

# ***Peculiarities of quasi-adiabatic dynamics of charged particles in current sheets with a magnetic shear***

**Helmi Malova<sup>1,2</sup>, Victor Popov<sup>2,3</sup>, and Elena Grigorenko<sup>2</sup>**

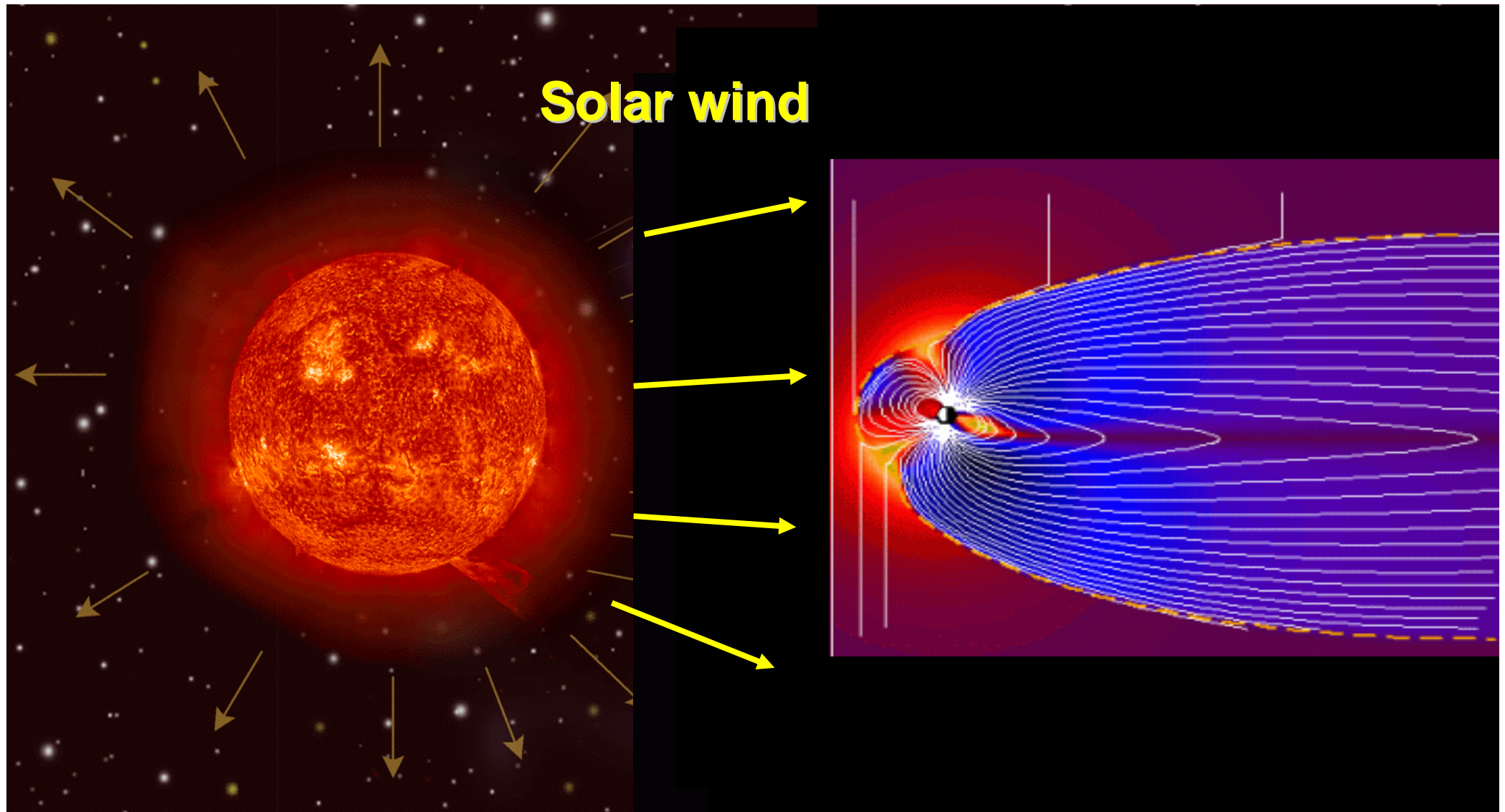
<sup>1</sup>Moscow State University, Scobeltsyn Institute of Nuclear Physics, Space Sciences, Moscow, Russian Federation  
(hmalova@yandex.ru)

<sup>2</sup>Space Research Institute, Russian Academy of Sciences, Moscow, Russia

<sup>3</sup>Faculty of Physics, Lomonosov Moscow State University, Moscow, Russia

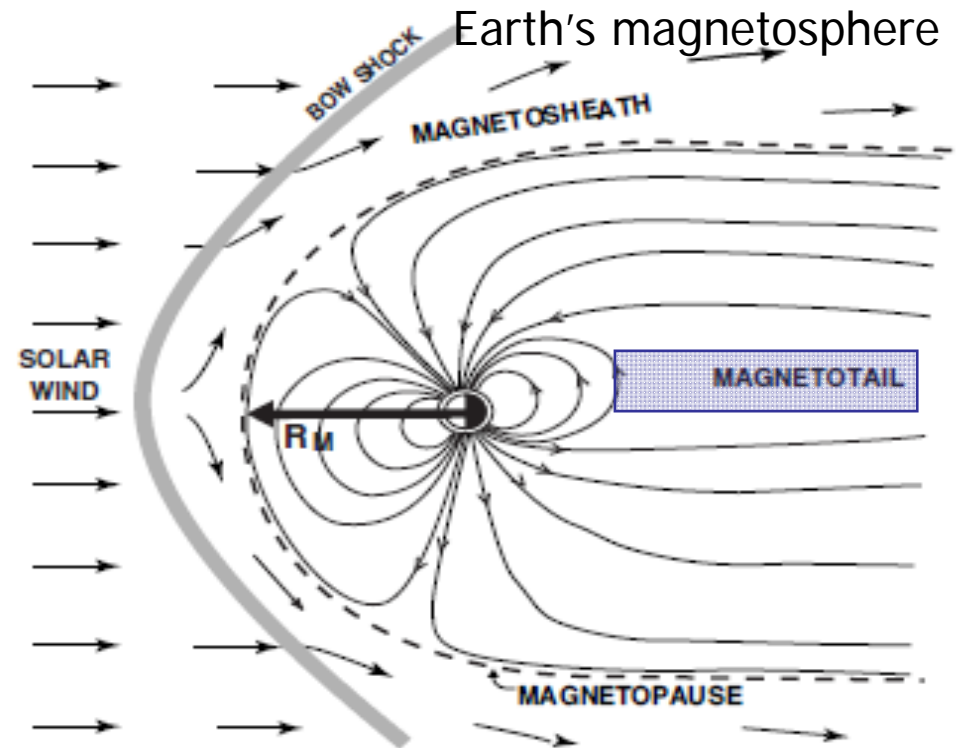
with the participation of Marina Belyalova (IKI RAS)

*Dynamical interaction of solar wind and planetary magnetospheres leads to the formation of thin current sheets in the magnetotails with thicknesses about one to several Larmor radii  $\rho_L$*

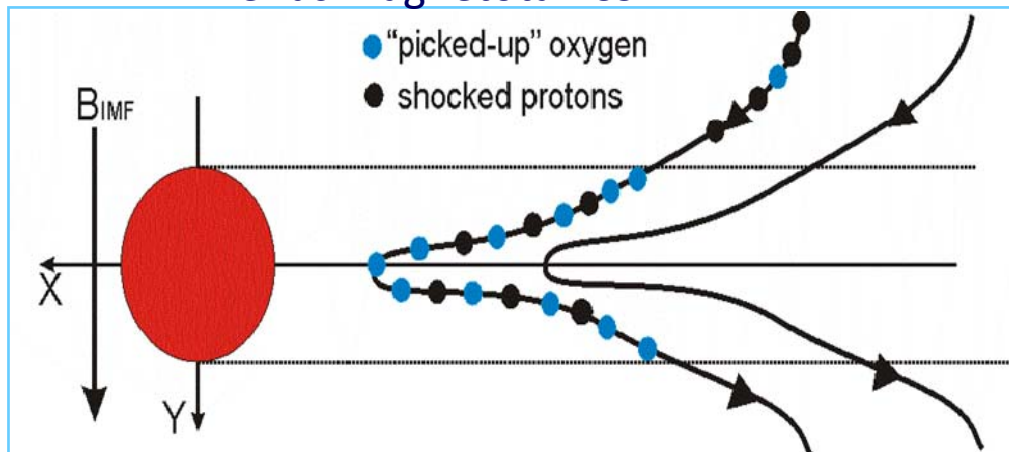


# Magnetotail current sheets

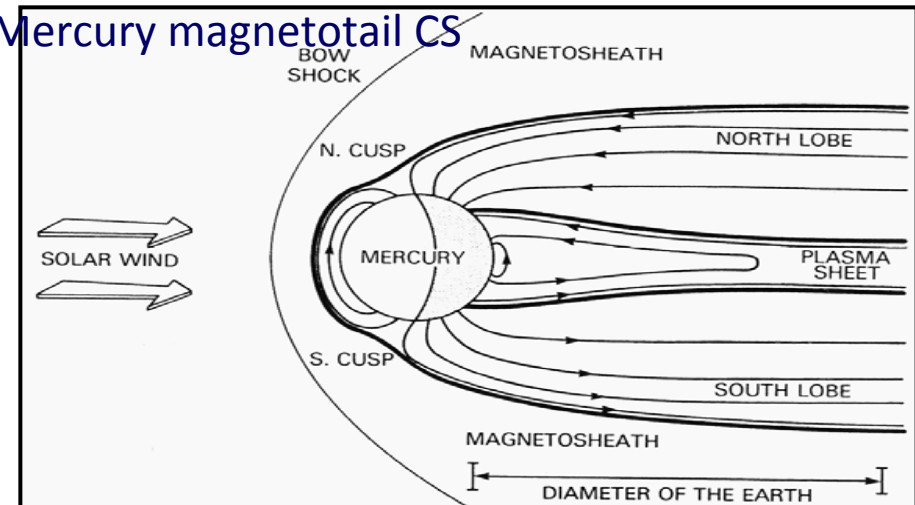
It is shown that not only the Earth, but Mercury, Venus and other planets can have magnetospheres with thin current sheets in their magnetotails



Venus magnetotail CS



Mercury magnetotail CS

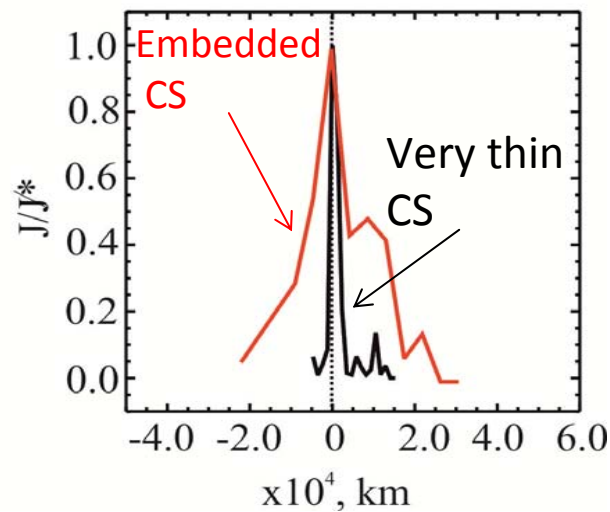
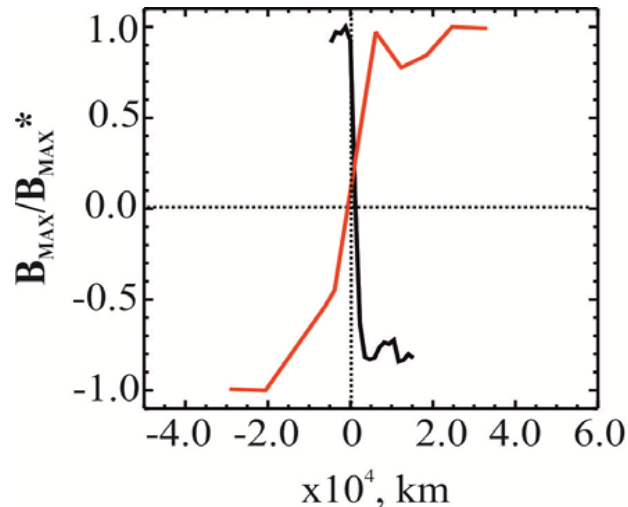


# The similar features are observed also in current sheets in:

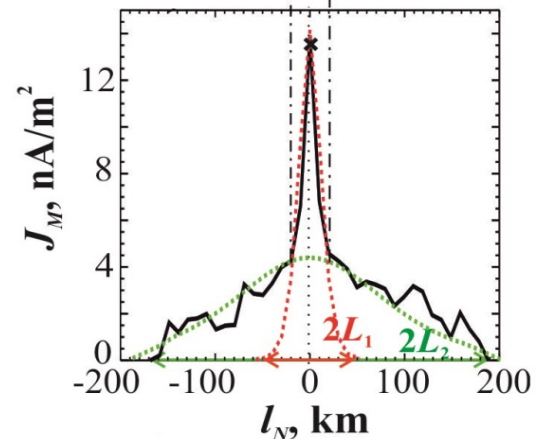
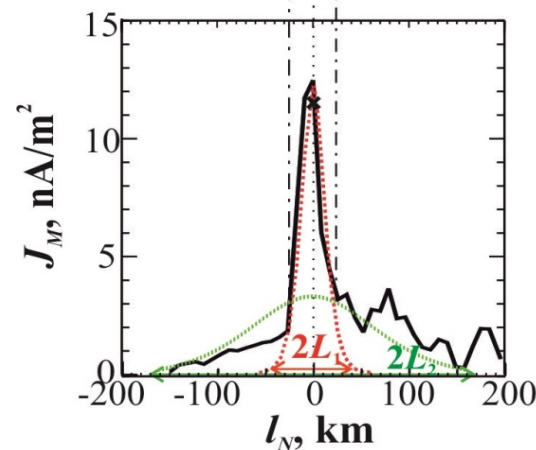
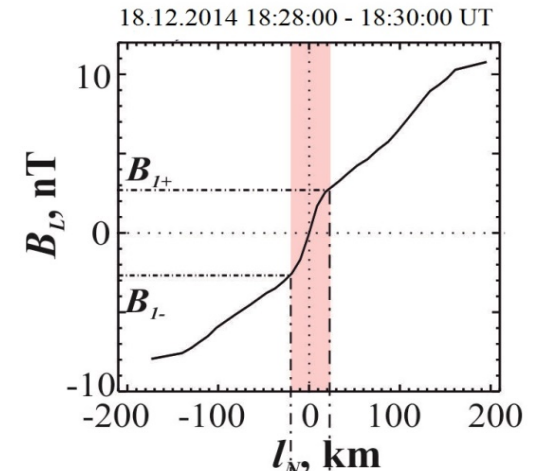
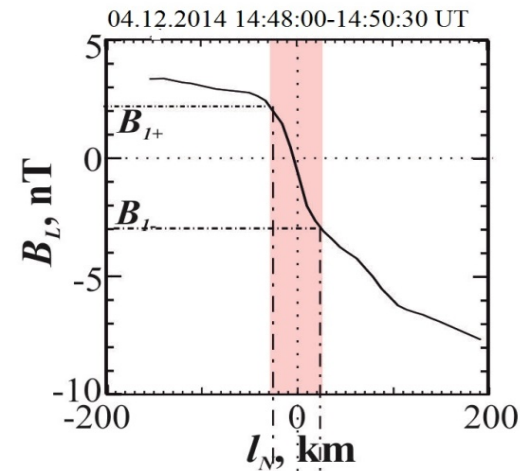
## the Solar Wind

## In the Martian magnetotail

STEREO-A observations of two CS crossings



MAVEN observations of two CS crossings



A thin embedded layer has a thickness  $L_1$  less than gyroradius of thermal protons

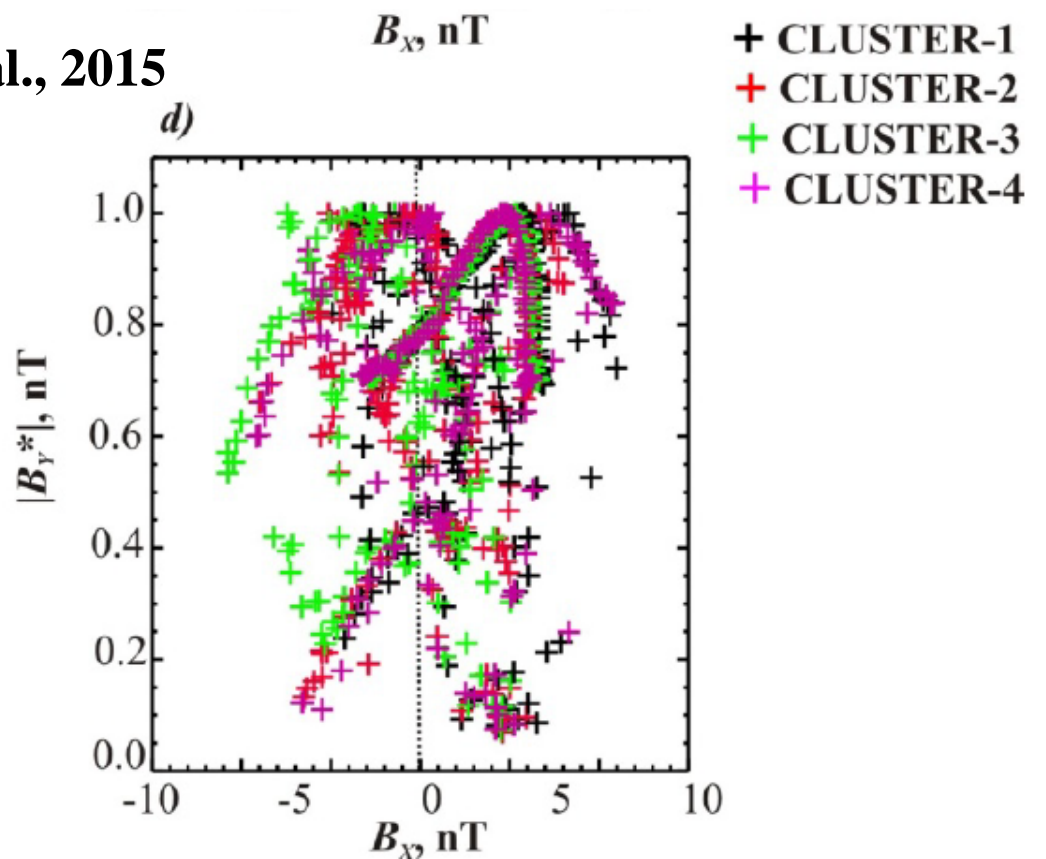
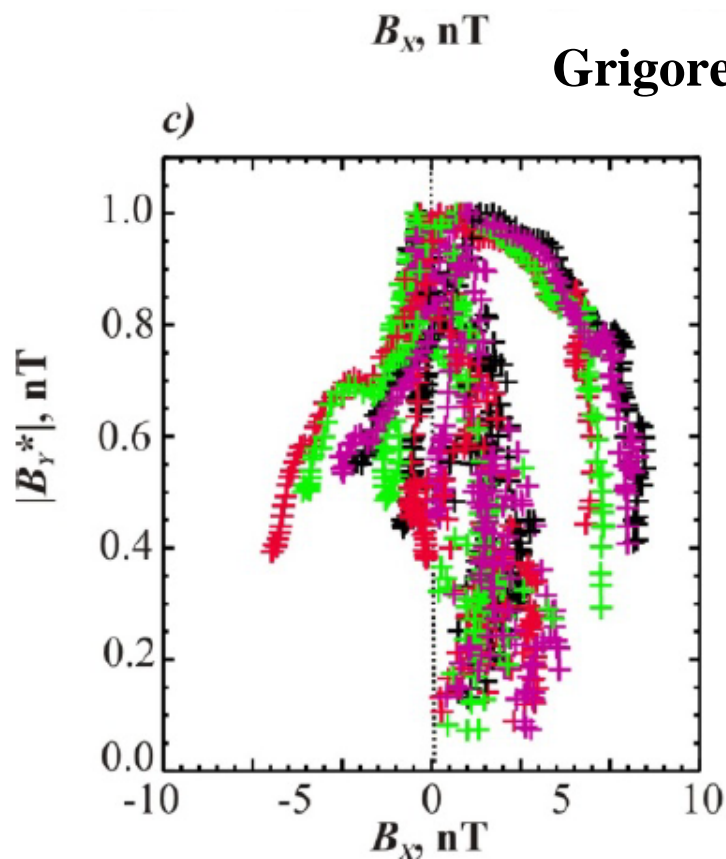
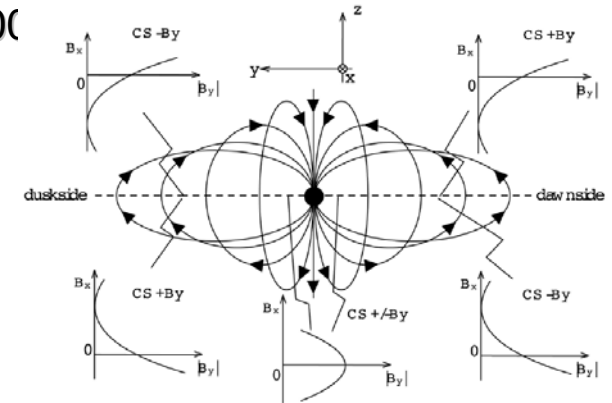
(Malova et al., ApJ, 2017)

(Grigorenko et al., JGR, 2017)



# After subtracting the flaring effect, a bell-shaped magnetic field $B_y$ is demonstrated in TCS (Rong, 2006)

Sometimes observations of  $B_y$  component in magnetotail current sheet do not correlate with the global magnetic shear in the solar wind (Petrukovich, 2011; Rong et al., 2012).



## Questions

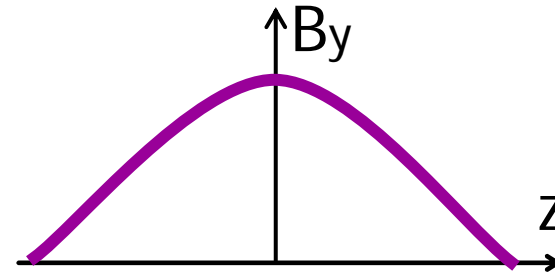
- What is the particle dynamics in thin current sheet with different shapes of magnetic shear components?
- How particle dynamics influence the structure and evolution of thin current sheet?

The aim of this work is to investigate a quasi-adiabatic particle dynamics in thin current sheet with a guiding field and to understand its influence to the current sheet structure

# *Three possible shapes of the shear magnetic component in thin current sheet*

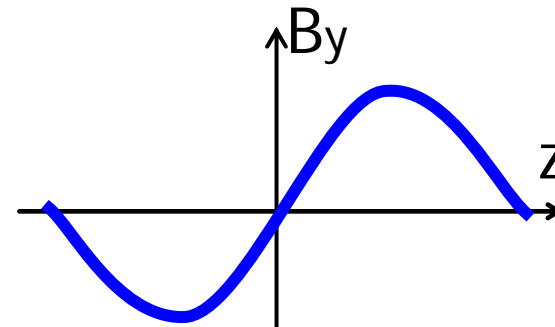
## 1) Symmetric $B_y(z)$

(planetary magnetotails)



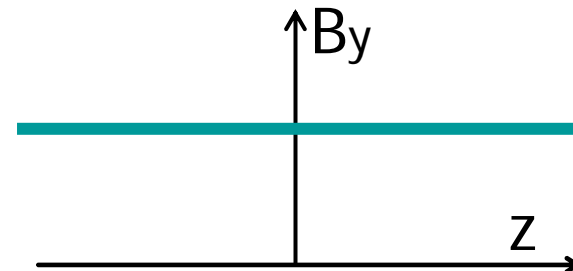
## 2) Antisymmetric $B_y(z)$

(Heliospheric and strong current sheets, planetary magnetotails)



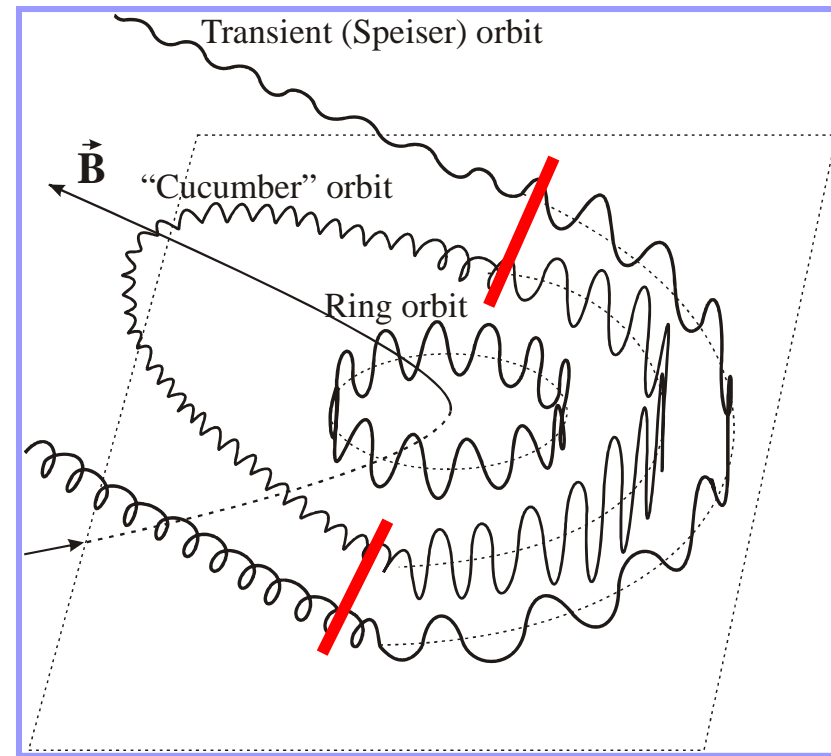
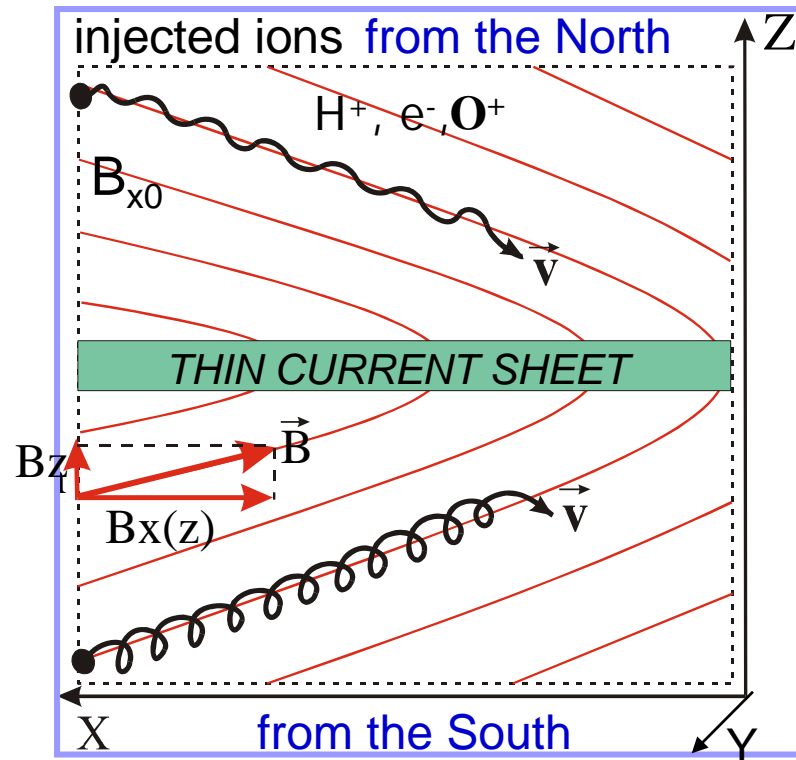
## 3) $B_y = \text{const}$

(planetary magnetotails, magnetopauses, Current structures in the solar wind)



Remark: In the case of a different sign of the shear component  $B_y(z)$  all results presented below have mirror symmetry relatively the neutral plane  $Z=0$

## Quasi-adiabatic (QA) ion dynamics in the field $\mathbf{B} = \{B_{x0}(z/L), B_y(z), B_z\}$



$$m \frac{d^2 \vec{\mathbf{R}}}{dT^2} = \frac{e}{c} [\vec{\mathbf{V}} \times \vec{\mathbf{B}}] + e \vec{\mathbf{E}}$$

$$I_z \equiv \frac{1}{2\pi} \oint m v_z dz \approx \text{const}, E = \text{const}, P_y = \text{const}$$

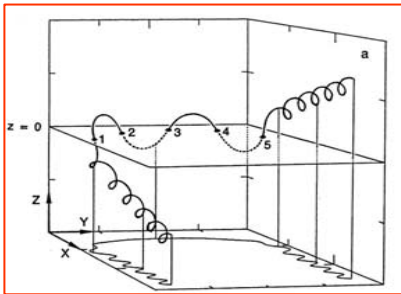
The quasi-adiabatic integral  $I_z$  is approximately conserved during ion motion. Its jumps during separatrix crossing are much smaller than the value of the QA integral itself

$$\Delta I_z \ll I_z$$



# Poincaré map: partitioning of a phase space for quasi-adiabatic ion dynamics in thin current sheets

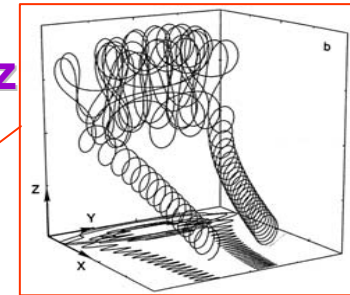
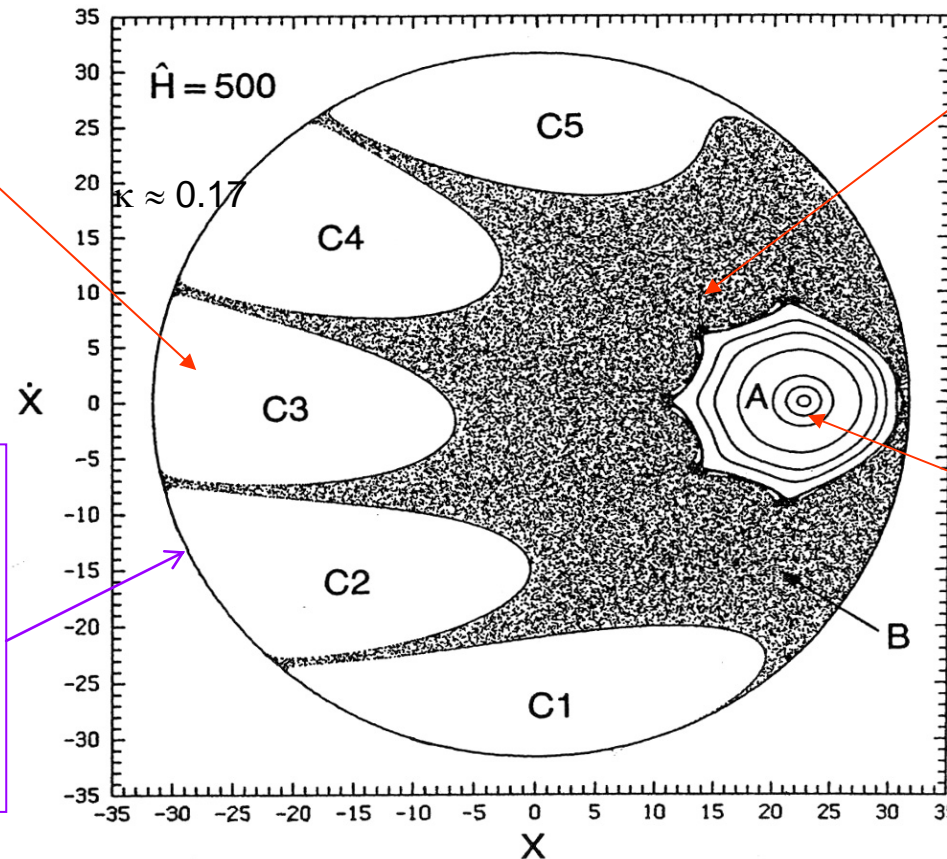
$$I_z \equiv \frac{1}{2\pi} \oint m v_z dz \approx \text{const}$$



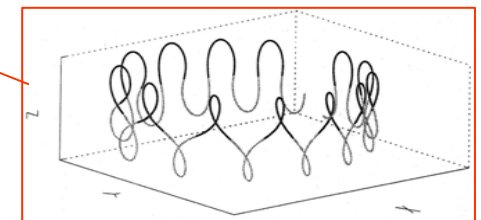
**transient particle**  
("Speiser")

**Surface of Section Plot**  
(Poincaré map  
is the set of  
normalized crossing  
points of CS plane  
 $z=0$ )

**Approximate quasi-adiabatic integral of motion  $I_z$**



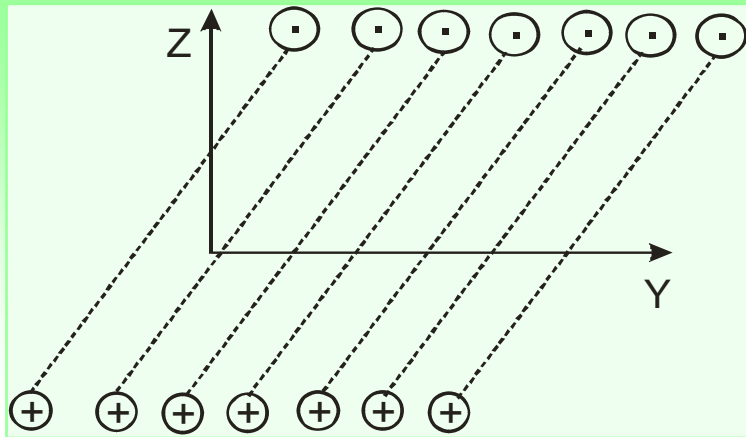
**quasi-trapped trajectory**  
("cucumber")



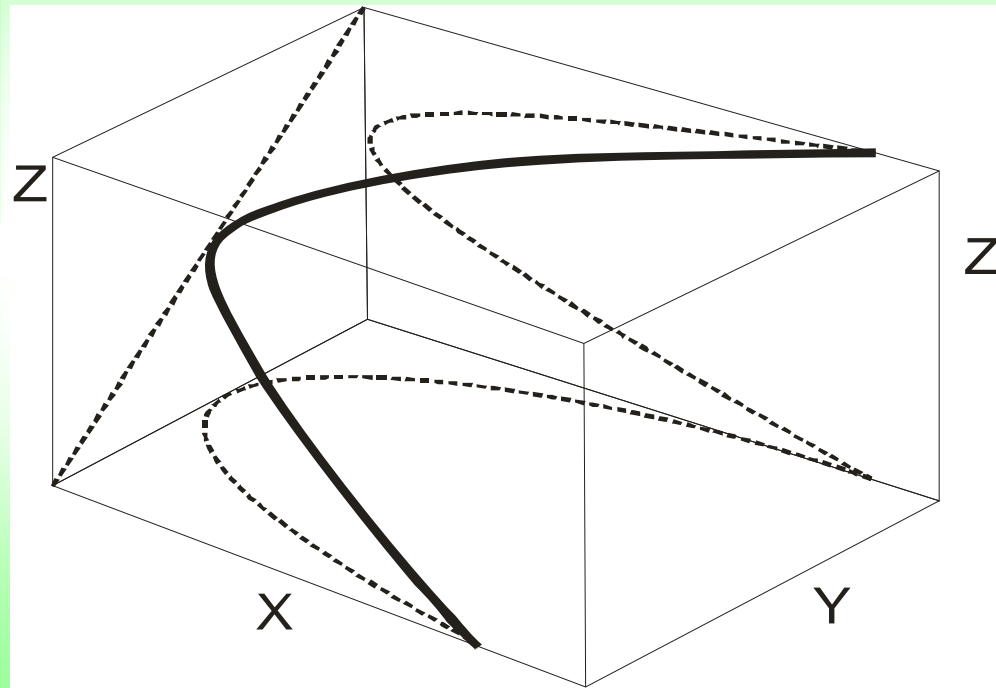
**Trapped ion trajectory**

Chen, 1992  
Buchner and Zelenyi,  
JGR, 1989

## Magnetic field lines configuration at the constant $B_y$



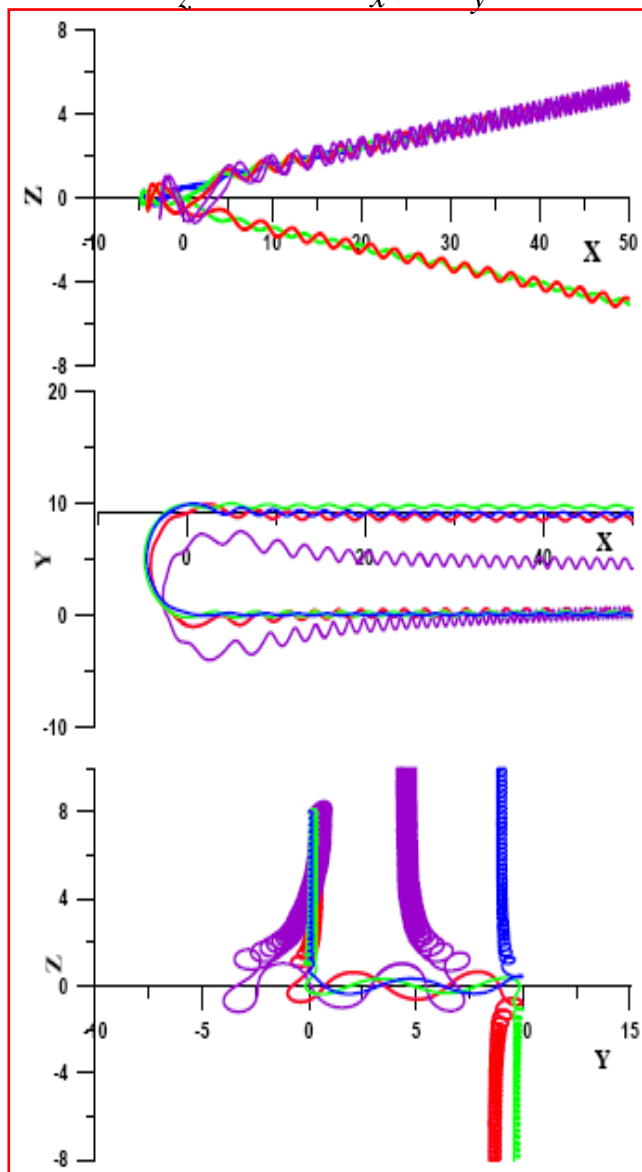
$$B_y = \text{const}$$



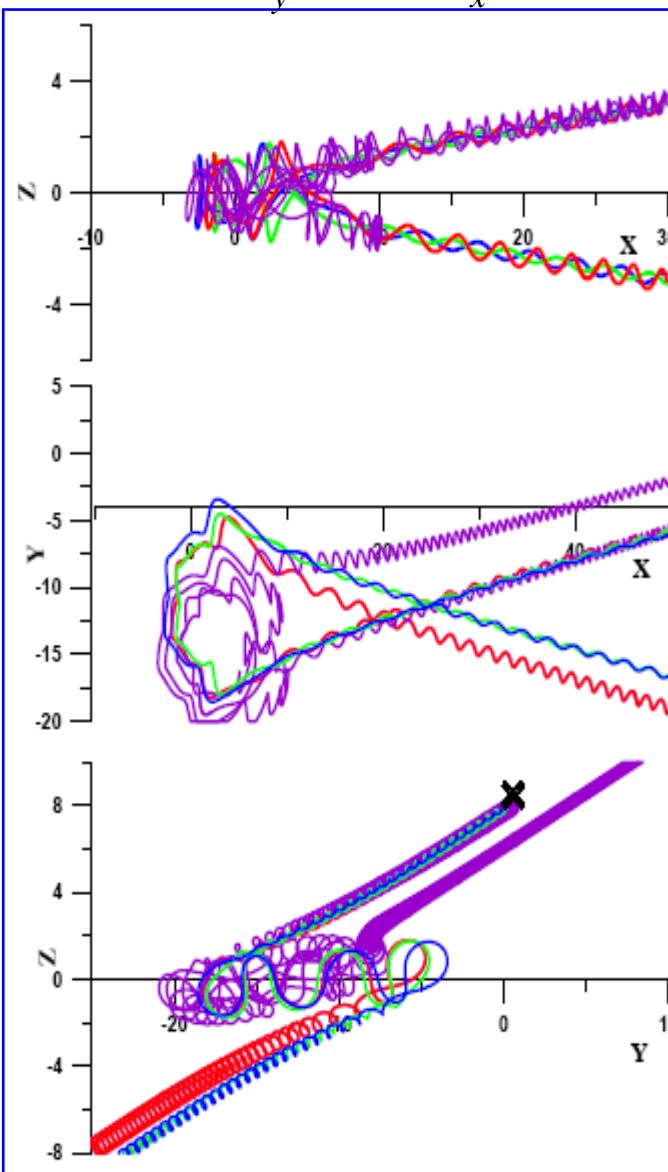
During numerical experiment protons were traced from the Northern to the Southern (marked as N  $\rightarrow$  S) hemisphere toward the current sheet and from the Southern to the Northern hemisphere (S  $\rightarrow$  N)

# Trajectories of 4 particles launched from the North hemisphere (N->S) at different $B_y$ values

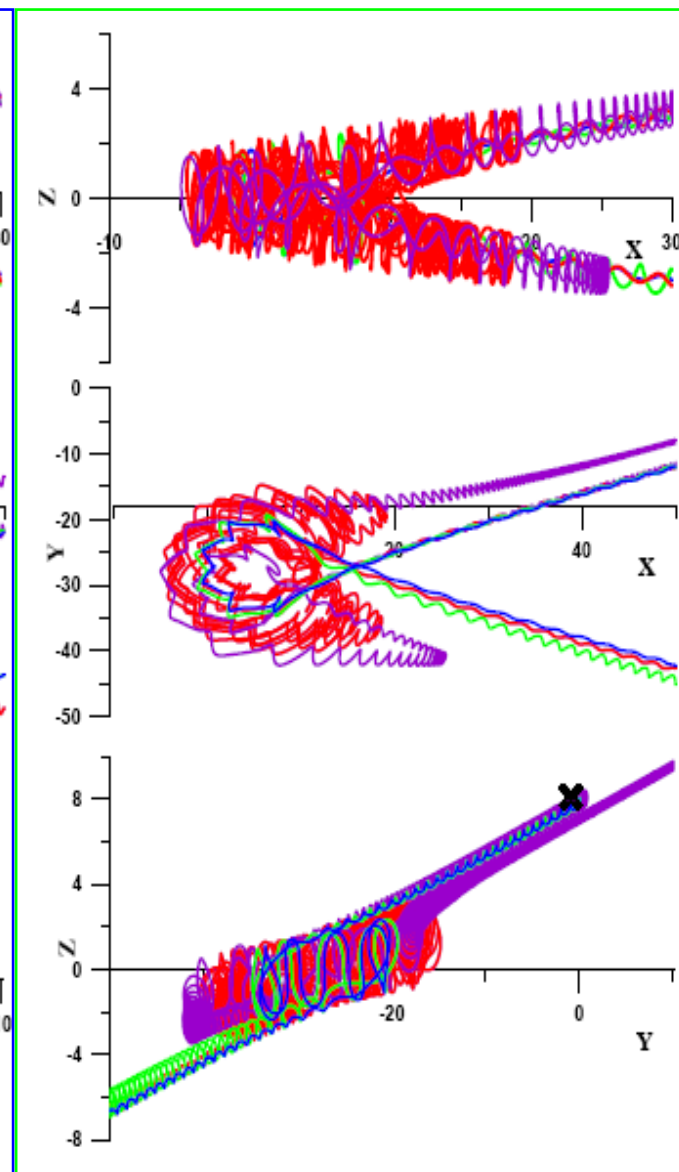
$$B_z = 0.1 B_x, B_y = 0$$



$$B_y = 0.2 B_x$$

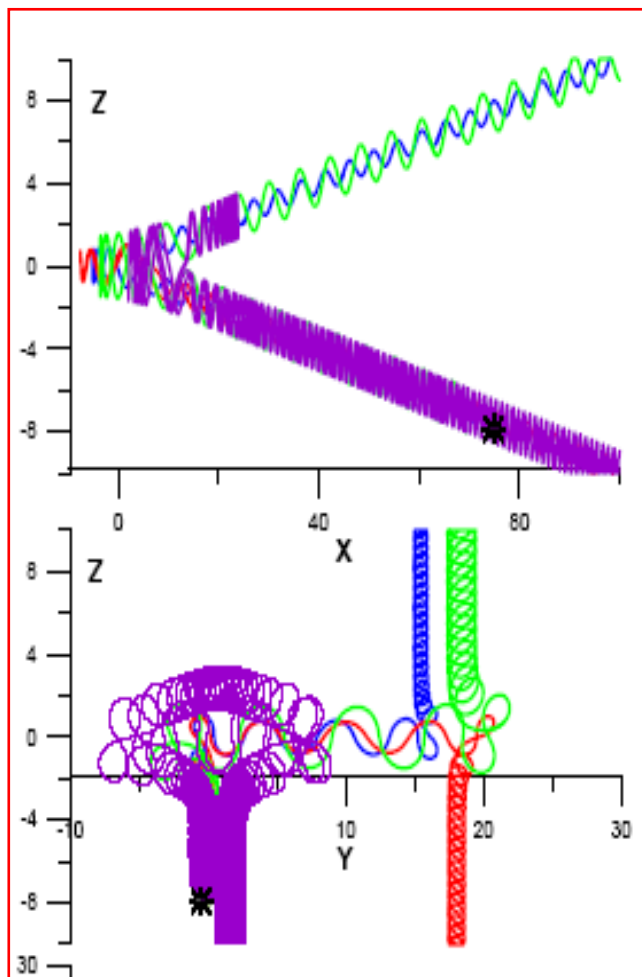


$$B_y = 0.4 B_x$$

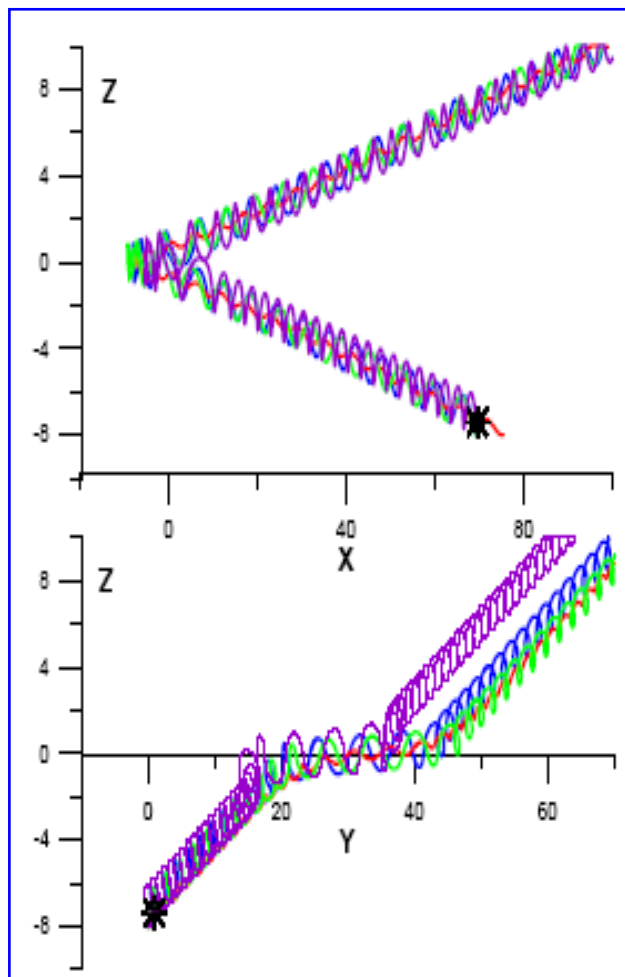


# Scattering of 4 particles launched from the Southern hemisphere (S->N) at different $B_y$ values

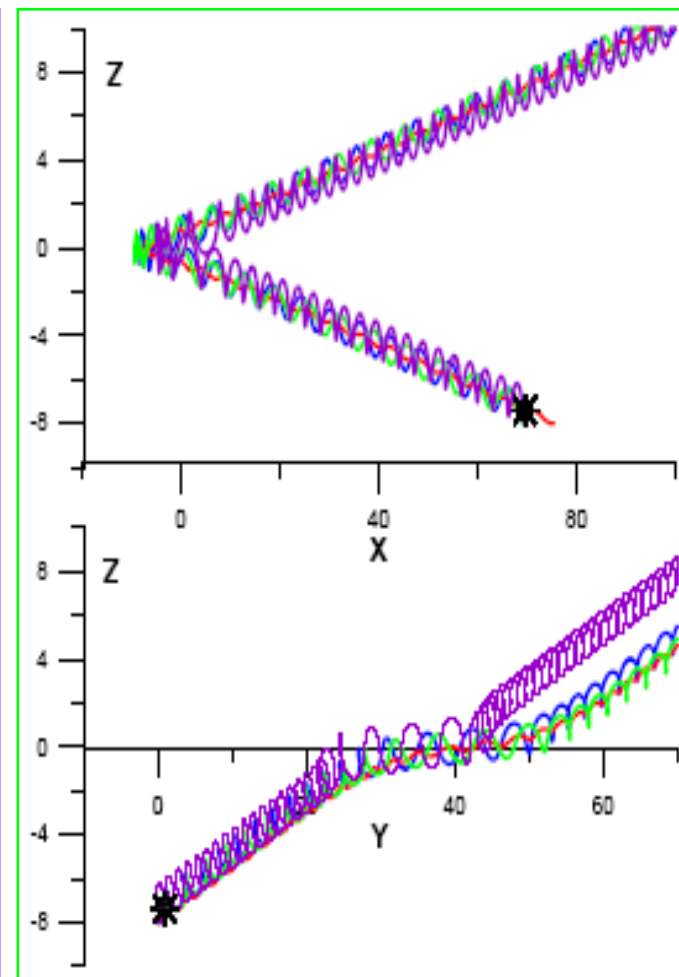
$$B_z = 0.1 B_x, \quad B_y = 0$$



$$B_y = 0.2 B_x$$

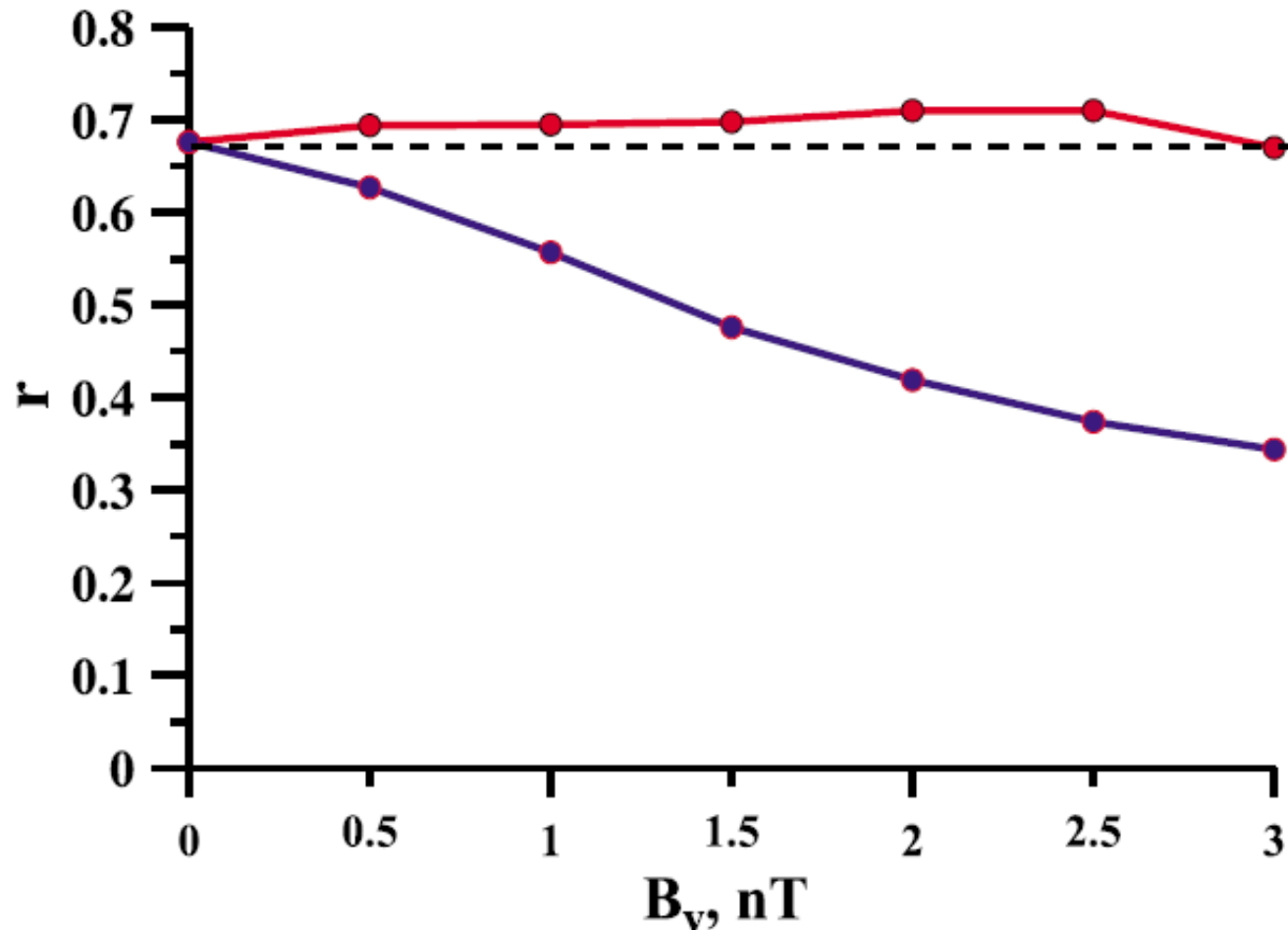


$$B_y = 0.4 B_x$$



# Asymmetry of particle scattering at $B_y$ -const

Coefficient of plasma reflection  $r=n_{\text{ref}}/n_0$  as a function of the value of magnetic component  $B_y$

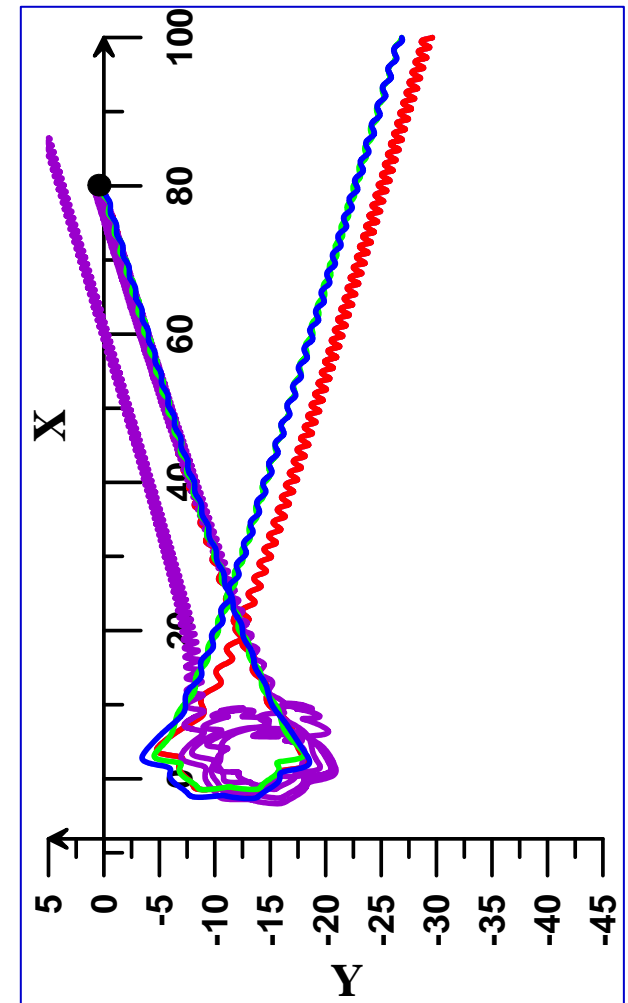
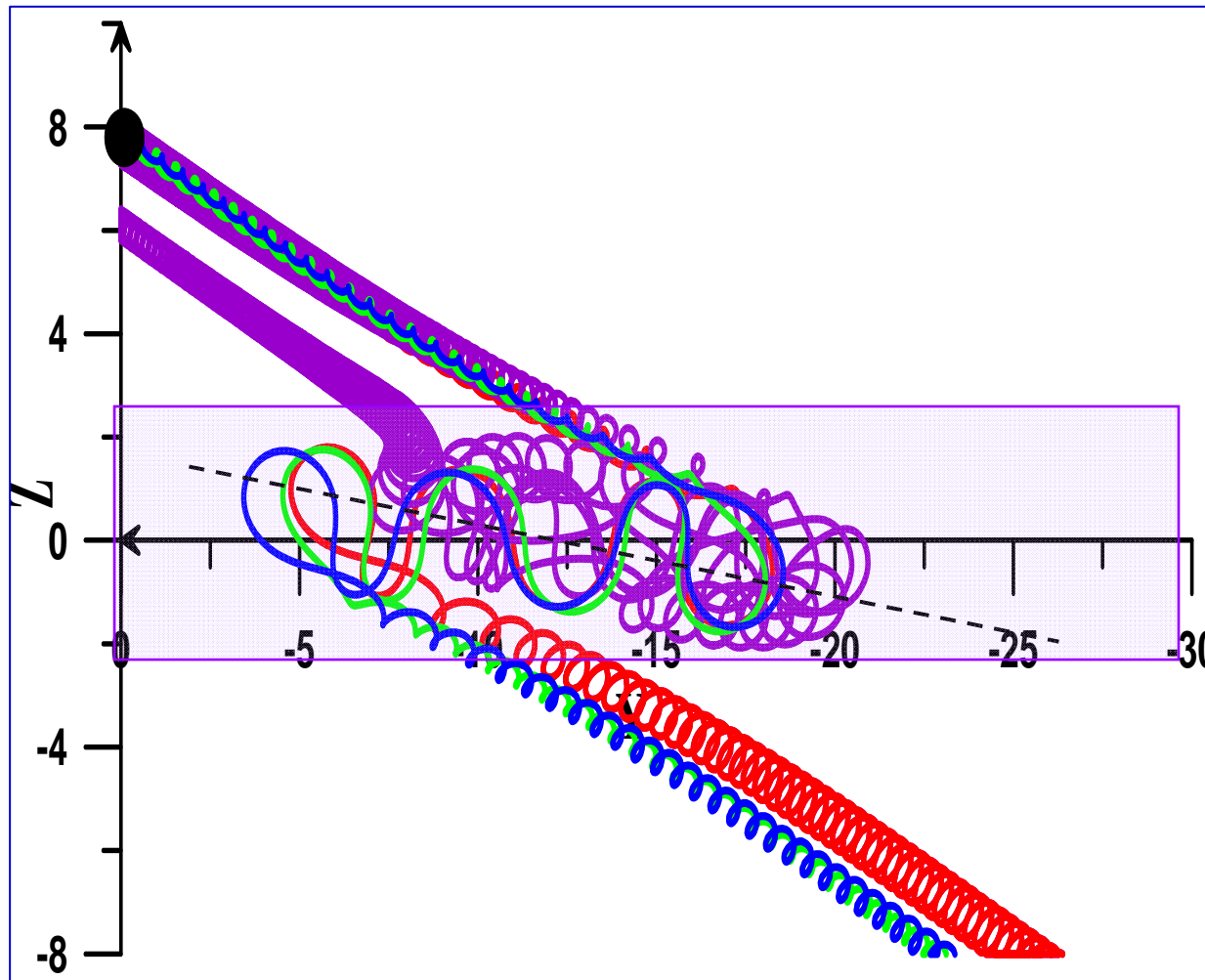


Coefficient  $r=n_{\text{ref}}/n_0$  is the relative density of ions  $n_{\text{ref}}$  reflected from the neutral plane relatively the total density  $n_0$

- Particles were launched from the Northern hemisphere (N->S)
- Particles were launched from the Southern hemisphere (S->N)

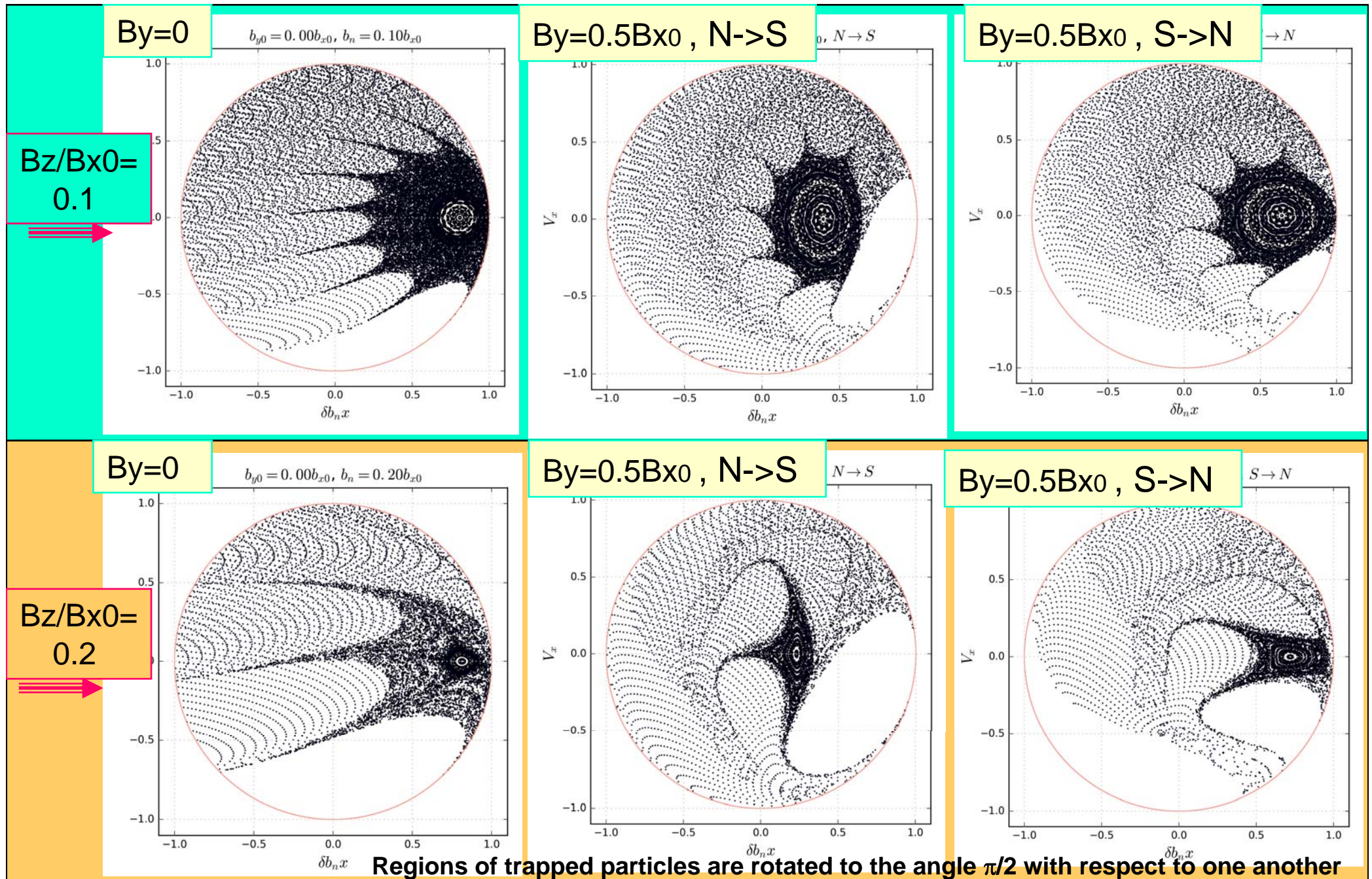


Particle dynamics at  $B_y/B_{x0}=\text{const}$ : Current sheet is thickened due to the geometric effect (tilting of the meandering plane)



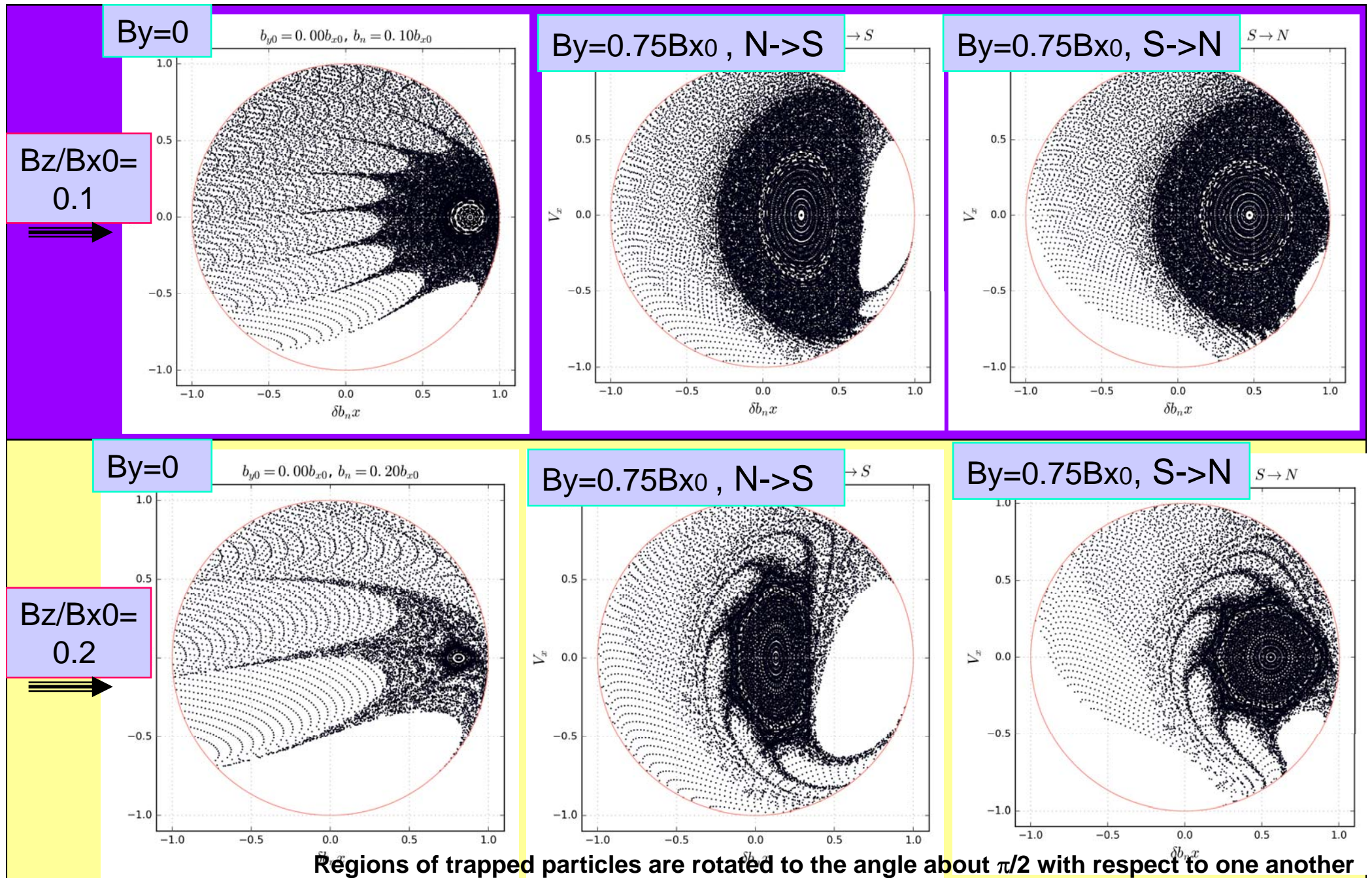
$B_{x0}$  is the tangential magnetic field value at the edges of current sheet

# Poincaré sections of quasi-adiabatic particle motion ( $B_y = \text{const}$ )



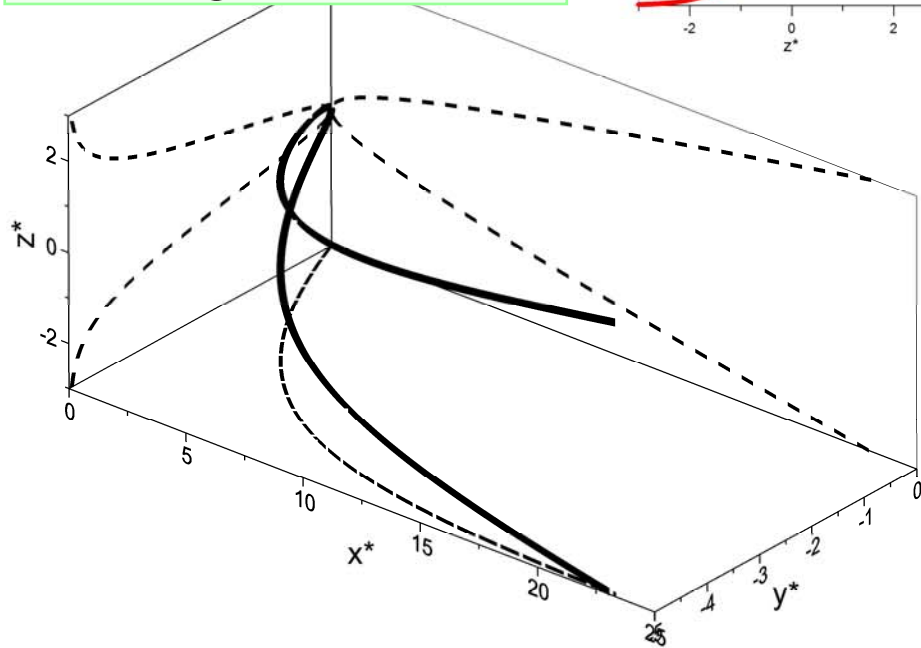
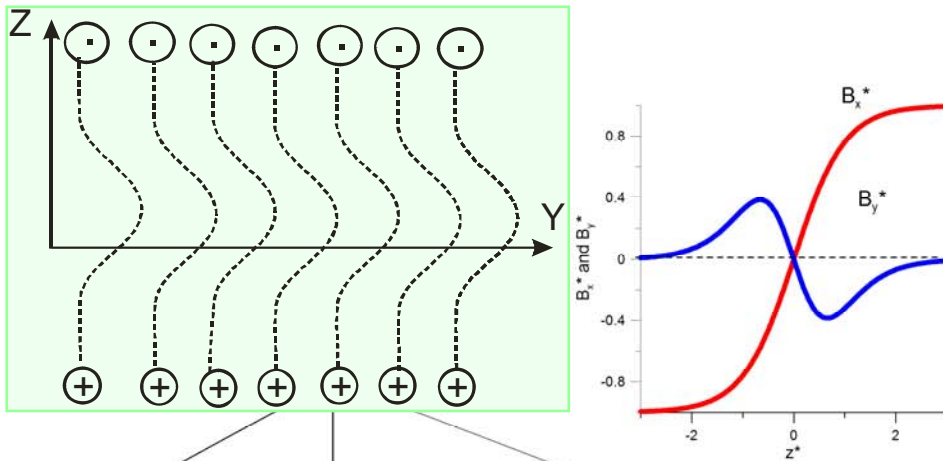


# Poincaré sections of quasi-adiabatic particle motion ( $B_y = \text{const}$ )

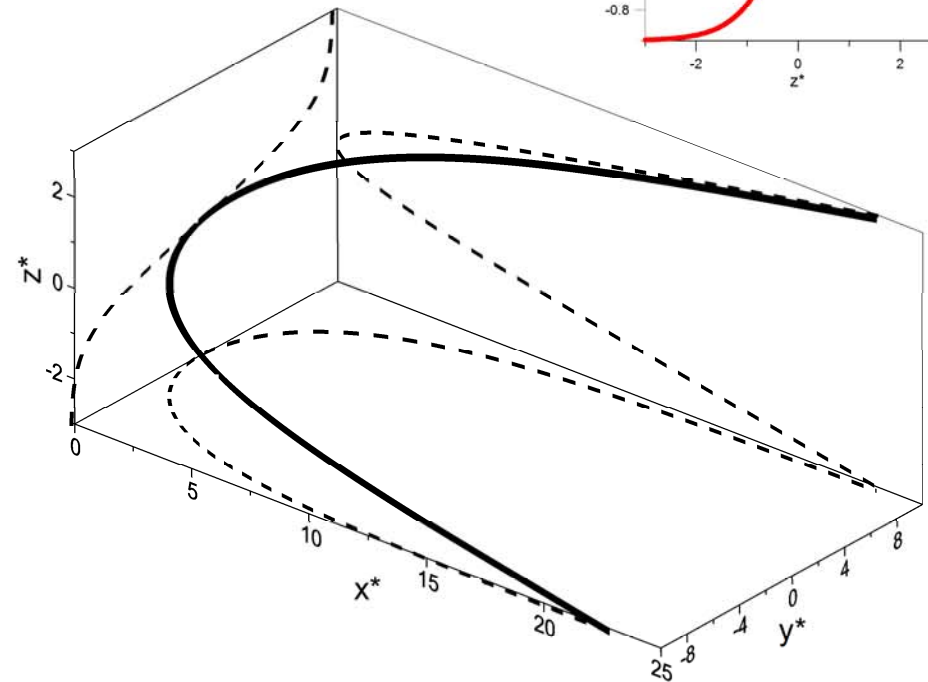
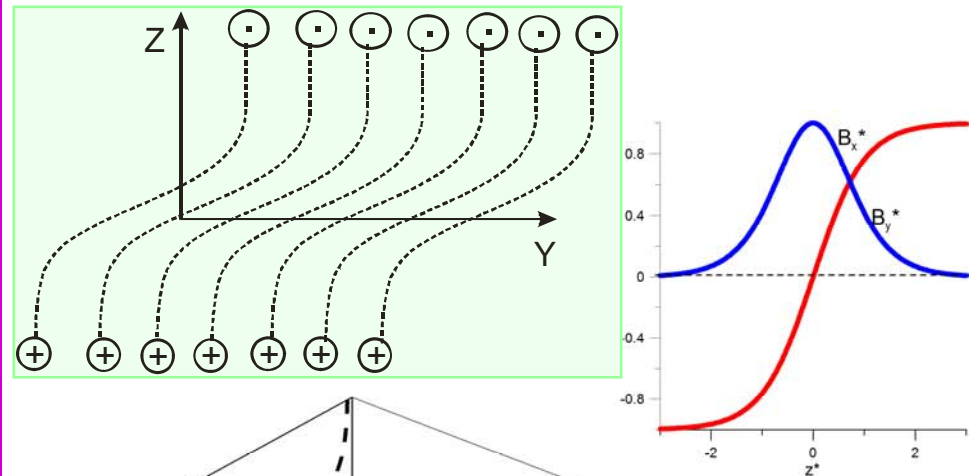


## Two magnetic field lines configurations $B_y(z)$

$B_y = B_y(z)$  antisymmetric

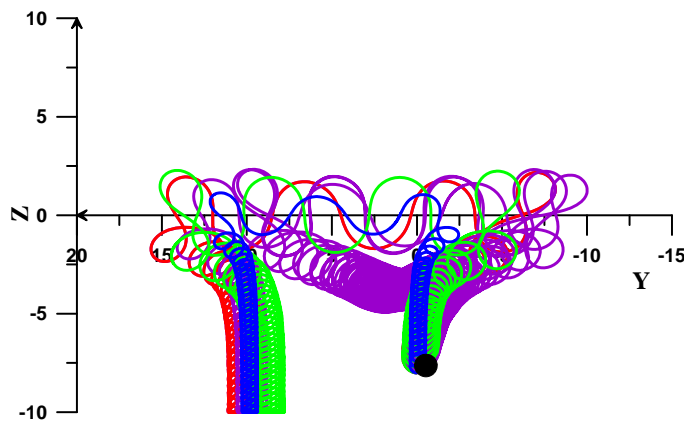
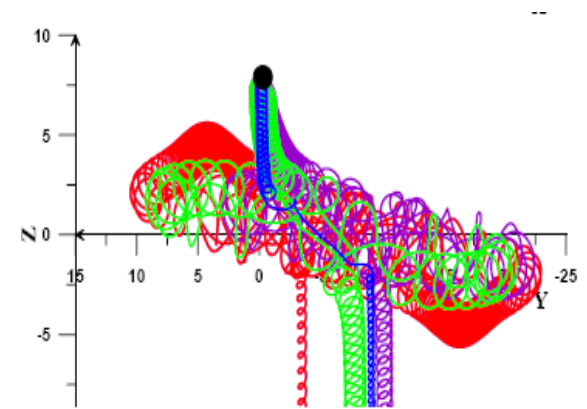
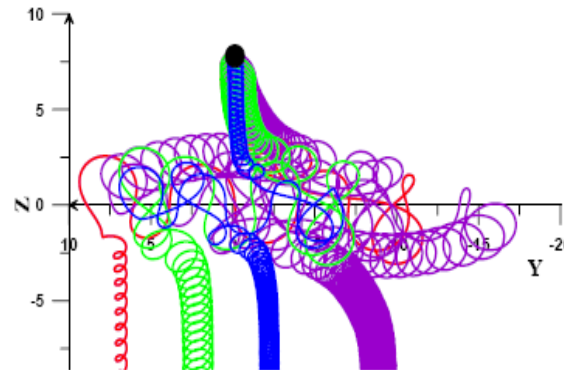
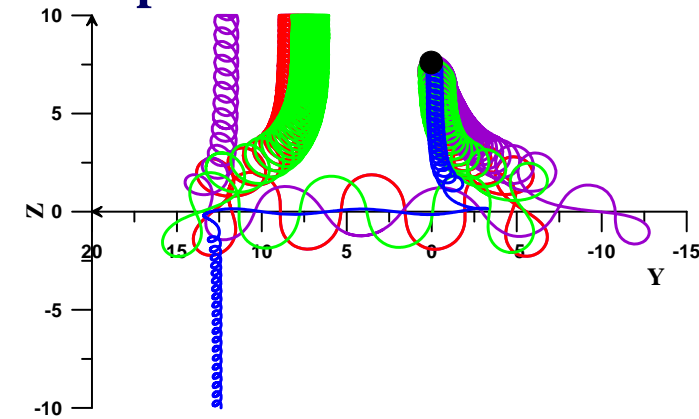


$B_y = B_y(z)$  symmetric

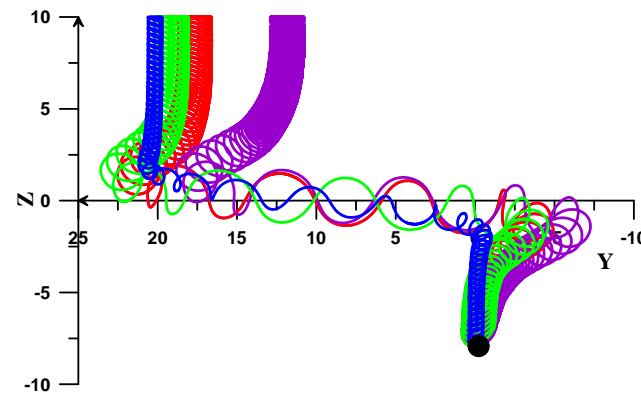


# Particle trajectories in thin current sheet with bell-shape (symmetric) $B_y(z)$

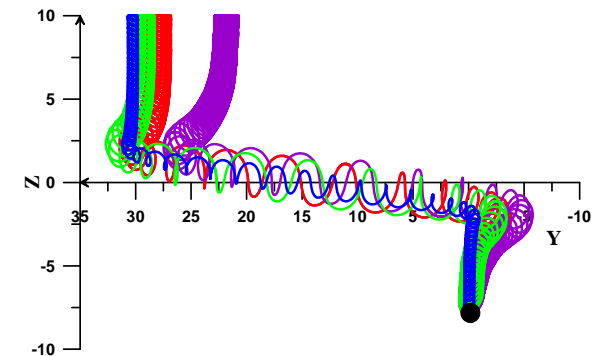
Amplitude of  $B_y(z)$  increases from the left to the right:  $B_{y0}/B_{x0}=0; 0.4; 0.8$ . The upper set of trajectories corresponds to the Northern source, the bottom set corresponds to the Southern one.



$B_{y0}/B_{x0}= 0;$



$0.4;$

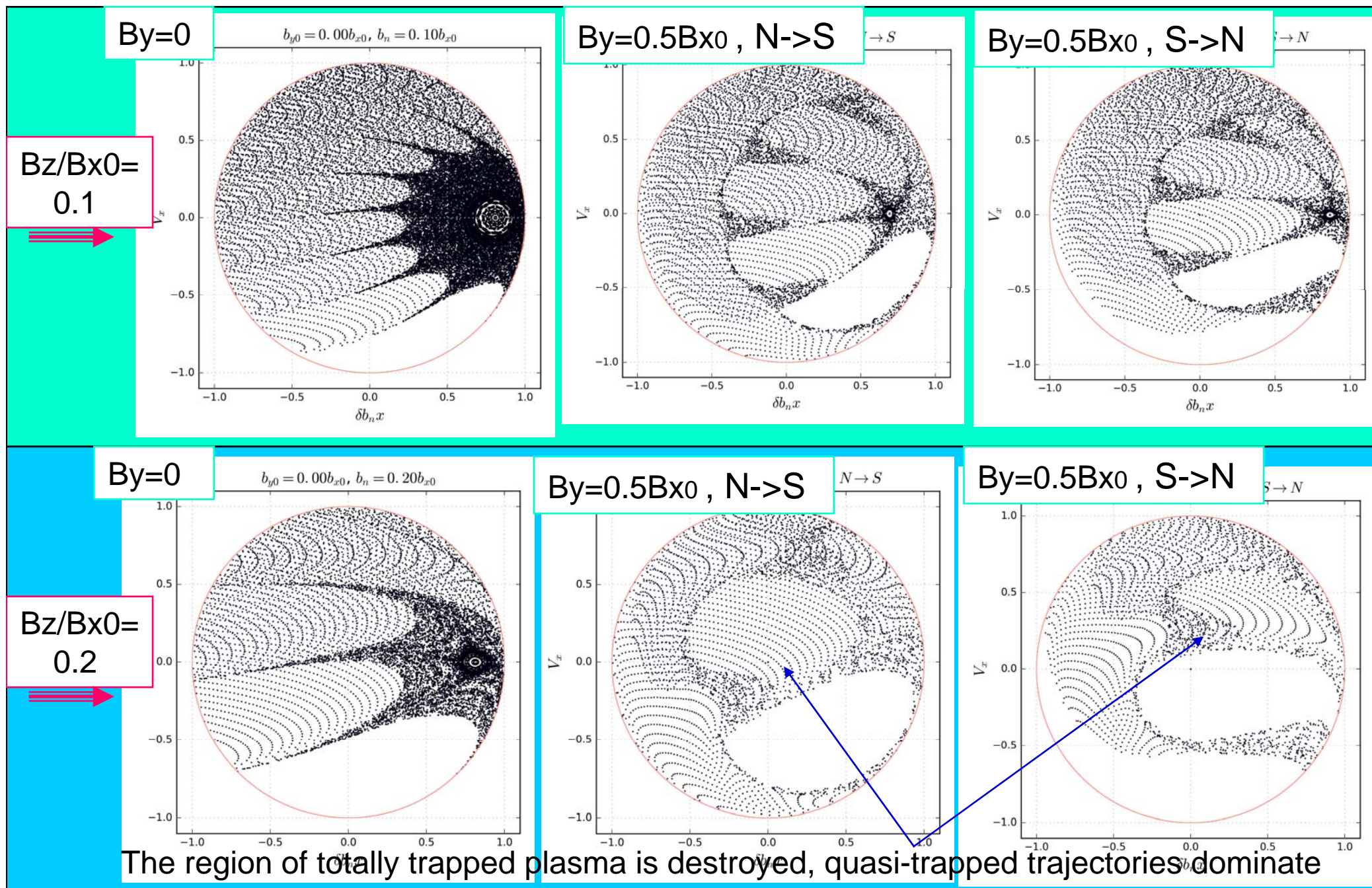


$0.8$

General effect: TCS is thickened due to the tilt of meandering plane relatively the neutral plane

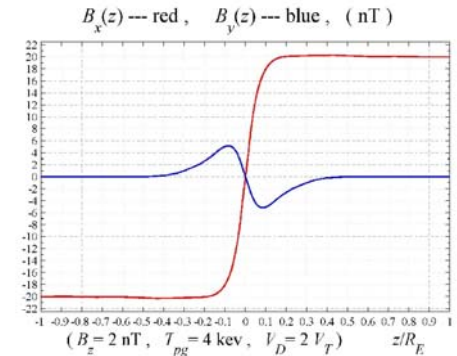


# Poincaré sections of quasi-adiabatic particle motion (By is bell-shaped)

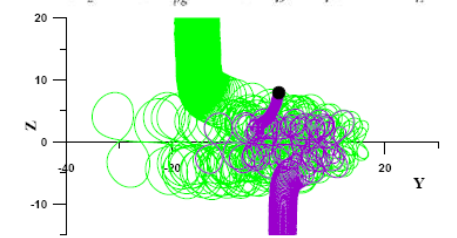
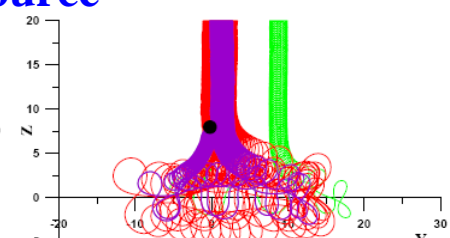
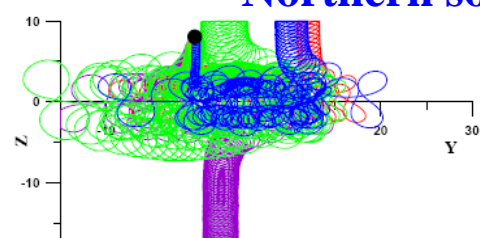
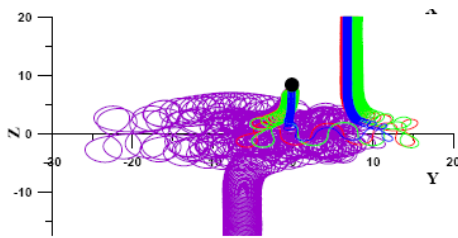


## Particle motion in the presence of the asymmetric $B_y(z)$ component

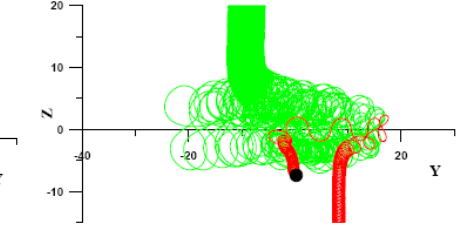
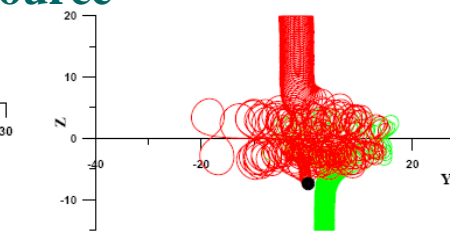
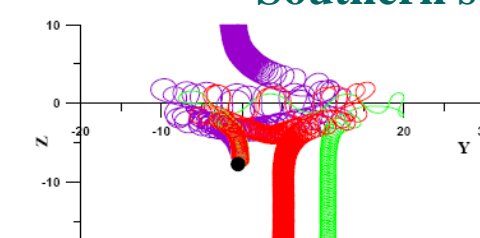
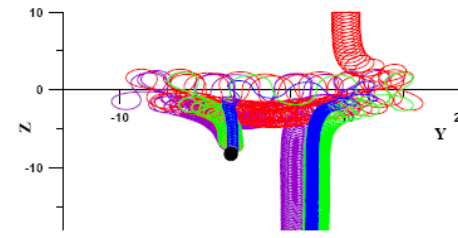
In the presence of asymmetrical  $B_y$  shear component the characteristic changes of  $B_y \sim \sin(z/L)$  are less than the transverse scale of meandering motion  $\rho \sim L$ . As a result most of particles with large Larmor radii almost do not feel the changes of magnetic field and their motion is not influenced strongly by the shear magnetic component. As a result the QA character of motion is conserved and the general part of current in thin current sheet is carried by ions in meandering (serpentine-like) orbits



Northern source



Southern source



$B_y/B_{x0} = 0.1,$

$0.2,$

$0.3$

$0.4$



# Poincaré sections of quasi-adiabatic particle motion (By is antisymmetric)

Particle scattering is the same in N->S and S->N directions

By=0

$b_{y0} = 0.00b_{x0}, b_n = 0.10b_{x0}$

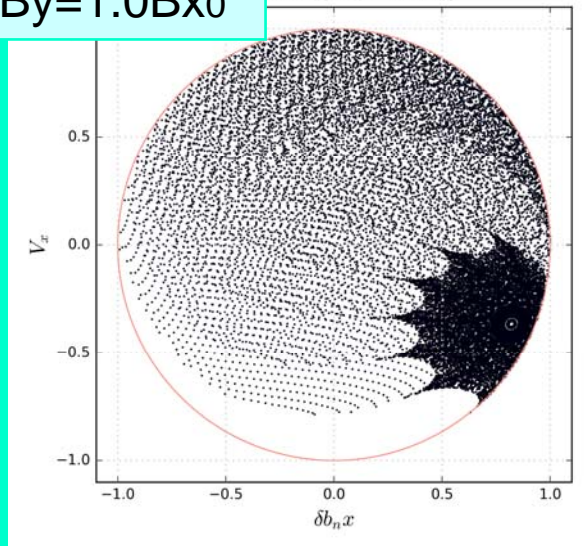
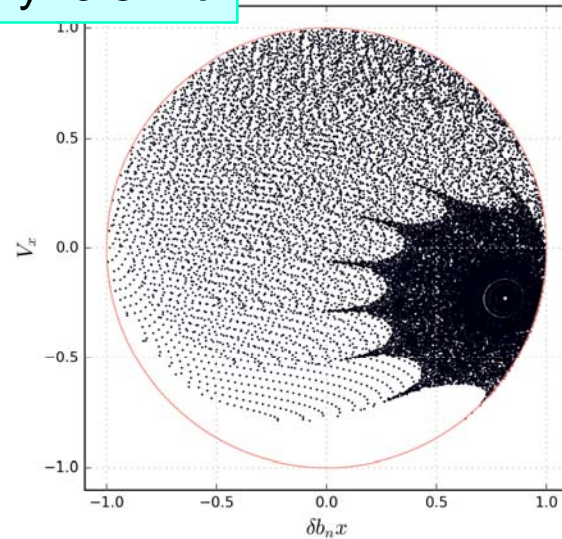
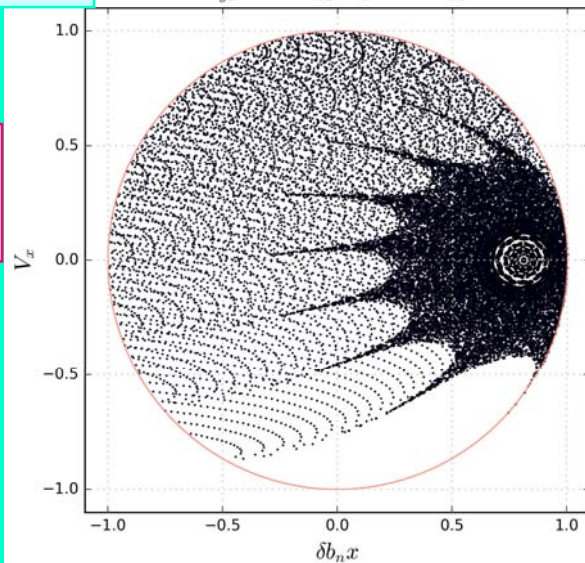
By=0.5Bx0

$b_{y0} = 0.50b_{x0}, b_n = 0.10b_{x0}$

By=1.0Bx0

$b_{y0} = 1.00b_{x0}, b_n = 0.10b_{x0}$

Bz/Bx0=  
0.1



By=0

$b_{y0} = 0.00b_{x0}, b_n = 0.20b_{x0}$

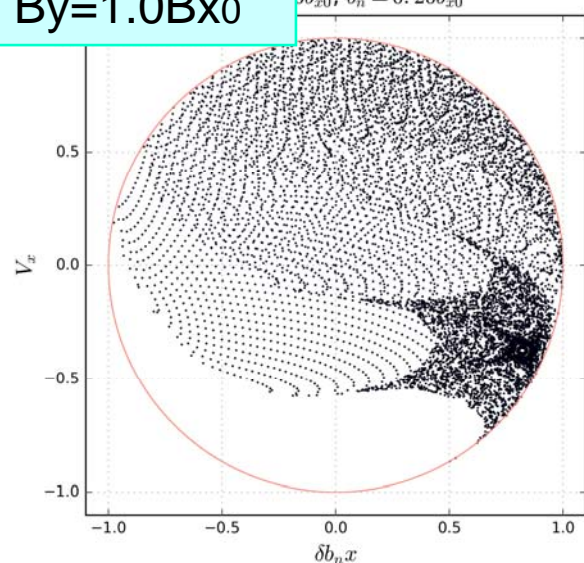
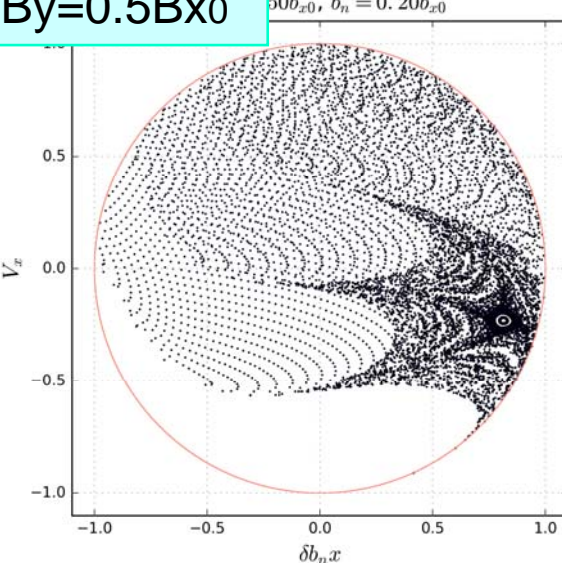
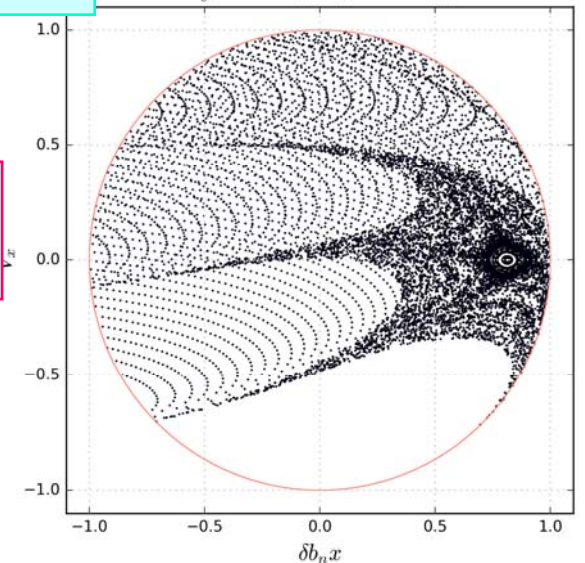
By=0.5Bx0

$b_{y0} = 0.50b_{x0}, b_n = 0.20b_{x0}$

By=1.0Bx0

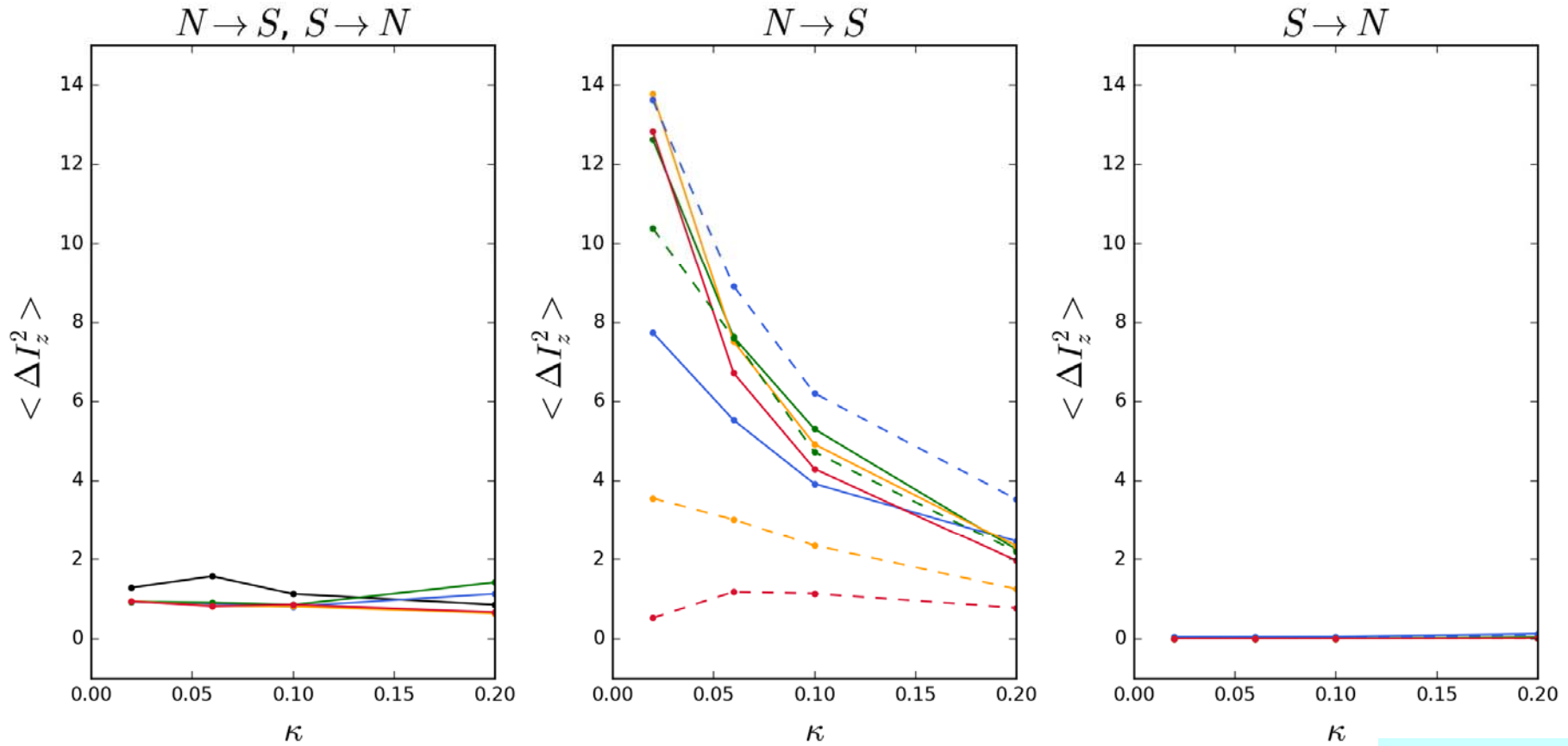
$b_{y0} = 1.00b_{x0}, b_n = 0.20b_{x0}$

Bz/Bx0=  
0.2

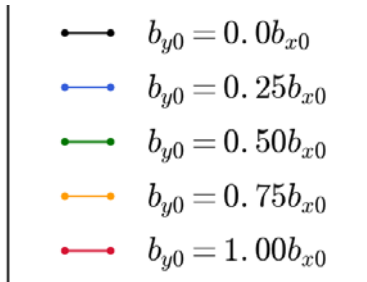


# The squares of the values of adiabatic invariant $I_z$ jumps

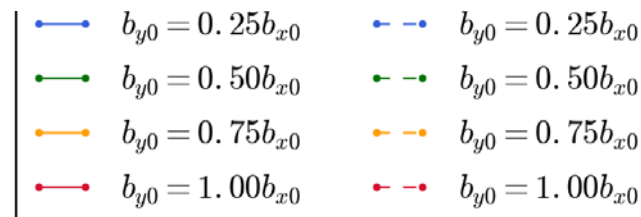
as functions of parameter of adiabaticity  $\kappa$



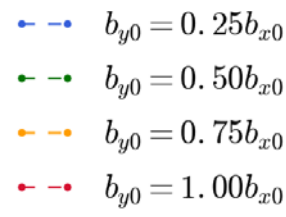
Shearless and antisymmetric By



Bell-shaped



constant By



*parameter of adiabaticity*

$$\kappa = \sqrt{\frac{R_{\min}}{\rho_{\max}}}$$

$$R_{\min} = \left. \frac{b_n}{db/dz} \right|_{z=0}$$

## Conclusions:

Numerical analysis of particle dynamic in thin current sheets (TCSs) taking (or do not taking) into account the magnetic shear component has shown that:

- 1) Particle dynamics is different in TCSs having three general shapes of magnetic shear component: (A) constant; (B) symmetric; (C) antisymmetric relatively neutral plane; and in (D) shearless configuration.
- 2) In the cases (A) and (B) particle scattering near the neutral plane is asymmetric and depends on the location of the plasma source in Northern or Southern hemispheres. Ions originated from the Northern hemisphere are scattered stronger in comparison with particles from the Southern hemisphere, their jumps of quasi-adiabatic invariants  $I_z$  tends to small value proportional to the value of  $B_y$  component. The direction of scattering depends on the sign of  $B_y$  (and has mirror symmetry relatively  $z=0$  plane). As a result the plasma density from both sides of CS plane can be different.
- 3) In the cases (C) and (D) particle scattering does not depend on plasma source location, and there is no difference between the jumps of particles going from the North or the South toward current sheet. The Poincaré sections and  $I_z$  jumps are identical for particles with the same parameters going in N->S and S->N directions.
- 4) Current sheet thickness is larger in configurations (A) and (B) due to geometric effect (the tilting of a meandering plane of ions). CS thickness does not change in the configuration (C) in comparison with (D).



**Thank you for your attention!**