×10^{-10}	3.7×10^{-9}
Atmospheric drag	$\frac{1}{2}C_D\frac{A}{M}\rho V^2$	-2.6×10^{-13}	-1.3×10^{-11}
Solar radiation pressure	$C_R \frac{A}{M} \frac{\Phi_{\odot}}{c}$	3.2×10^{-9}	1.2×10^{-9}
Albedo radiation pressure	$C_R \frac{A}{M} \frac{\Phi_{\odot}}{c} A_{\bigoplus} \left(\frac{R_{\oplus}}{r}\right)^2$	3.5×10^{-10}	2.4×10^{-10}
Thermal emission	$\frac{4}{9}\frac{A}{M}\frac{\Phi_{\odot}}{c}lpha \frac{\Delta T}{T_{0}}$	2.8×10^{-11}	not available
Dynamic solid tide	$3k_2 \frac{GM}{r} \left(\frac{R_{\oplus}}{r}\right)^2 \frac{R_{\oplus}^3}{r^4}$	3.7×10^{-6}	2.2×10^{-5}
Dynamic ocean tide	~ 0.1 of the dynamic solid tide	3.7×10^{-7}	2.2×10^{-6}

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Our main goals in the field of fundamental physics measurements fall in the following main targets:

- Schwarzschild precession (gravitoelectric field)
- Lense-Thirring precession (gravitomagnetic field)
- Geodetic (de Sitter) precession
- Post-Newtonian parameter (β, γ, α1, α2, ...)
- Constraints and limits to alternative theories of the gravitational interaction (Yukawa, non-symmetric/torsional ...)

We are now ready to start new refined measurements of the above relativistic effects with laser-ranged satellites. As said, there are two main aspects to satisfy:

- 1. obtain very precise measurements from the analysis of the post-fit residuals (after the POD)
- 2. provide a very reliable estimate of the systematics, i.e., accurate measurements

Gravito-electromagnetism: linearized theory of General Relativity (GR)

In the Weak-Field and Slow-Motion (WFSM) limit of the theory of GR, Einstein's equations reduce to a form quite similar to those of electromagnetism. Following this approach we have a:

- gravitoelectric field produced by masses, analogous to the electric field produced by charges
- gravitomagnetic field produced by mass currents, analogous to the magnetic field produced by electric currents.

$$G_{\alpha\beta} = 8\pi \frac{G}{c^4} T_{\alpha\beta}$$

$$\begin{bmatrix} \overline{h}^{\alpha\beta} \\ ,\beta \end{bmatrix} = 0$$

$$g_{\alpha\beta} = \eta_{\alpha\beta} + h_{\alpha\beta}$$

$$\Delta \overline{h}_{\alpha\beta} = 16 \pi \frac{G}{c^4} T_{\alpha\beta}$$

$$\begin{vmatrix} h \end{bmatrix}$$

$$\begin{cases} -\frac{1}{h_{\alpha\beta}} \equiv h_{\alpha\beta} - \frac{1}{2}\eta_{\alpha\beta} h\\ h \equiv h_{\alpha}^{\alpha} = \eta^{\alpha\beta}h_{\alpha\beta} \end{cases}$$
$$\begin{vmatrix} h_{\alpha\beta} \end{vmatrix} \cong \begin{vmatrix} \Phi\\ c^2 \end{vmatrix} \le 10^{-6}$$

$$\begin{cases} \overline{h}^{00} = 4 \frac{\Phi}{c^2} \\ \overline{h}^{0l} = -2 \frac{A^l}{c^2} \\ \overline{h}^{ij} = O(c^{-4}) \end{cases}$$

$$\eta_{\alpha\beta} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

Gravitoelectric potential

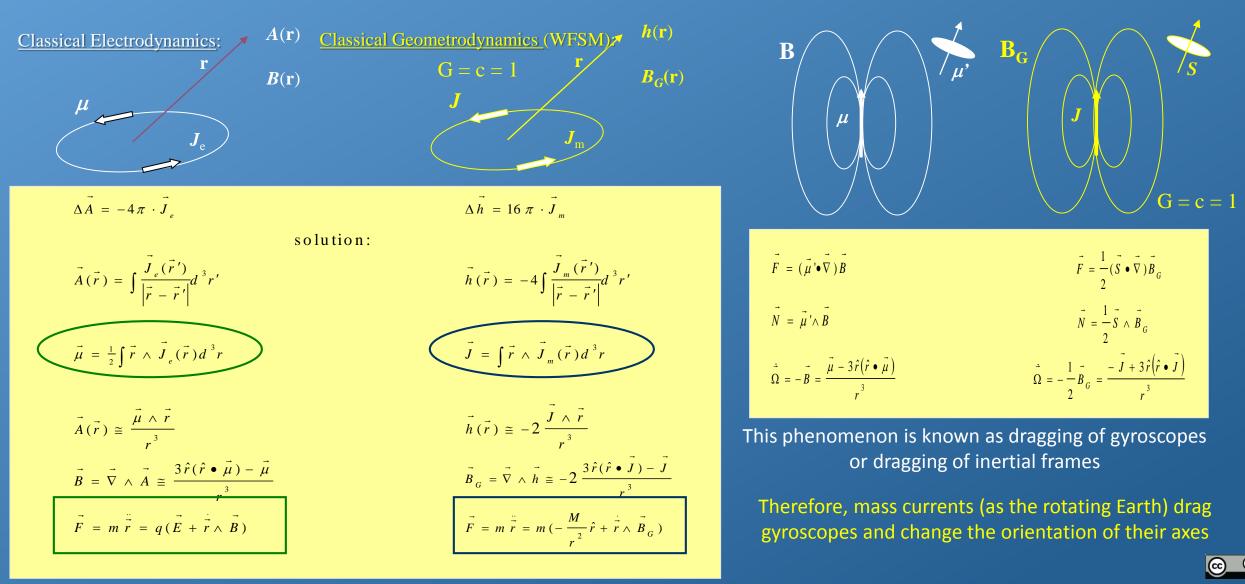
 $A^{l} = \frac{G}{2} \frac{J^{n} x^{k}}{3} \varepsilon^{l}_{nk}$

 GM_{\odot}

Gravitomagnetic potential



Formal analogy with electrodynamics: linearized theory of General Relativity (WFSM limit)



Gravitomagnetism

- Mass currents contribute to the curvature of spacetime
- Gravitomagnetism may be thought of as a manifestation of the way inertia originates in Einstein geometrodynamics ...
 "inertia here arises from mass there" ...
- The dragging of inertial frames or Lense-Thirring effect represents a weak manifestation (within GR) of Mach's Principle (the experimental proof of the origin of local inertial forces, interpreted as gravitational forces)
- The full inclusion of Mach Principle in GR is still debated ...
- Anyway, the astrophysical and cosmological consequences are very significant ...

See "Gravitation and Inertia", Ciufolini and Wheeler, 1995 for a deep insight into gravitomgnetism

Model for GR

Huang et al., Celest. Mech. & Dyn. Astron. 48, 1990

Relativistic perturbations $Gm_{\odot}\left[\left(\begin{array}{c}Gm_{\odot}\\Gm_{\odot}\end{array}\right)\right]$

$$\vec{A}_E = \frac{1}{c^2 r^3} \left[\left(4 \frac{1}{r} - v^2 \right) \vec{r} + 4(\vec{r} * \vec{v}) \vec{v} \right]$$
$$\vec{A}_{as} = 2 \left(\vec{\Omega} \wedge \vec{v} \right)$$

$$\vec{A}_{LT} = 2 \frac{Gm_{\oplus}}{c^2 r^3} \left[\frac{3}{r^2} (\vec{r} \wedge \vec{v}) (\vec{r} * \vec{J}) + (\vec{v} \wedge \vec{J}) \right]$$

Einstein or Schwarzschild component

De Sitter (or geodetic) component

Lense–Thirring component

ith:
$$\vec{\Omega} \approx \frac{3}{2} (\vec{V}_E - \vec{V}_S) \wedge \left(-\frac{GM_S}{c^2 R_{ES}^3} \right) \vec{X}_{ES}$$

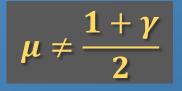
Where, capital letters refer to position, velocity, acceleration and mass in the barycentric reference frame, while small letters refer to the same quantities in the non–inertial geocentric reference system (E=Earth, S=Sun)

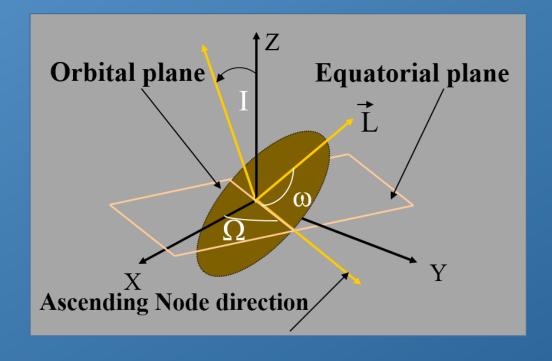
Gravitomagnetism: orbit precession

Lense-Thirring, Phys. Z, 19, 1918

$$\dot{\Omega}_{LT} = \mu \frac{2G}{c^2 a^3} \frac{J_{\oplus}}{(1 - e^2)^{3/2}}$$

$$\dot{\omega}_{LT} = -\mu \frac{6G}{c^2 a^3} \frac{J_{\oplus}}{(1 - e^2)^{3/2}} \cos I$$





mas/yr	LAGEOS	LAGEOS II	LARES
$\dot{\Omega}_{LT}$	30.7	31.5	118.5
$\dot{\omega}_{LT}$	31.2	- 57.3	- 124.5

30 mas/yr at LAGEOS altitude (≅5900 km) corresponds to a displacement of about 1.8 m/yr!

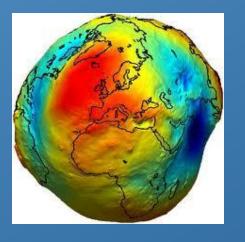
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In order to reach a precise and accurate measurement of a relativistic effect we need to perform a:

- precise orbit determination (POD) of the orbit of the satellites
 - observations (SLR data)
 - dynamical model (software)
- careful (robust) evaluation of the main systematic error sources
 - gravitational and non-gravitational
 - ad hoc analyses of the errors (formal and calibrated) of the various models considered





Big problem with the even zonal harmonics uncertainties: systematic errors

2

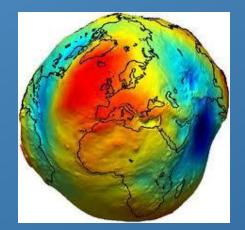
$$V(r,\varphi,\lambda) = -\frac{GM_{\oplus}}{r} \left[1 + \sum_{\ell=2}^{\infty} \sum_{m=0}^{\ell} \left(\frac{R_{\oplus}}{r} \right)^{\ell} P_{\ell_m}(\sin\varphi) (C_{\ell_m} \cos m\lambda + S_{\ell_m} \sin m\lambda) \right]$$
$$\prod_{m=0}^{\ell} m = 0 \quad \rightarrow \text{ zonal harmonics}$$
$$M_{\oplus} \left[1 + \frac{R_{\oplus}}{2} \frac{23\cos^2\vartheta - 1}{2\cos^2\vartheta - 1} \right] \quad \text{Departure}$$

r

Spherical harmonics development of the Earth's potential V(r)

Dependency from the even zonal harmonics only, their uncertainties mimics a secular effect in the right ascension of the node and also in the argument of pericenter

$$\delta \dot{\Omega}_{class} = -\frac{3}{2}n\left(\frac{R_{\oplus}}{a}\right)^2 \frac{\cos i}{(1-e^2)^2}\delta J_2 +$$



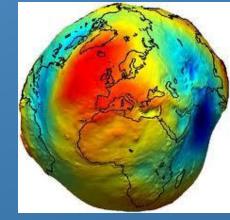
Big problem with the even zonal harmonics uncertainties: systematic errors

$$V(r,\varphi,\lambda) = -\frac{GM_{\oplus}}{r} \left[1 + \sum_{\ell=2}^{\infty} \sum_{m=0}^{\ell} \left(\frac{R_{\oplus}}{r} \right)^{\ell} P_{\ell_m}(\sin\varphi) (C_{\ell_m} \cos m\lambda + S_{\ell_m} \sin m\lambda) \right]$$
$$m = 0 \rightarrow \text{ zonal harmonics}$$

$$V(r) = -\frac{GM_{\oplus}}{r} \left[1 - J_2 \left(\frac{R_{\oplus}}{r}\right)^2 \frac{3\cos^2\vartheta - 1}{2} + \cdots \right]$$

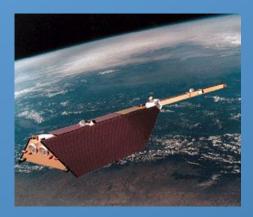
Dependency from the even zonal harmonics only, their uncertainties mimics a secular effect in the right ascension of the node and $\dot{\Omega}^{Class} \simeq -\frac{3}{2}n\left(\frac{R_{\oplus}}{a}\right)^2 \frac{\cos I}{(1-e^2)^2} \left\{J_2 + J_4\right\} \frac{5}{8} \left(\frac{R_{\oplus}}{a}\right)^2 \left(7\sin^2 I - 4\right) \frac{\left(1+\frac{3}{2}e^2\right)}{(1-e^2)^2} + \dots\right\}$ gument of pericenter

$$\dot{\omega}^{Class} \cong -\frac{3}{4}n\left(\frac{R_{\oplus}}{a}\right)^2 \frac{\left(1-5\cos^2 I\right)}{\left(1-e^2\right)^2} \left\{J_2 + J_4 \left[\frac{5}{256}\left(\frac{R_{\oplus}}{a}\right)^2 \left(7\sin^2 I - 4\right)\frac{C(e,I)}{\left(1-5\cos^2 I\right)}\right] + \dots\right\}$$



Spherical harmonics development of the Earth's potential V(r)

The CHAMP, GRACE and GOCE missions: high tech. space missions





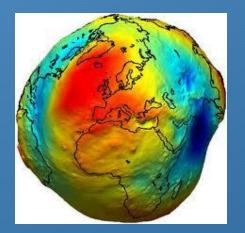


CHAMP carries:

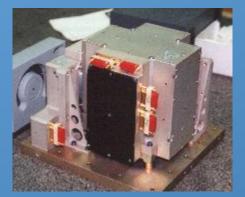
- GPS receiver for precise positioning, atmospheric profiling and bistatic altimetry;
- 2. STAR accelerometer for non–gravitational accelerations measurement
- 3. Advanced Stellar Compass for attitude control
- 4. magnetometers to measure the Earth's magnetic field;
- 5. laser retro-reflector array enabling laser ranging

GRACE carries:

- 1. Black–Jack GPS receiver for precise positioning
- 2. SuperStar accelerometer for non–gravitational accelerations measurement;
- 3. Star Tracker (ST) for attitude control;
- 4. K–Band Ranging (KBR) system;
- 5. laser retro–reflector array enabling laser ranging



The CHAMP, GRACE and GOCE missions: high tech. space missions



STAR accelerometer (ONERA/CNES, France)

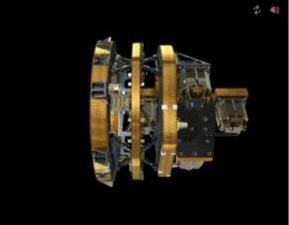


LASER Retroreflector

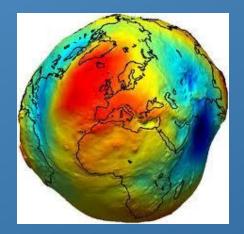
$$a \approx 10^{-9} m/s^2 / \sqrt{Hz} \quad (10^{-1} - 2 \cdot 10^{-4}) Hz$$
$$a \approx 10^{-10} m/s^2 / \sqrt{Hz} \quad (10^{-1} - 2 \cdot 10^{-4}) Hz$$



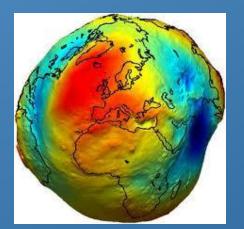
Multiple tracking



Gradiometer (EGG) $(10^{-1} - 5 \cdot 10^{-3})Hz$ $a \approx (10^{-12} - 10^{-10}) m/s^2/\sqrt{Hz}$



The CHAMP, GRACE and GOCE missions



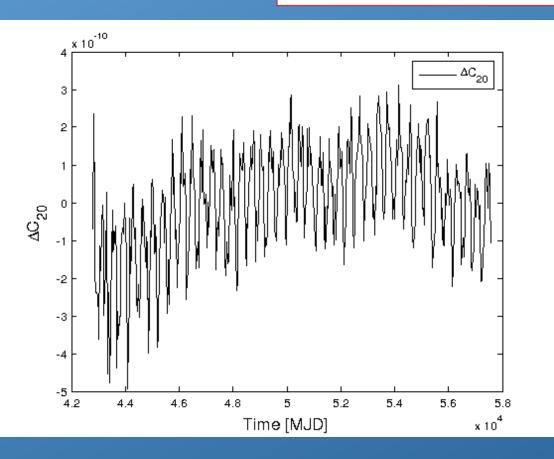
- huge improvement in the medium-wavelength and long-wavelength of the global gravity field of the Earth
 - but no significant improvements in the low degree coefficients, since the low altitude of the involved satellites (~ 500 km and ~ 260 km)
 - conversely, LAGEOS satellites are very sensitive to these low degree coefficients because of their much higher height (~ 6000 km)
 - SLR data are used to compute the low degree coefficients and their time variation (also including data from the two LAGEOS)

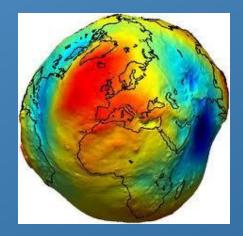


The Earth's quadrupole coefficient C20: UT/CSR monthly time series (SLR)

- SLR data from 8 satellites from 1976 up to 2016.5 (corrected to the tide-free system)
- $\Delta C_{20} = C_{20} \langle C_{20} \rangle$

Cheng et al., Deceleration in Earth's oblateness. JGR, 118 (2013)

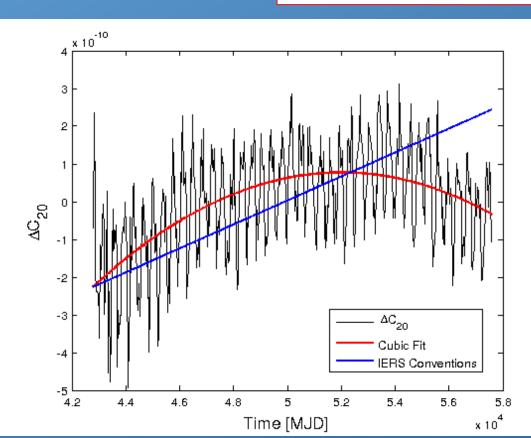




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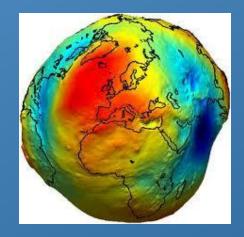


Cubic fit

$$\Delta C_{20} = a(t_0) + b(t - t_0) + c(t - t_0)^2 + d(t - t_0)^3$$

IERS Conventions
 $\Delta C_{20} = C_{20}(t_0) + \dot{C}_{20}(t - t_0)$

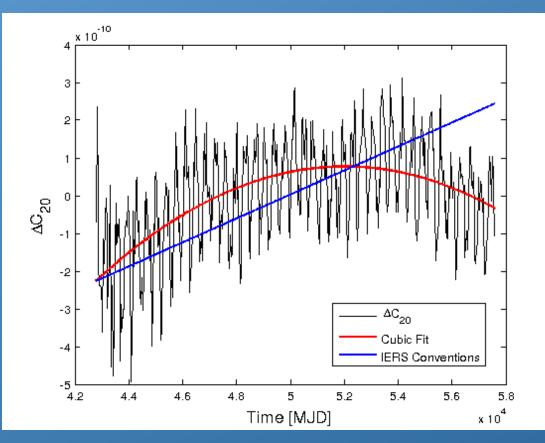
 $t_0 = J2000 = MJD 51544$

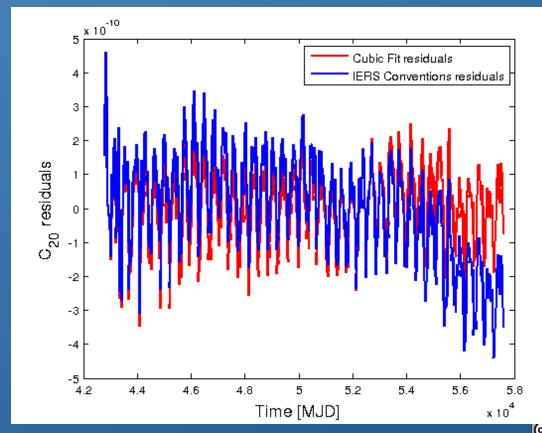


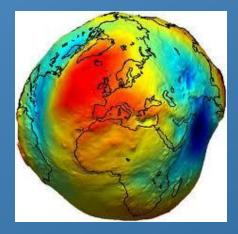
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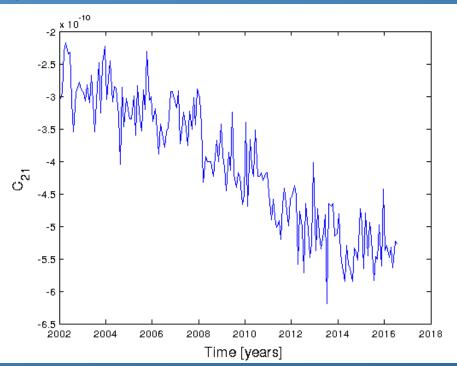


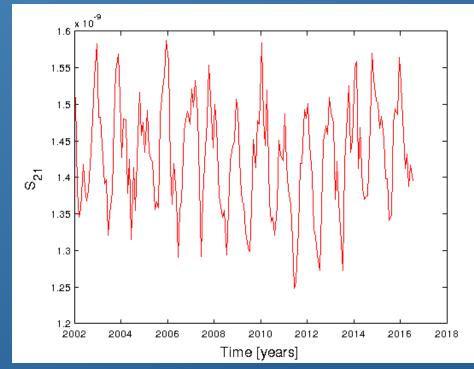
x) (

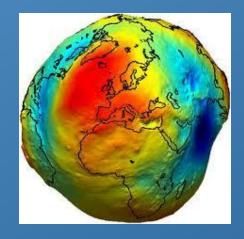
Similar considerations are valid for other low degree coefficients, as:

- C₃₀, C₅₀, ...
- C₄₀, C₆₀, ...

As well as for the coefficients that define the *figure axis* of the Earth: UT/CSR monthly time series (SLR)







Big problem with the even zonal harmonics uncertainties: systematic errors

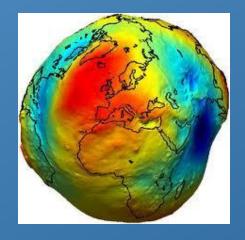
We have two main unknowns:

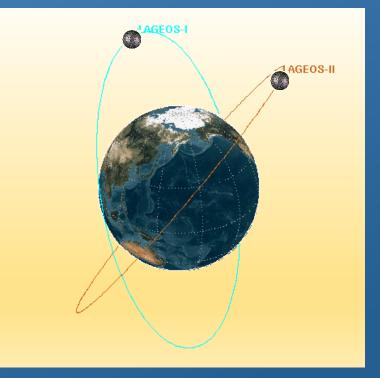
- 1. the precession on the node due to the LT effect: $\mu_{\rm LT}$;
- 2. the J_2 uncertainty: δJ_2 ;

Hence, we need two observables in such a way to eliminate the uncertainty of the first even zonal harmonic and solve for the LT effect. These observables are:

- 1. LAGEOS node:
- 2. LAGEOS II node:







 $\mu = \delta \dot{\Omega}_{I}^{res} + k \delta \dot{\Omega}_{II}^{res}$

represents the solution of a system of two equations in two unknowns

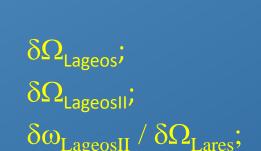
Big problem with the even zonal harmonics uncertainties: systematic errors

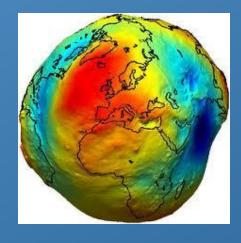
With three observables we can eliminate the uncertainty of the first and second even zonal harmonics and solve for the LT effect. These observables are:

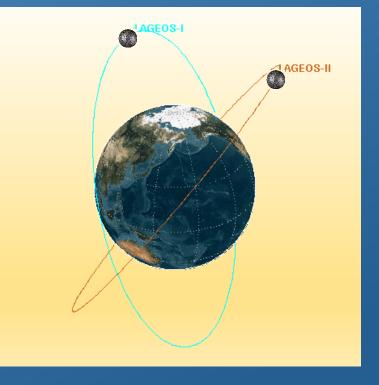
- 1. LAGEOS node:
- 2. LAGEOS II node:
- 3. LAGEOS II pericenter/LARES node

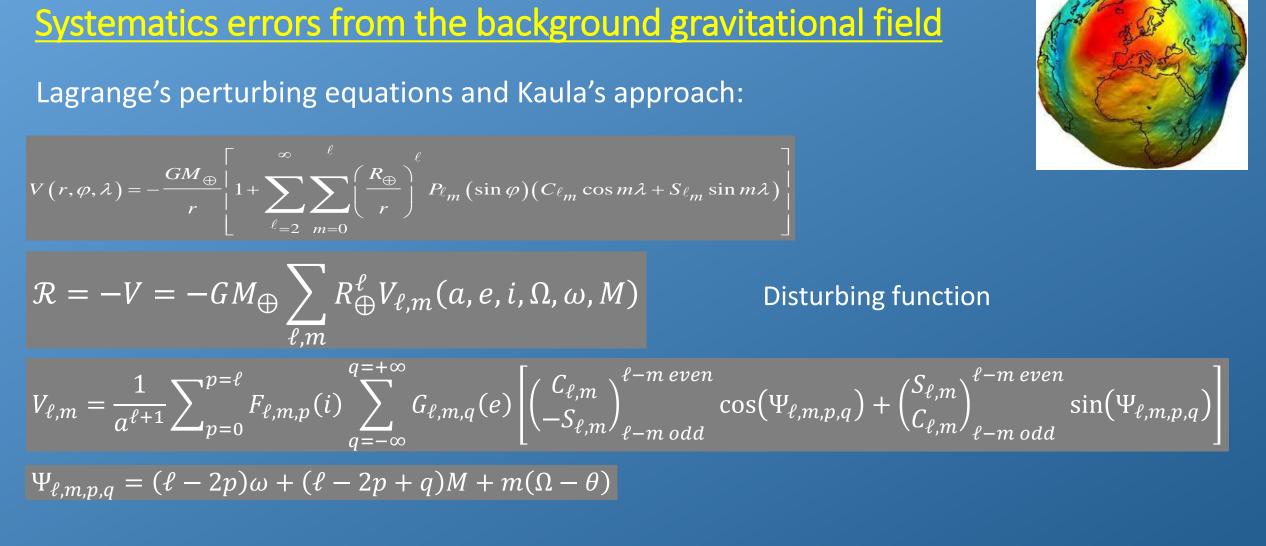
$$\begin{split} \mu &= \delta \dot{\Omega}_{I}^{res} + k \delta \dot{\Omega}_{II}^{res} \\ \mu &= \delta \dot{\Omega}_{I}^{res} + h_1 \delta \dot{\Omega}_{II}^{res} + h_2 \delta \dot{\omega}_{II}^{res} \\ \mu &= \delta \dot{\Omega}_{I}^{res} + k_1 \delta \dot{\Omega}_{II}^{res} + k_2 \delta \dot{\Omega}_{LR}^{res} \end{split}$$

Of course, including the pericenter, we have three observables: LAGEOS II perigee has been considered thanks to its larger eccentricity ($\cong 0.014$) with respect to that of LAGEOS ($\cong 0.004$) and LARES ($\cong 0.001$)









where $F_{\ell mpg}(i)$ and $G_{\ell mg}(e)$ are the inclination and eccentricity functions of Kaula

Lagrange's equations for the node and pericenter:

$$\mathcal{R} = -GM_{\bigoplus} \sum_{\ell,m} R_{\bigoplus}^{\ell} V_{\ell,m}(a, e, i, \Omega, \omega, M)$$

$$\frac{d\Omega}{dt} = -\frac{1}{na^2\sqrt{1-e^2}\sin(i)}\frac{\partial\mathcal{R}}{\partial i} \qquad \qquad \frac{d\omega}{dt} = -\frac{\sqrt{1-e^2}}{na^2e}\frac{\partial\mathcal{R}}{\partial e} + \frac{\cot(i)}{na^2\sqrt{1-e^2}}\frac{\partial\mathcal{R}}{\partial i}$$

Equations for the errors due to the even zonal harmonics (J_{ℓ} with ℓ =2,4,6,...):

$$\delta \frac{d\Omega}{dt} = \sum_{\ell=2}^{\ell=\ell_{max}} \left| \frac{\partial}{\partial J_{\ell}} \frac{d\Omega}{dt} \right|_{\ell} \delta J_{\ell}$$

a

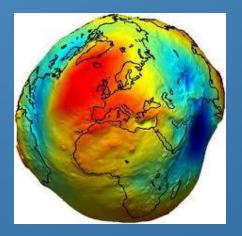
$$\delta \frac{d\omega}{dt} = \sum_{\ell=2}^{\ell=\ell_{max}} \left| \frac{\partial}{\partial J_{\ell}} \frac{d\omega}{dt} \right|_{\ell} \delta J_{\ell}$$



Big problem with the even zonal harmonics uncertainties: systematic errors

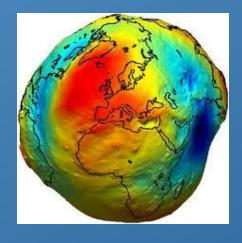
We have considered several models for the gravitational field of the Earth in our analyses from CHAMP, GRACE and GOCE missions (and from the older multi-satellite models):

□ EIGEN2S	GOCO03S
□ EIGEN-GRACE02S	GGM05S
GGM03S	□ EGM2008
GIF48	JGM-3
DGM-1S	EGM96



Big problem with the even zonal harmonics uncertainties: systematic errors

- We present the results obtained for just two of these different fields:
 - o GGM05S
 - EIGEN-GRACE02S
- i.e. for the fields used for the measurements performed so far for the precession of the orbital plane due to the Lense-Thirring effect

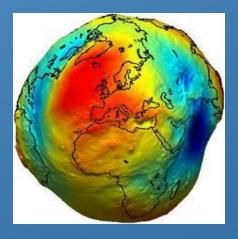


Uncertainties in the nodal rates of LAGEOS, LAGEOS II and LARES: GGM05S

 $\mu = \delta \dot{\Omega}_{I}^{res} + k_1 \delta \dot{\Omega}_{II}^{res} + k_2 \delta \dot{\Omega}_{LR}^{res}$

Calibrated errors from GGM05S

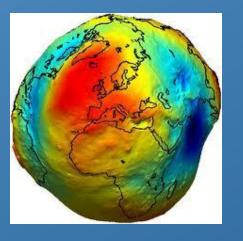
Degree	Nodal rate errors [mas/yr] – GGM05S				
l	LAGEOS	LAGEOS II	LARES		
2	-109.221282927	+201.337942030	+545.564947870		
4	-3.139224492	+1.140909687	+37.669353493		
6	-0.950852596	-3.610259385	+25.849276956		
8	-0.021732869	-0.100739170	+0.854809443		
10	+0.010691329	+0.016779287	-2.344019659		
12	+0.004927116	-0.007506301	-2.276900184		
14	+0.001096146	-0.000009023	-1.108699479		
16	+0.000133200	-0.000357766	-0.244315825		
18	-0.000007894	-0.000053184	+0.146261335		
20	-0.000012748	+0.000019595	+0.246592188		



Uncertainties in the nodal rates of LAGEOS, LAGEOS II and LARES: GGM05S

 $\mu = \delta \dot{\Omega}_{I}^{res} + k_1 \delta \dot{\Omega}_{II}^{res} + k_2 \delta \dot{\Omega}_{LR}^{res}$

Calibrated errors from GGM05S



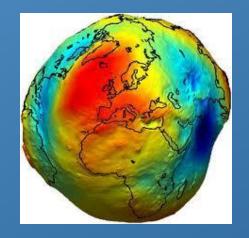
Degree	Nodal ra	Error [%]		
l	LAGEOS	LAGEOS II	LARES	δμ/μ
2	-109.221282927	+201.337942030	+545.564947870	0
4	-3.139224492	+1.140909687	+37.669353493	0
6	-0.950852596	-3.610259385	+25.849276956	0.380
8	-0.021732869	-0.100739170	+0.854809443	0.012
10	+0.010691329	+0.016779287	-2.344019659	0.308
12	+0.004927116	-0.007506301	-2.276900184	0.316
14	+0.001096146	-0.00009023	-1.108699479	0.159
16	+0.000133200	-0.000357766	-0.244315825	0.035
18	-0.000007894	-0.000053184	+0.146261335	0.021
20	-0.000012748	+0.000019595	+0.246592188	0.036



Uncertainties in the nodal rates of LAGEOS, LAGEOS II and LARES: GGM05S

 $\mu = \delta \dot{\Omega}_{I}^{res} + k_1 \delta \dot{\Omega}_{II}^{res} + k_2 \delta \dot{\Omega}_{LR}^{res}$

Calibrated errors from GGM05S



Degree	Nodal ra	Error [%]		
l	LAGEOS	LAGEOS II	LARES	δμ/μ
2	-109.221282927	+201.337942030	+545.564947870	0
4	-3.139224492	+1.140909687	+37.669353493	0
6	-0.950852596	-3.610259385	+25.849276956	0.380
8	-0.021732869	-0.100739170	+0.854809443	0.012
10	+0.010691329	+0.016779287	-2.344019659	0.308
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16	+0.000133200	-0.000357766	-0.244315825	0.035
18	-0.000007894	-0.000053184	+0.146261335	0.021
20	-0.000012748	+0.000019595	+0.246592188	0.036

Error in the combination:

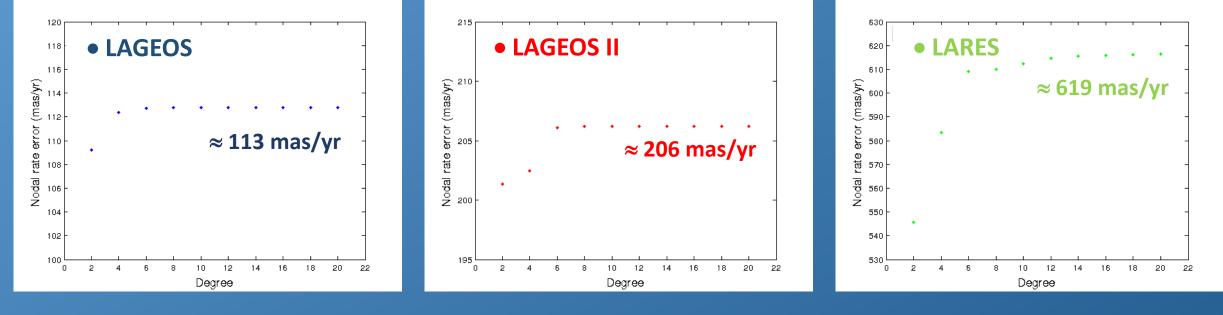
Sum Absolute Values $SAV = \sum_{\ell} \delta \mu_{\ell} \cong 1.3\% \mu$

Root Sum Squares

$$RSS = \sqrt{\sum_{\ell} \partial \mu_{\ell}^2} \cong 0.6\%\mu$$

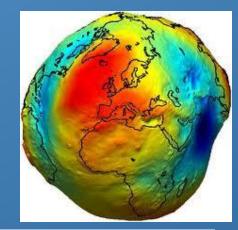


Systematics errors from the background gravitational field Uncertainties in the nodal rates of LAGEOS, LAGEOS II and LARES: GGM05S Plots from previous Table



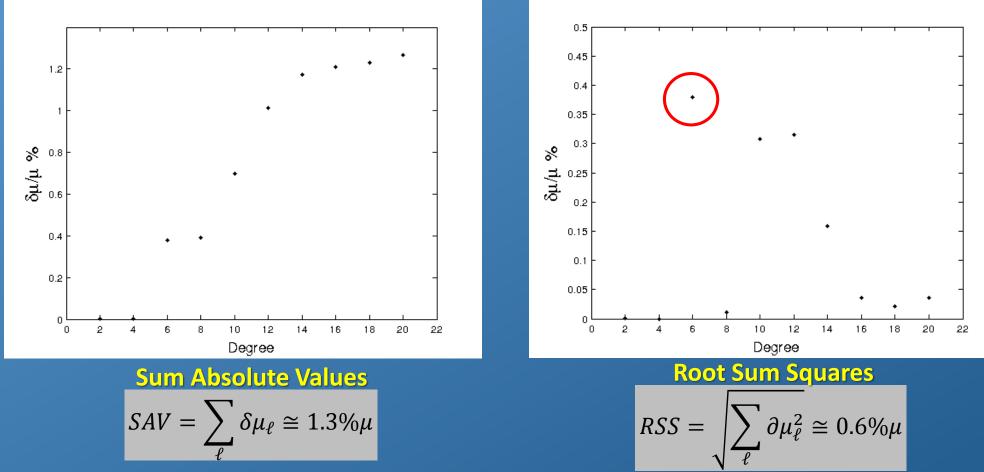
Nodal rate of the three satellites

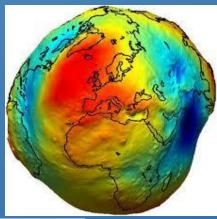




Uncertainties in the combination of the three nodal rates: **GGM05S** <u>Plots from previous Table</u>

Systematics errors from the background gravitational field

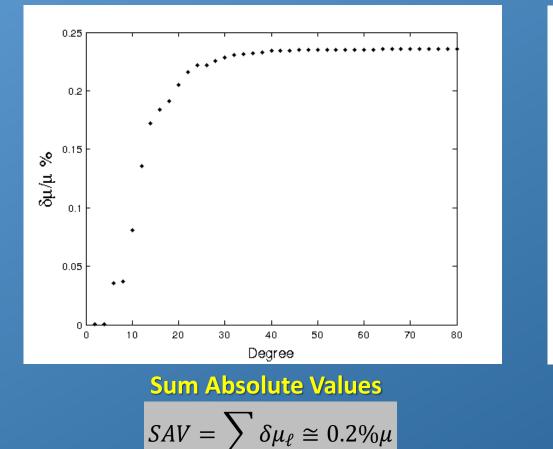


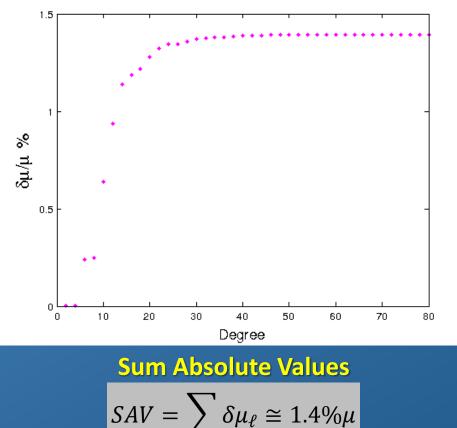


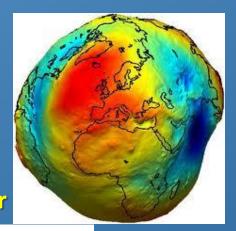
Uncertainties in the combination of the three nodal rates: **EIGEN-GRACE02S**

Cumulative Formal Error

Cumulative Calibrated Error

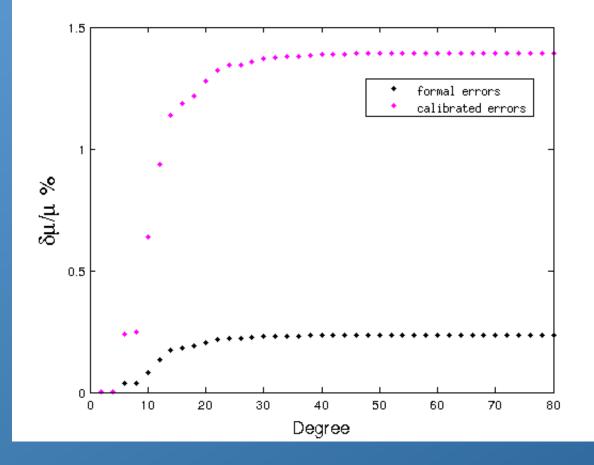


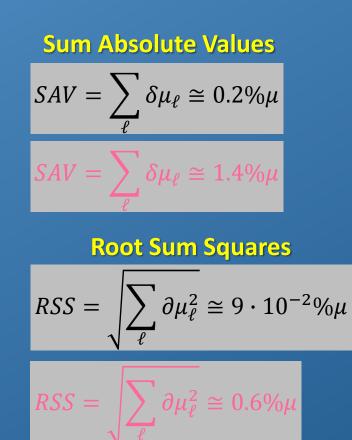




Systematics errors from the background gravitational field

Uncertainties in the combination of the three nodal rates: **EIGEN-GRACE02S**

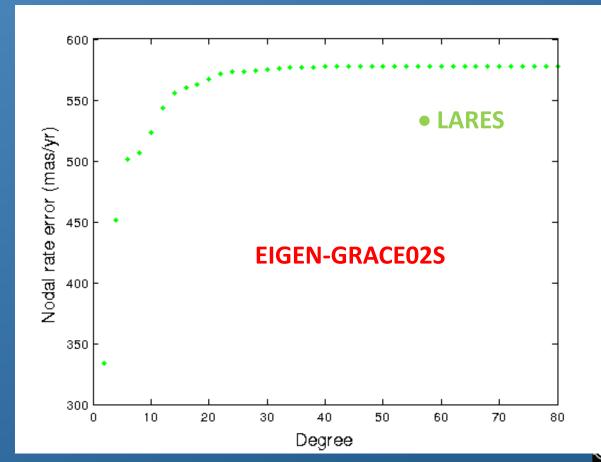


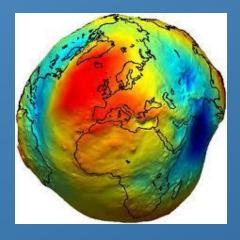


Systematics errors from the background gravitational field

Preliminary (partial) conclusion after the comparison of different models for the Earth's gravitational field:

- in order to estimate the error budget for the Lense-Thirring effect measurement due to the uncertainties of the even zonal harmonics coefficients it is enough to consider the contribution of the harmonics up to the degree & ≈ 20 or 30
- the contribution from the higher degrees is negligible, also in the case of LARES

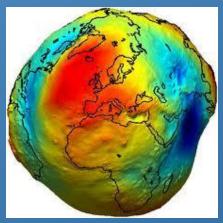




Systematics errors from the background gravitational field

Preliminary (partial) conclusion after the comparison of different models for the Earth's gravitational field:

- we remark that, even if, in the case of LARES, the mis-modelled precession does not decrease at high degrees, we dot not report any runaway inflation of the uncertainty, as instead claimed by some authors in the literature
- possible sources of errors in the evaluation of nodal precession contribution may originate in the evaluation of the derivatives of eccentricity and inclination functions of the Kaula expansion of the potential
- to overcome computational issues of high-degree contributions, we have avoided to simplify partial sums of trigonometric terms in the form of rational expressions (with the risk of huge errors when performing derivatives)
- rather, we left expanded the trigonometric polynomials which, even if quite long, are well within the range of operation of any algebraic manipulator, resulting in accurate determination of high-degree terms



Summary

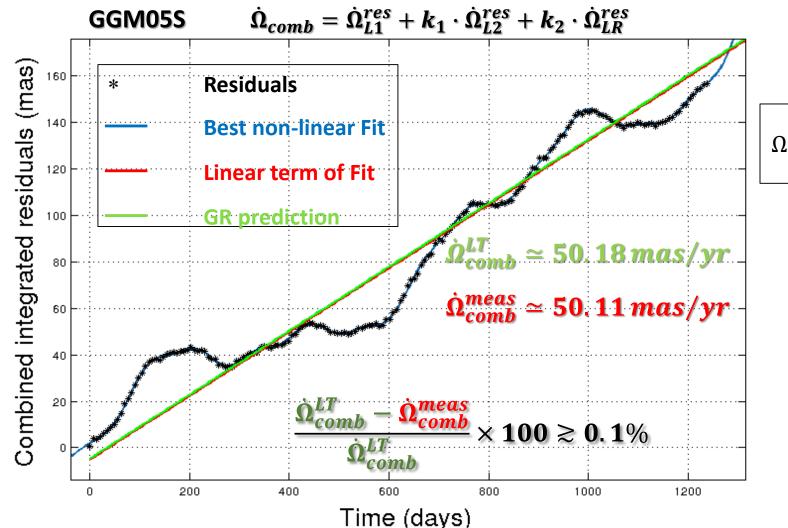
- The LARASE experiment and its goals
- Relativistic effects to be measured
- Systematic errors from the background gravitational field
- Some recent measurements of relativistic effects
- Conclusions and future work



LARASE measurements of relativistic precessions:

- 1. A new preliminary measurement of the Lense-Thirring precession with the two LAGEOS satellites and LARES (2016)
- 2. Measurement of the overall GR precession of LAGEOS II pericenter (2014)
- 3. Constraints on alternative theories of gravitation (2014)

A very preliminary new measurement of the Lense-Thirring effect with the two LAGEOS and LARES (3.4 yr)



We fitted also for a minimum of three to a maximum of twelve tidal waves (both solid and ocean):

$$\Omega^{Fit} = a + b \cdot t + \sum_{i=1}^{n} A_i \cdot \sin\left(\frac{2\pi}{P_i} \cdot t + \Phi_i\right)$$

Indeed, tides mismodelling plus unmodelled nongravitational forces due to thermal effects may corrupt the measurement of the relativistic effect.

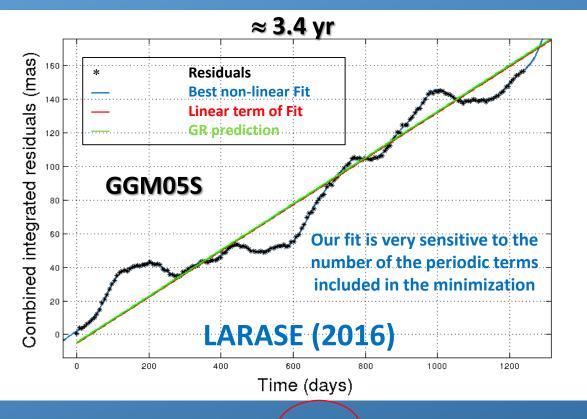
For instance, the (both solid and ocean) K1 tides have the same periods of the right ascension of the node of the satellites:

≈1044 days, ≈569 days and ≈224 days

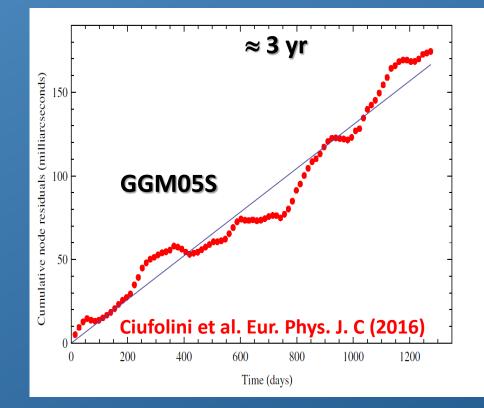
$\dot{\Omega}_{comb} = \dot{\Omega}_{L1}^{res} + k_1 \cdot \dot{\Omega}_{L2}^{res} + k_2 \cdot \dot{\Omega}_{LR}^{res}$

Some recent measurements of relativistic effects

Comparison with an independent measurement



 $\mu = (0.999 \pm 0.001) \pm \varepsilon(sys)$



$\mu = (0.994 \pm 0.002) \pm 0.05$

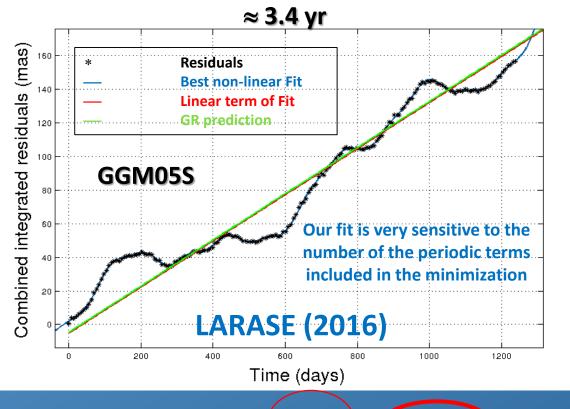
0.2% formal error of the fit (1-sigma) plus 5% preliminary estimate of systematics (4% grav. + 1% non-grav.)

Up to a few % from a sensitivity analysis of the main tidal waves

$\dot{\Omega}_{comb} = \dot{\Omega}_{L1}^{res} + k_1 \cdot \dot{\Omega}_{L2}^{res} + k_2 \cdot \dot{\Omega}_{LR}^{res}$

Some recent measurements of relativistic effects

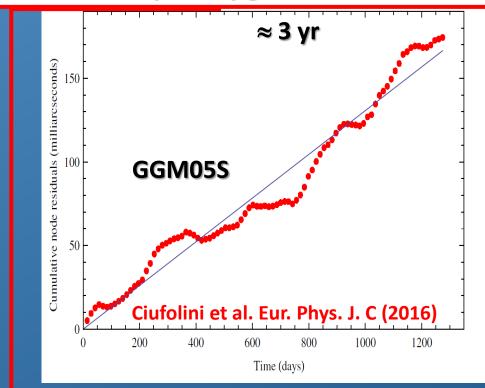
Comparison with an independent measurement



 $\mu = (0.999 \pm 0.001) \pm \varepsilon(sys)$

Up to a few % from a sensitivity analysis of the main tidal waves

Indeed, a robust and reliable estimate of systematics is one of the primary goals of LARASE !



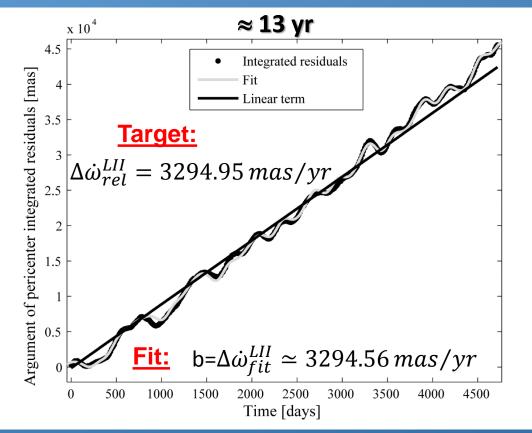
$\mu = (0.994 \pm 0.002) \pm 0.05$

0.2% formal error of the fit (1-sigma) plus 5% preliminary estimate of systematics (4% grav. + 1% non-grav.)

 \sim

A precise and accurate measurement performed in the recent pass (2014):

Fit to the pericenter residuals of LAGEOS II



Lucchesi, Peron, Phy. Rev. D, 89, 2014

Fitting function for the pericenter:

$$\Delta \omega^{FIT} = a + b \cdot t + c \left(t - t_0\right)^2 + \sum_{i=1}^n D_i \sin\left(\frac{2 \cdot \pi}{P_i} \cdot t + \Phi_i\right)$$

- We obtained b ≅ 3294.6 mas/yr, very close to the prediction of GR
- The discrepancy is just **0.01%**
- From a sensitivity analysis, with constraints on some of the parameters that enter into the least squares fit, we obtained an upper bound of 0.2%

 $\varepsilon = 1 - (0.12 \pm 2.10) \cdot 10^{-3} \pm 2.5 \cdot 10^{-2}$

Lucchesi, Peron, Phy. Rev. D, 89, 2014

Summary of the constraints in gravitational theories so far obtained

TABLE XVIII. Summary of the results obtained in the present work; together with the measurement error budget, the constraints on fundamental physics are listed and compared with the literature.

Parameter	Values and uncertainties (this study)	Uncertainties (literature)	Remarks
$\epsilon_{\omega} - 1$	$-1.2 \times 10^{-4} \pm 2.10 \times 10^{-3} \pm 2.54 \times 10^{-2}$		Error budget of the perigee precession measurement in the field of the Earth
$\frac{ 2+2\gamma-\beta }{3}-1$	$-1.2 \times 10^{-4} \pm 2.10 \times 10^{-3} \pm 2.54 \times 10^{-2}$	$\pm (1.0 \times 10^{-3}) \pm (2 \times 10^{-2})^{a}$	Constraint on the combination of PPN parameters
$ \alpha $	$\lesssim 0.5 \pm 8.0 \pm 101 \times 10^{-12}$	$\pm 1 imes 10^{-8b}$	Constraint on a possible (Yukawa-like) NLRI
$\begin{array}{c} \mathcal{C}_{\oplus \text{LAGEOSII}} \\ 2t_2 + t_3 \end{array}$	$ \leq (0.003 \text{ km})^4 \pm (0.036 \text{ km})^4 \pm (0.092 \text{ km})^4 \\ \lesssim 3.5 \times 10^{-4} \pm 6.2 \times 10^{-3} \pm 7.49 \times 10^{-2} $	$\pm (0.16 \text{ km})^{4^{\circ}}; \pm (0.087 \text{ km})^{4^{\circ}}$ $3 \times 10^{-3^{\circ}}$	Constraint on a possible NSGT Constraint on torsion

^aFrom the preliminary estimate of the systematic errors of [166] for the perihelion precession of Mercury.

^bFrom [167] with Lunar-LAGEOS *GM* measurements.

^cFrom [5] and based on a partial estimate for the systematic errors.

^dFrom [7] and based on the analysis of the systematic errors only.

^eFrom [168] with no estimate for the systematic errors.

Lucchesi, Peron, Phy. Rev. D, 89, 2014

Summary of the constraints in gravitational theories so far obtained

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$ \alpha $	$\lesssim \! 0.5\pm 8.0\pm 101 \times 10^{-12}$	$\pm 1 \times 10^{-8b}$	Constraint on a possible (Yukawa-like) NLRI
$\frac{C_{\oplus \text{LAGEOSII}}}{ 2t_2 + t_3 }$	$ \leq (0.003 \text{ km})^4 \pm (0.036 \text{ km})^4 \pm (0.092 \text{ km})^4 \\ \lesssim 3.5 \times 10^{-4} \pm 6.2 \times 10^{-3} \pm 7.49 \times 10^{-2} $	$\pm (0.16 \text{ km})^{4c}; \pm (0.087 \text{ km})^{4d}$ 3×10^{-3e}	Constraint on a possible NSGT Constraint on torsion

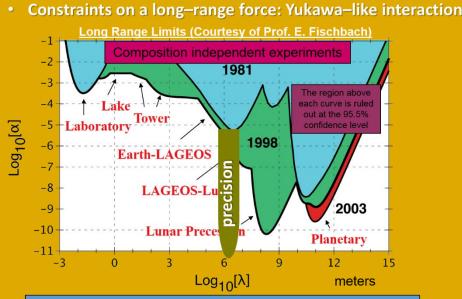
^aFrom the preliminary estimate of the systematic errors of [166] for the perihelion precession of Mercury.

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^dFrom [7] and based on the analysis of the systematic errors only.

^eFrom [168] with no estimate for the systematic errors.



Reference: Coy, Fischbach, Hellings, Standish, & Talmadge (2003)

Conclusions and future work



The LARASE (LAser RAnged Satellites Experiment) activities, in terms of orbit modelling improvements and relativistic measurements, are ongoing. In particular:

- the preliminary measurements of relativistic effects are very promising and very precise (thanks to both models and POD improvements)
- a deep evaluation of the main systematic error sources has been started
- we presented the results for the even zonal harmonics uncertainties (calibrated and not) of various models of the Earth's gravitational field up to a very high degree for its expansion in spherical harmonics
- in order to have a reliable evaluation for the error it is enough to consider harmonics up to degree ≈30, at least in the case of the Lense-Thirring effect measurement
- of course, other additional analysis are also necessary to finally reach a very robust error budget



Conclusions and future work



Future work:

- a new study has been started in order to improve the thermal models of the two LAGEOS and to develop a thermal model for LARES by means of a Finite Element Model (FEM)
- thermal effects are the main candidate in order to explain a residual along-track deceleration on LARES of the order of 2×10^{-13} m/s² (not explained by the neutral drag), but charged drag will be also investigated
- the POD set up (stations position/velocity and biases, International Conventions/Reference frames, etc.) is in line with that of the Analysis Centers of the ILRS and is continuously improving:
 - ✓ POD is very good for the two LAGEOS
 - ✓ but some improvement is still expected for LARES
- new measurements of relativistic effects in the field of the Earth are foreseen



More on LARASE activities

@ Poster Session

Poster X3.120 EGU2017-13124



LARASE website

http://larase.roma2.infn.it

Classical and Quantum Gravity

Class. Quantum Grav. 32 (2015) 155012 (50pp)

doi:10.1088/0264-9381/32/15/155012

Testing the gravitational interaction in the field of the Earth via satellite laser ranging and the Laser Ranged Satellites Experiment (LARASE)

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