

BY

Technique to inverse the planetary dipolar magnetic field based on the current loop model: Sounding the interior via magnetic feld Zhaojin Rong (戎昭金)

Email: rongzhaojin@mail.iggcas.c.cn

Key Laboratory of Earth and Planetary Physics Institute of Geology and Geophysics, Chinese Academy of Sciences



## Earth has a dipolar magnetic field





### Geomagnetic field is generated by Geodynamo currents



- Where is the dynamo zone?
- How is the scale?
- Current orientation?
- Current strength?

It's hard to probe the real dynamo currents since nobody can cut open Earth

### Inversion of Geomagnetic field



BΥ

Since geomagnetic field is generated by the dynamo current, geomagnetic field could be the most direct means to diagnose the **real** geodynamo current



### **Geomagnetic field model**

**Spherical Harmonic Analysis** 

 $\nabla \times \mathbf{B} = 0$  $\mathbf{B} = -\nabla V$ 

Geomagnetic observatories





### SHC tells few about the source of dynamo

- > The potential field is required for the solution by SHC
- The dominant term of dipole component implies infinite value at Earth center
- An offset of a dipole center is mathematically indistinguishable from a dipole plus a series of higher orders multipole components [Merrill, McElhinny, & McFadden, 1996]



### Models with "more physics" are required



#### **Current loop**



### **Advantages of loop model**

It can roughly indicate the geometric information of dynamo



- Loop center(x0,y0,z0): dynamo location
- Axis ( $\vartheta_0, \varphi_0$ ) orientation of dynamo current
- Radius (a) dynamo scale
- Carried current (I) dynamo strength

Loop models are tried in many studies [e.g. Zidarov and Petrova, 1974; Peddie, 1979; Zidarov, 1985; Demina and Farafonova, 2016]



### **Difficulties of loop model**

**Difficulties:** the simultaneous fitting of these loop parameters would result in multiple error minima in the space of parameters. The fitting results are strongly dependent on the initial input values.

The sum-of-squares function for eccentric current loops may have multiple minima in the parameter space. That is, in general, least squares solutions for current loops are not unique. Unfortunately, no general technique exists for determining whether a particular minimum is the desired universal minimum, that is, the best possible solution. To increase the likelihood of finding the universal minimum, a random search technique was used. For each model, 20 preliminary computer

> by N.W.Peddie [JGR,1979], the generator of 3<sup>rd</sup> IGRF model

More physics, but more complicated calculations!





Motivation: Trying to develop a new technique to separate and solve the loop parameters successively based on the sampled magnetic field data.

**Target**: quick, effective, and applicable. It could be applied generally to planetary dipole field and paleomagnetic field





### **Method: Setup of loop coordinates**



Origin is at planet center O

Current loop is in the Xloop-Yloop plane



### **Method: Parameter separation**

The loop parameters can be separated according to the field geometry of loop current

- Dipole center,r0 (x0,y0,z0), related with displacement of field structure.
- Axis orientation,  $M(\vartheta_0, \varphi_0)$ , related with the orientation of field structure;
- Loop radius, a, related with the geometric configuration of magnetic field;
- Current, I, only related with the field strength



**Assumption:** Sampled field should be purely internal field, or the external field can be evaluated and subtracted.



### Method: 1. axis and loop center

The best axis orientation should make the projected field lines radially orientated for best

The projected position vector of S/C is

$$\mathbf{r}_{ip}\left(x_{i}^{'}\hat{\mathbf{x}}_{loop}, y_{i}^{'}\hat{\mathbf{y}}_{loop}\right)$$

The projected unit field vector is

$$\mathbf{b}_{ip}\left(b_{xi}^{'}\hat{\mathbf{x}}_{loop},b_{yi}^{'}\hat{\mathbf{y}}_{loop}\right)$$

The projected loop center is

$$\mathbf{r}_{0p}\left(x_{0}^{'}\hat{\mathbf{x}}_{loop}, y_{0}^{'}\hat{\mathbf{y}}_{loop}\right)$$



$$\sin \alpha_{i} = \frac{\mathbf{r}_{ip} \times \mathbf{b}_{ip}}{\left|\mathbf{r}_{ip}\right| \left|\mathbf{b}_{ip}\right|} = \frac{\left(x_{i}^{'} - x_{0}^{'}\right) b_{yi}^{'} - \left(y_{i}^{'} - y_{0}^{'}\right) b_{xi}^{'}}{\sqrt{b_{xi}^{'2} + b_{yi}^{'2}} \sqrt{\left(x_{i}^{'} - x_{0}^{'}\right)^{2} + \left(y_{i}^{'} - y_{0}^{'}\right)^{2}}}$$

$$\sigma(x_0', y_0', \theta_0, \varphi_0) = \frac{1}{N} \sum_i a \sin(|\sin \alpha_i|)$$

the optimum of **M** and  $\mathbf{r}_{0p}$ should make  $\sigma$  reaches the global minimum



#### Method: 1. axis and loop center

Each projected field line should cross the loop center

$$b'_{yi}\left(x'_{0} - x'_{i}\right) - b'_{xi}\left(y'_{0} - y'_{i}\right) = 0$$

$$A = \begin{pmatrix} b'_{y1} & -b'_{x1} \\ b'_{y2} & -b'_{x2} \\ \dots & \dots \\ b'_{ym} & -b'_{xm} \end{pmatrix}, X = \begin{pmatrix} x'_{0} \\ y'_{0} \end{pmatrix}, Y = \begin{pmatrix} x'_{1}b'_{y1} - y'_{1}b'_{x1} \\ x'_{2}b'_{y2} - y'_{2}b'_{x2} \\ \dots \\ x'_{m}b'_{ym} - y'_{m}b'_{xm} \end{pmatrix}$$



The optimum solution of X is  $\downarrow$ 

 $X = \left(A^T A\right)^{-1} A^T Y \leftrightarrow$ 

The optimum loop center is function of axis orientation ( $\vartheta_0, \varphi_0$ ).

$$\sigma(\theta_0,\varphi_0) = \frac{1}{N} \sum_i a \sin(|\sin\alpha_i|)$$

ų,

 $\sigma$  is only the function of axis orientation, the global minimum of  $\sigma$  can be quickly searched in the 2-D map constituted by  $\vartheta_0$  and  $\varphi_0$ . Correspondingly, the optimum axis and loop center can be found.



The ideal loop field component

$$\tilde{B}_{ir} = \frac{\mu_0 I}{4\pi} \frac{2\cos\theta_i}{\sin\theta_i \sqrt{a^2 + R_i^2 + 2aR_i \sin\theta_i}} \left[ \frac{a^2 + R_i^2}{a^2 + R_i^2 - 2aR_i \sin\theta_i} E - K \right] \qquad \tilde{B}_{iz} = \frac{\mu_0 I}{4\pi} \frac{2}{\sqrt{a^2 + R_i^2 + 2aR_i \sin\theta_i}} \left[ \frac{a^2 - R_i^2}{a^2 + R_i^2 - 2aR_i \sin\theta_i} E + K \right]$$

 $\tilde{\mathbf{b}}_i = \frac{\tilde{B}_{ir}}{B} \hat{\mathbf{r}} + \frac{\tilde{B}_{ir}}{B} \hat{\mathbf{M}}$ 

The ideal unit field vector is function of radius *a* and the shift component of loop center *z0*'.

 $\mathbf{b}_i = \frac{\mathbf{B}_i}{\left|\mathbf{B}_i\right|}$ 

#### The sampled unit field vector

 $\gamma_i = a \cos\left(\mathbf{b}_i \cdot \tilde{\mathbf{b}}_i\right)$  $\varepsilon\left(z_0, a\right) = \frac{1}{N} \sum_i \gamma_i$ 

The global minimum of  $\mathcal{E}$  can be quickly searched in the plane of  $(z'_0, a)$ . Correspondingly, the optimum radius a and the axis-shift component  $z'_0$  can be solved



### **Method: current and error**

#### current

$$\begin{split} \tilde{B}_{ir} &= IF_{ir} \qquad F_{ir} = \frac{\mu_0}{4\pi} \frac{2\cos\theta_i}{\sin\theta_i \sqrt{a^2 + R_i^2 + 2aR_i\sin\theta_i}} \left[ \frac{a^2 + R_i^2}{a^2 + R_i^2 - 2aR_i\sin\theta_i} E - K \right] \\ \tilde{B}_{iz} &= IF_{iz} \qquad F_{iz} = \frac{\mu_0}{4\pi} \frac{2}{\sqrt{a^2 + R_i^2 + 2aR_i\sin\theta_i}} \left[ \frac{a^2 - R_i^2}{a^2 - R_i^2 - 2aR_i\sin\theta_i} E + K \right] \end{split}$$

Since magnetic axis, loop radius and loop center have been solved, current can be derived as

$$I = \frac{1}{N} \sum_{i} \frac{B_i}{\sqrt{F_{ir}^2 + F_{iz}^2}}$$

#### **Method error**

$$\delta = \frac{1}{N} \sum_{i} \frac{|\mathbf{B}_{it} - \mathbf{B}_{it0}|}{|\mathbf{B}_{it}|} \qquad \mathbf{B}_{it} : \text{ sampled field} \\ \mathbf{B}_{it0} : \text{ field of loop}$$





#### **Input loop parameters**

Loop center: (x0=0.1;y0=0.2;z0=0.5) m Axis orientation:  $(\theta_0=60^\circ, \phi_0=40^\circ)$ radius: a=0.5m current: I=0.35A

#### along arbitrary trajectory

$$(x=-2, y=-2, z=-2)$$
 m to  $(x=2, y=2, z=2)$  m  
Sampled 20 data points evenly.



Sampled field vector

**Target:** Applying our technique to the sampled data to inverse the input parameters





#### Test results:

Loop center: (x0=0.1;y0=0.2;z0=0.5) m Axis orientation:  $(\theta_0=60^\circ, \phi_0=40^\circ)$ radius: a=0.5m current: I=0.35A

#### It is able to inverse parameters exactly

Result: Yes! Quick, Stable, and Effective!





# **Application-IGRF**



IGRF: 2015-01-01 00:00:004 orbits, variable altitudeEach orbit, 20 data pointsIn total, 80 data points

# **Application-IGRF**

Altitude	<i>x</i> 0+'	<i>Y</i> 0	<i>Z</i> 0	<b>a</b> *	<i>I</i> *	$M^{\mathtt{a}_{\mathtt{s}^{j}}}$	<b>θ</b> 0↔	<i>\$</i> 0₽	σ <sub>min</sub> ₄	<b>€</b> ₽	$\delta^{\!\scriptscriptstyle  m P}$	*
(km).	(km) <sub>*</sub>	(km) <sub>*</sub>	(km) <sub>"</sub>	(km) <sub>"</sub>	(*10 <sup>10</sup> A) <sub>°</sub>	(*10 <sup>22</sup> Am <sup>2</sup> ) <sub>*</sub>	(°),	(°),	(deg),.	(deg)		
0+2	<b>-286</b> @	309.0	111+2	856₽	3.4.	7.82*	<b>167.7</b> ₽	113.1.	7.465.	<b>8.079</b> ₽	0.19	*
100.0	-274	320 ¢	<mark>96</mark> ₽	817.	3.7.	7.81₽	167.8	118.0	7.261.	7 <b>.907</b> ₽	0.180	*
<b>5</b> 00₽	-293	<b>334</b> 4	71₽	745₽	4.50	7.7 <b>9</b> ₽	168.6	<b>117.4</b> ₽	<b>6.558</b> ₽	<b>6.876</b> ₽	0.16	*
1000+2	-311@	<b>339</b> ¢	<mark>8</mark> 3₽	<b>701</b> ₽	5.0 <i>-</i>	7.77₽	<b>169.1</b> @	115.60	<b>5.856</b> <sub>°</sub>	<b>6.016</b> ₽	0.14	*
2000	-334 @	348 @	<b>105</b> ₽	754₽	4.3.	7.75₽	<b>169.7</b> ₽	113.3.	4.843.	4.756₽	0.11	*
<b>5</b> 000₽	-364 @	<u>356</u> ∉	17 <b>0</b> ₽	720₽	4.7₽	7.72₊∘	171.1	111.8 <sub>°</sub>	3.333+	<b>2.979</b> ₽	0.07	*
10000+3	-381	356 .	<mark>206</mark> ⊬	<b>631</b> #	<b>6.2</b> <i>∉</i>	7.72≠	170.9	<b>109.4</b> ⊷	2.145.	1.833	0.05	ļ
20000	-390 @	<b>357</b> ∉	<b>212</b> ₽	353₽	<b>19.7</b> ₽	7.72₽	170.5	<b>108.0</b> ₽	1.276.	1.031	0.05	÷
50000 <sup>b</sup> .	Ŷ	$\sim$	$\sim$	$\sim$	~2	~	<b>170.4</b> ₽	<b>107.5</b> ₽	0.581+	$\sim$	~~	*
IGRF C <sub>e</sub>	<b>-</b> 400₽	352.0	221.0	$\sim$	~2	7.72₽	170.4~	107.4~	~	~ĵ	~~	*

Table1. The inversed loop parameters for the IGRF model of the year 2015.

Inversed results quickly converged to the eccentric-dipole of IGRF as the increase of altitude.





Four orbits with altitude 5000 km sample 80 data point at 1960:00:00:00



New technique is applicable and reasonable



#### 100 years variation of IGRF-loop model

Four orbits with altitude 5000 km sample 80 data point





#### 100 years variation of IGRF-loop model





The moment of the sampled data is arbitrarily set at 2015-01-01 00:00:00, and the originally sampled data of 123 observatories are used.

The original recorded data set subjected to the disturbance of external space currents and local field anomaly



Model₄	<i>x</i> ₀⊷	$y_0$	$z_0$	<i>a</i>	I e	$M$ $_{ m e}$	$ heta_{0^{*'}}$	<i>\$</i> 0₽	$\sigma_{\min^{\omega}}$	E₽	8	]
	(km).	(km),	(km),	(km),	(*10 <sup>10</sup> A) <sub>*</sub>	(*10 <sup>22</sup> Am <sup>2</sup> ),	(°),	(°),	(deg) <sub>e</sub>	(deg) <sub>e</sub>		
Loop model₊≀	<b>-21</b> 3 @	<b>403</b> *	128.0	<b>892</b> ₽	3.1.0	7.88₽	172.0	114.6	5.313.	8.423.	0.1	
											8.0	
IGRF 🖉	<b>-400</b> ¢	352₽	221.0	~	~~	7.72₽	170.4~	107.4.	~7	~	~	1



## **Further modifications**

#### $\checkmark$ 1. It can be reduced to a unconstrained dipole model

6 parameters are involved.

Loop center(3 parameters), axis orientaion (2 parameters),

and dipole moment(1 parameters)

Axis orientation and dipole center

Same with loop model

Loop radius=0; The shift component of dipole center is

$$\gamma_i = a \cos\left(\mathbf{b}_i \cdot \tilde{\mathbf{b}}_i\right)$$

 $\mathcal{E}\left(z_{0}^{'}\right) = \frac{1}{N} \sum_{i} \gamma_{i}$ 

The dipole moment is

$$B_{i} = MF_{i}$$

$$F_{i} = \frac{\mu_{0}}{4\pi R_{i}^{5}} \sqrt{1 + 3\left(\frac{\tilde{z}_{i}}{R_{i}}\right)^{2}}$$

$$M = \frac{1}{N} \sum_{i} B_{i} F_{i}^{-1}$$



## **Application to Jupiter**



Connerney et al.,2018,GRL

4 circular orbits with variable altitude Each orbit, sampled 20 data points In total, 80 data points obtained

1

x 10<sup>4</sup>



Results of loop model for the sampled data when r=3Rj,4j,and 5j, code terminated, and returned a=0.

# Results of reduced dipole model

#### Jupiter dipole center is eccentric !

Orbit distance (RJ)	x <sub>0</sub> km	y₀ km	z <sub>o</sub> km	M (*10 <sup>27</sup> Am²)	θ <sub>ο</sub> (°)	Φ <sub>0</sub> (°)	σ <sub>min</sub> (deg)	ε (deg)	δ
3RJ	-5594	-4903	263	1.52	8.9	168.0	3.2	2.3	0.05
4RJ	-5629	-4809	560	1.52	9.9	167.9	2.6	1.5	0.03
5RJ	-5709	-4734	920	1.52	10.3	166.9	2.3	1.2	0.02
wiki	~	~	~	1.52	10	159			



### Comparison with sampled data

In this example, the orbit radius is 5 RJ





# Application of the reduced dipole model -Mercury



Paper	M (Gaussian unit)	North Shift of dipole center	Dipole tilt angle
Anderson, et al.2011, science	195±10nT*RM^3	484 km	<3 degree
Anderson, et al.2012, JGR	190±10nT*RM^3	479 km	<0.8 degree
Johnson et al.,2012,JGR	190 nT*RM^3	478 km	0



# Application of the reduced dipole model -Mercury





## **Application to Mercury**

#### Scattering plot of the angle a=acos (**b**·**r**)



- 1. Night side data set is used to reduce the disturbance of space current
- 2. Check the inversion errors for different ranges of altitude
- 3. Check the inversion errors for different ranges of latitude



# **Application to Mercury**

To reach the minimum inversion error  $\delta$ :

- Altitude: <100km</p>
- night side data: 20:00< local time<04:00</p>
- Latitude range:40<lat<50 degree</p>

Paper	M (Gaussian unit)	North Shift of dipole center	Dipole tilt angle
Anderson, et al.2011, science	195±10nT*RM^3	484 km	<3 degree
Anderson, et al.2012, JGR	190±10nT*RM^3	479 km	<0.8 degree
Johnson et al.,2012,JGR	190 nT*RM^3	478 km	0
Our study	215nT*RM^3	480 km	3.7 degree



# **Further modifications**

# ✓ It could be modified to a model of spherical surface current

(7parameters are involved. 2 for the axis orientation, 3 for the spherical center, 1 for the spherical radius, and 1 for the surface current density).

This model could be closer to the real dynamo current pattern.

J=j0; J=j0\*sin(
$$\vartheta$$
); J=j0\*sin( $\vartheta$ )^2;

It will be our Next Study



## **Further modifications**



$$\mathbf{j}(\mathbf{r}) = j_0 \delta(r - r_s) \sin \theta \hat{\mathbf{\varphi}} \qquad r > rs,$$

$$\mathbf{B}(\mathbf{r}) = -\mu(\mathbf{r})\nabla\psi_{out}(\mathbf{r}) = \frac{-\mu(\mathbf{r})\mu_r r_s^3 j_0}{(\mu_r + 2)} \left(-2r^{-3}\cos\theta \hat{r} - r^{-3}\sin\theta \hat{\theta}\right)$$

It is ideal dipole field

We have to keep in mind that "The inversion is non-unique", we have to interpret the inversion results very carefully.



- 1. A new technique to diagnose the source of planetary dipolar field is developed based on a single circular current loop model.
- 2. The technique is able to separate the loop parameters successively.
- 3. This technique can be reduced to a dipole model or modified to a spherical surface model.
- 4. Tests and application to spacecraft's data demonstrate the effectiveness of this technique
- 5. It could be applied to study the local field anomaly after subtraction of main dipole/loop component.



# Thank You 谢谢!