

Lower hybrid waves at the magnetosheath separatrix region

Binbin Tang^{1*}, Wenya Li¹, Daniel Graham², Chi Wang¹, Yuri Khotyaintsev², Barbara L. Giles³, Per-Arne Lindqvist⁴, Robert E. Ergun⁵, Christopher T. Russell⁶, and James L. Burch⁷

(* bbtang@spaceweather.ac.cn)

¹ National Space Science Center, CAS, China

² Swedish Institute of Space Physics, Uppsala, Sweden

³ NASA Goddard Space Flight Center, Greenbelt, MD, USA

⁴ KTH Royal Institute of Technology, Stockholm, Sweden

⁵ Laboratory of Atmospheric and Space Physics, University of Colorado Boulder, Boulder, Colorado, USA

⁶ Department of Earth and Space Sciences, University of California, Los Angeles, California, USA.

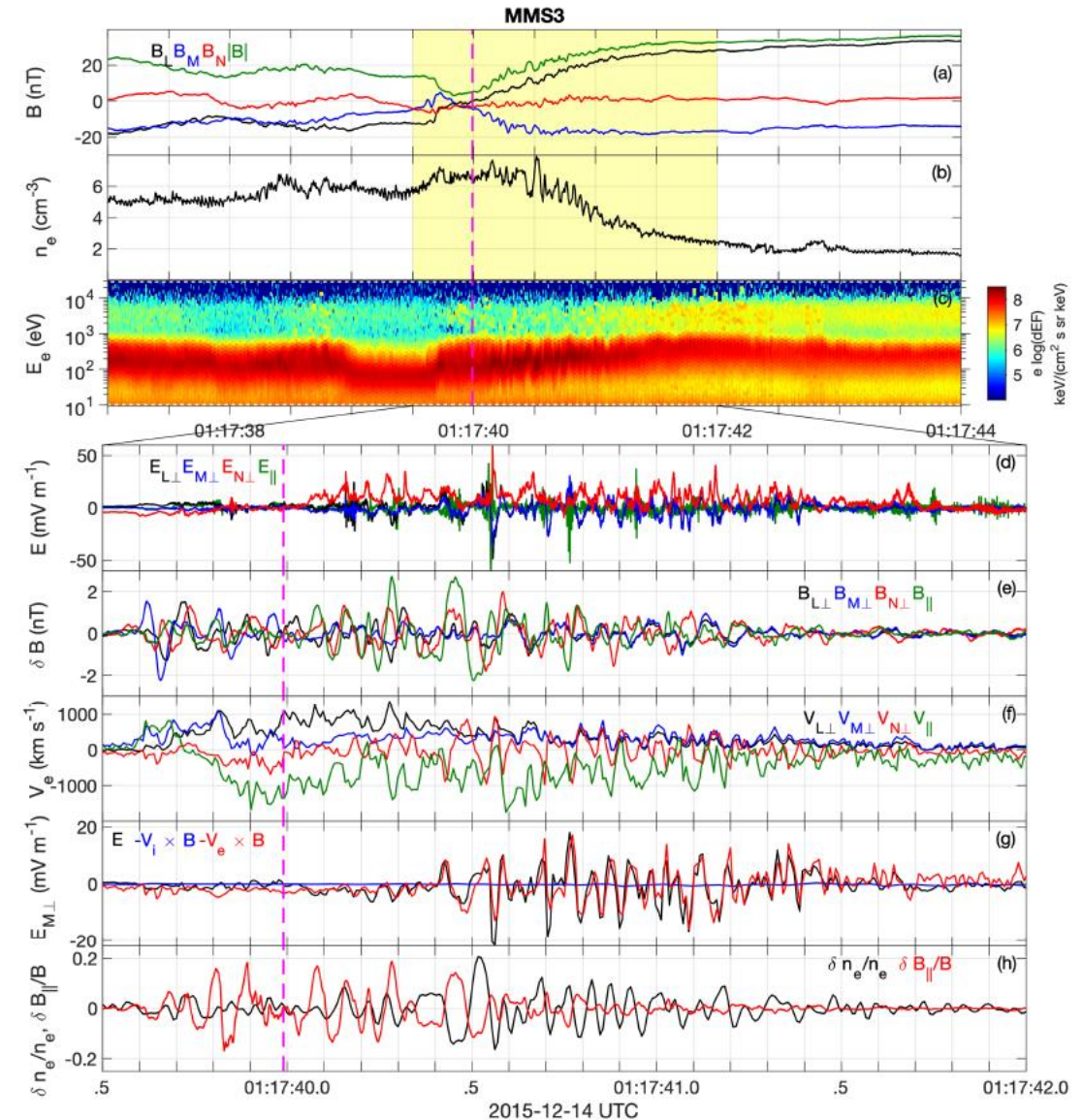
⁷ Southwest Research Institute, San Antonio, Texas, USA.

Basic properties of lower hybrid waves (LHWs)

- Frequency: $f_{ci} < f < f_{ce}$
- Wavelength: $\rho_i > \lambda > \rho_e$
 - electrons: approximated magnetized
 - ions: almost unmagnetized
- quasi-electrostatic, but $\delta \mathbf{B}$ are not zero (the ratio W_B/W_E is often >1 at magnetopause)
- $k_{\perp} \gg k_{\parallel}$
 - Linearly polarized and δB mainly in parallel directions
- More details can be found at Graham et al., [2019].

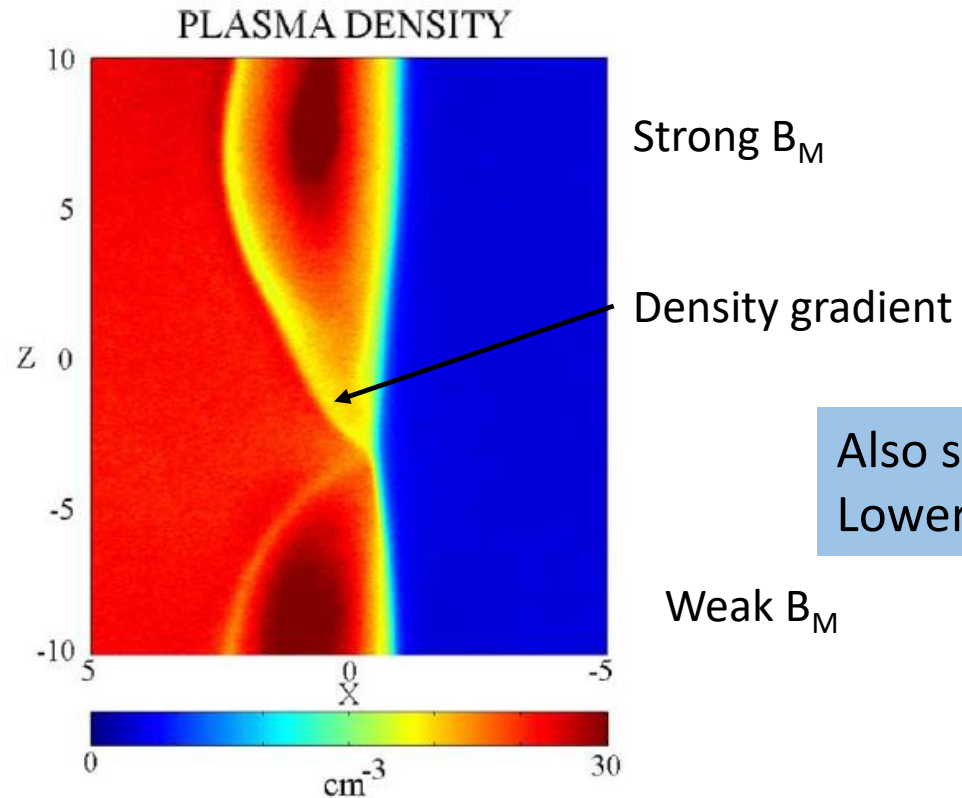
Lower hybrid waves (LHWs) at Earth's magnetopause

- LHWs are frequently observed at **magnetospheric side** of the magnetopause reconnection, which can be generated by
 - a. Lower hybrid drift instability: a density/temperature gradient
 - b. Modified two-stream instability: beam-like finite gyroradius ions
- LHWs are of particular importance
 - a. cross-field particle diffusion → built-up of LLBL
 - b. anomalous resistivity/ viscosity
 - c. ...



Asymmetric guide-field reconnection

- Asymmetry of the reconnection exhaust: $B_g \pm B_{\text{Hall}}$

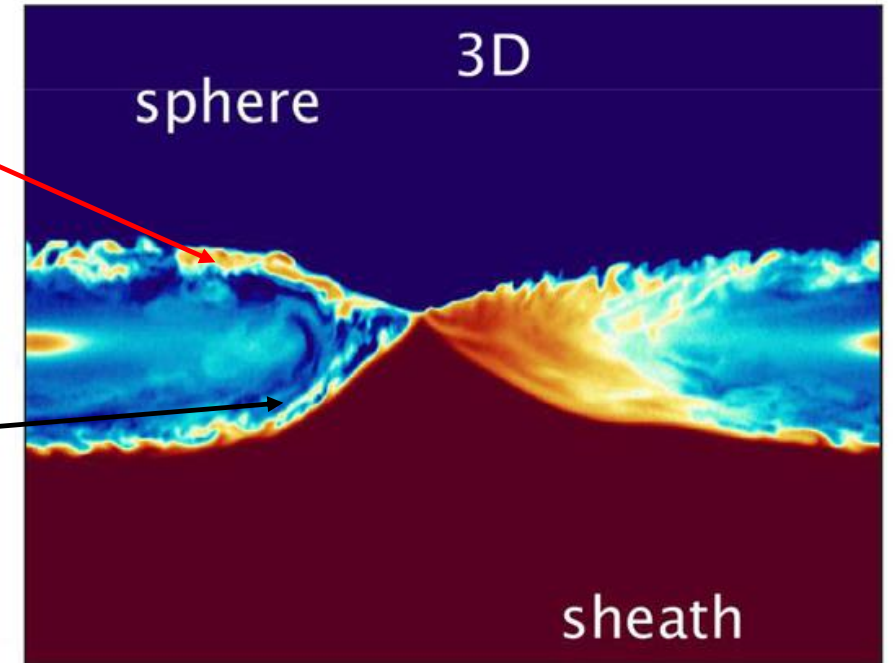


Mozer et al., 2008

Clear perturbations
of LHWs?

Also some perturbations.
Lower hybrid waves?

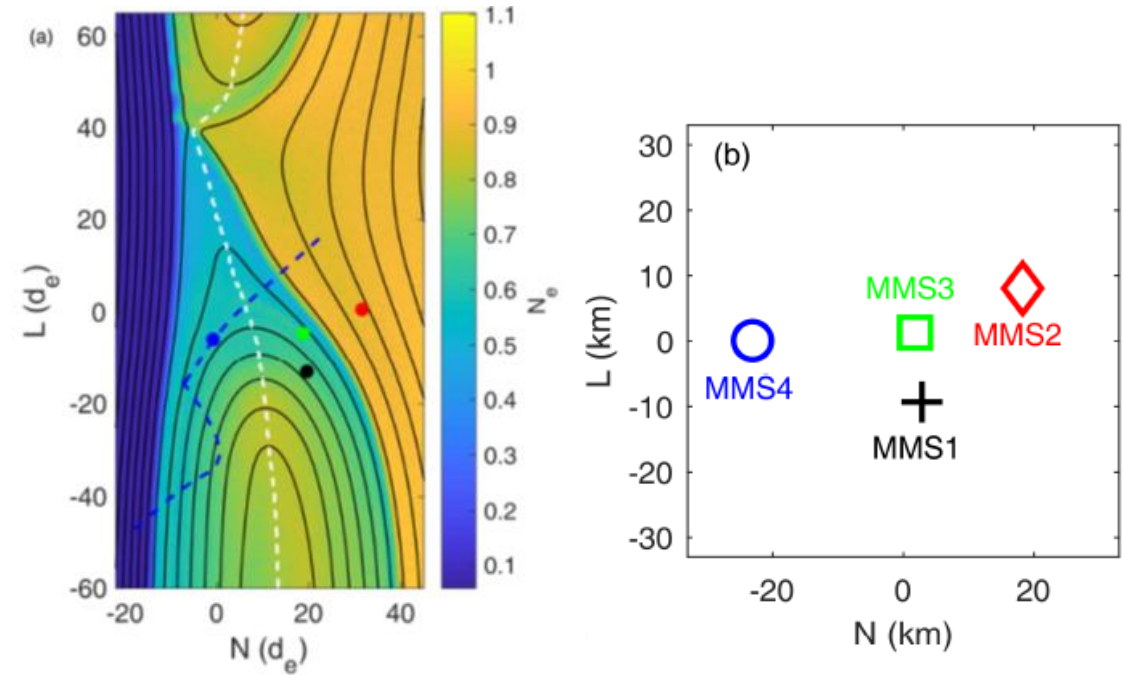
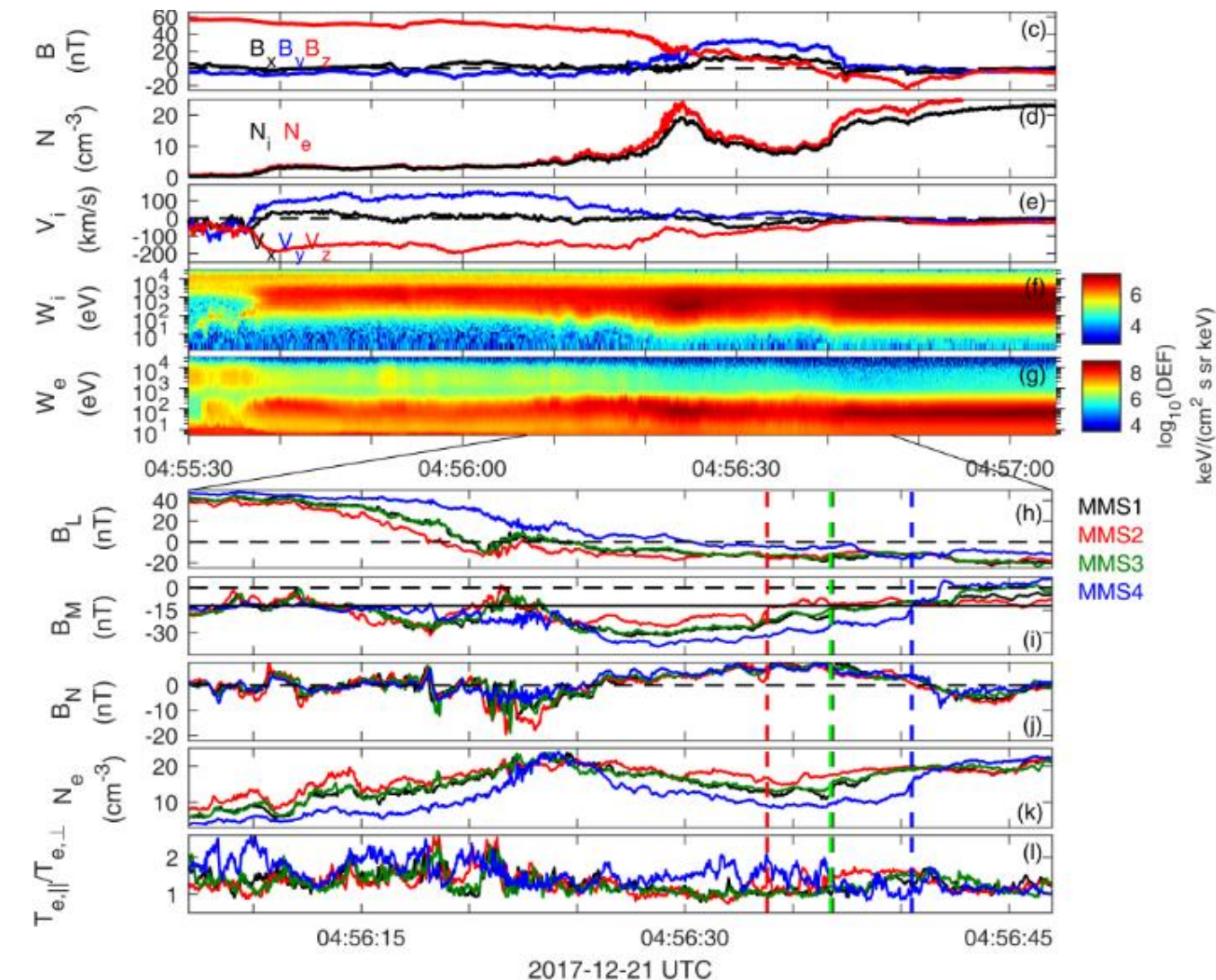
Run B $t\omega_{ci}=35$



$B_g = 0.4$

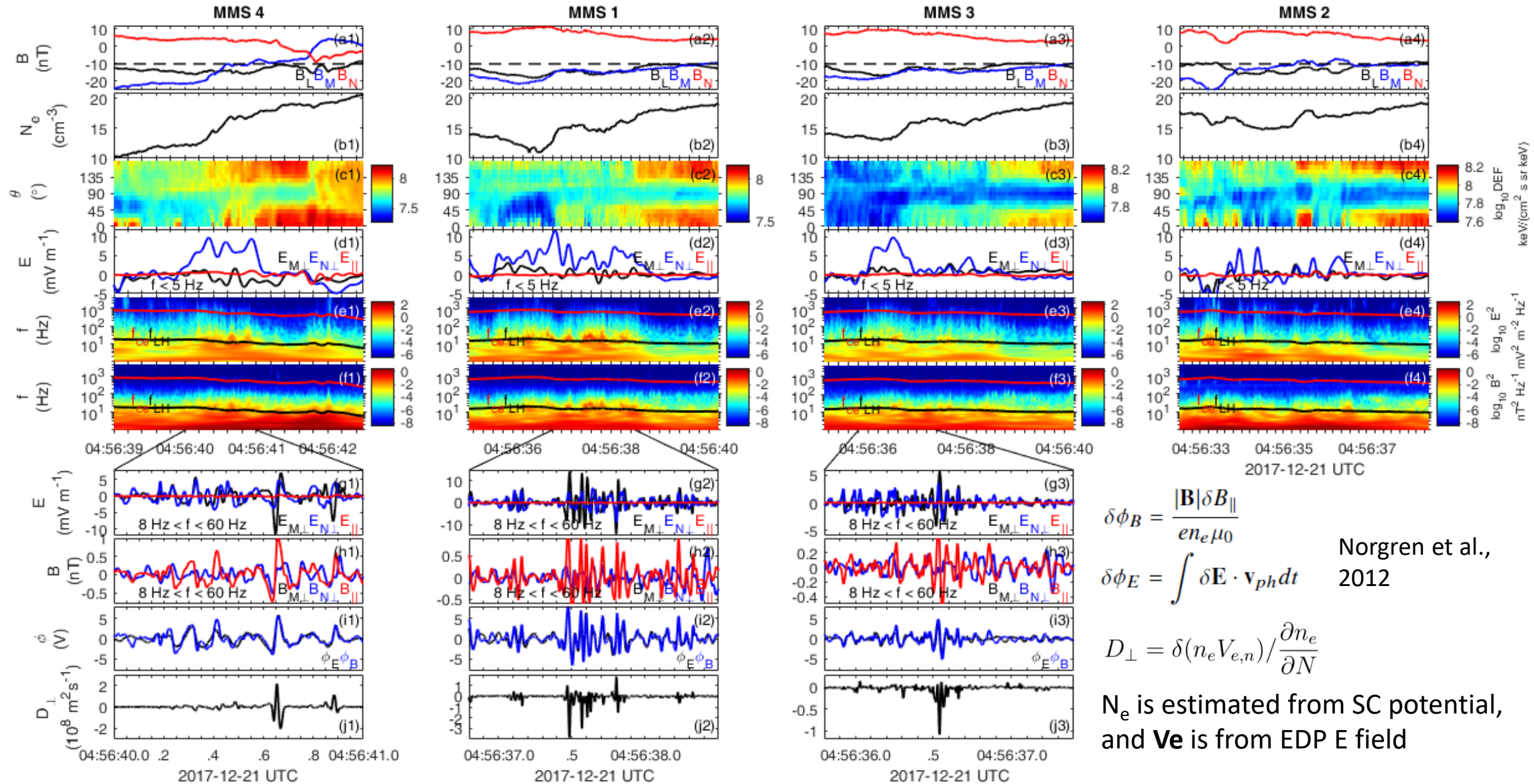
Le et al., 2018

Event overview



- From the x-line to downstream (due to the large tangential motion)
MMS 4 $\rightarrow$ MMS 1, 3 $\rightarrow$ MMS 2
- Slow normal motion ($V_{\text{mp},N} \sim 5 \text{ km/s}$)
- Significant different MMS observations
- Clear density gradient is found at the sheath separatrix (marked by the vertical dashed line)

Zoom-in at the magnetosheath separatrix



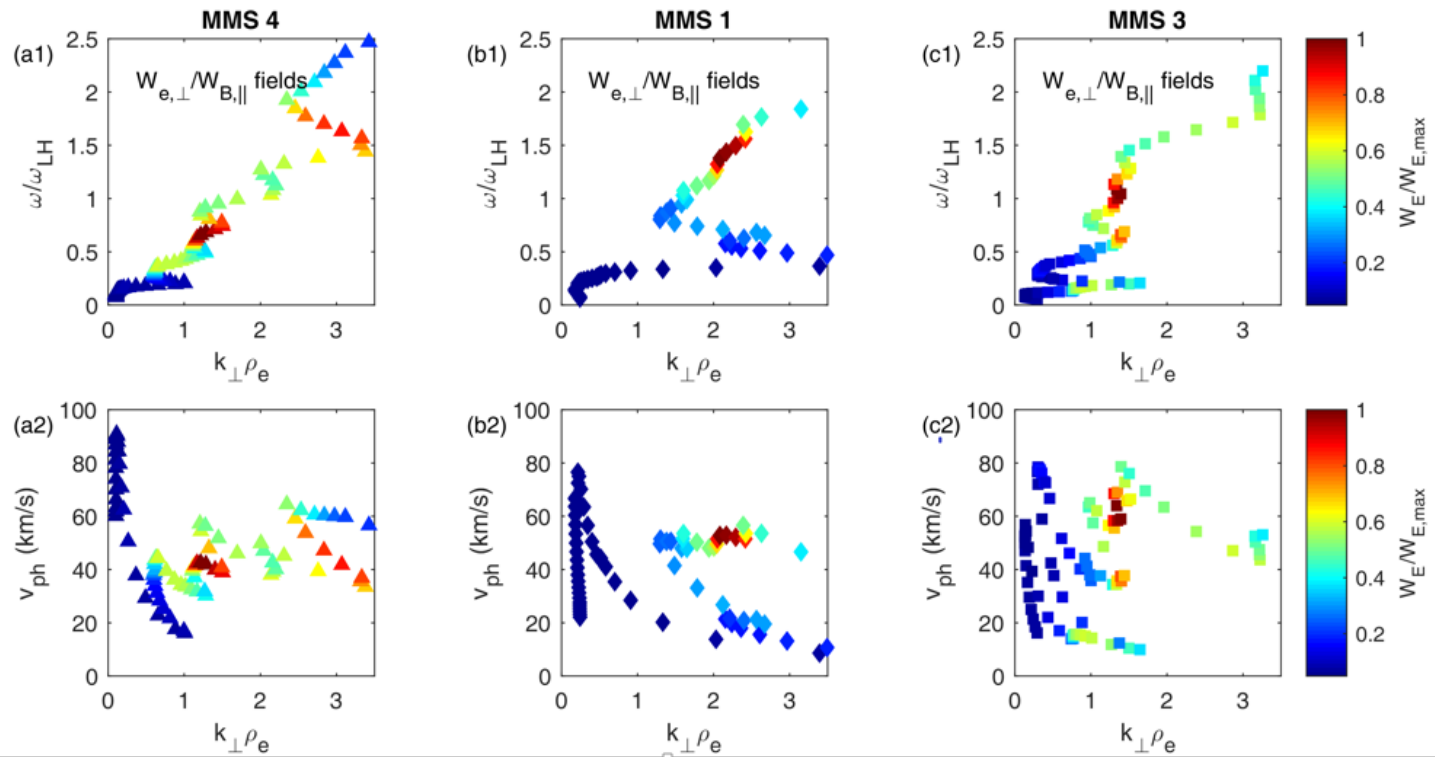
Properties of observed LHWs

	V_{ph} (km/s)	λ (km)	$k_{\perp}\rho_e$	Φ_B (V)	Φ_E (V)	C_{Φ}	$\langle D_{\perp} \rangle$ (m ² s ⁻¹)
MMS4	$50 \times [0.62 \ -0.71 \ -0.33]$	4.2	1.0	5.7	5.8	0.83	$\sim 10^5$
MMS1	$65 \times [0.72 \ -0.70 \ 0.06]$	3.1	1.6	7.8	8.4	0.87	-2.6×10^7
MMS3	$94 \times [0.67 \ -0.72 \ -0.18]$	6.3	0.9	5.2	4.5	0.81	-1.8×10^7

- Due to the presence of the magnetic B_M , the propagation direction of these waves is not closed aligned in -M direction, but has a component in L direction. → For comparison, LHWs at the magnetospheric side propagate approximately in M direction (in the ion rest frame).
- $k_{\perp}\rho_e \sim 1 - 1.6$
- The sheath electron temperature is about 30 eV, leading to $e\Phi/k_B T_e \sim 15\% - 25\%$
- $\langle D_{\perp} \rangle$ here is 1~2 orders in magnitude smaller than it in the magnetospheric side. Still it will only take several seconds to diffusely transport electrons if assuming the boundary is ~ 10 km thick.

Properties of observed LHWs

	k (LMN)	V_{ph} (km/s)	Φ_B (V)	Φ_E (V)	C_Φ	$k_\perp \rho_e$	V_{ph} (km/s)	f (Hz)	$k_\perp \rho_e$	$\langle D_\perp \rangle$ (m ² s ⁻¹)
	Norgren et al. [2012]						Graham et al. [2019]			
MMS4	[0.62 -0.71 -0.33]	50	5.7	5.8	0.83	1.0	43	8.6	1.2	$\sim 10^5$
MMS1	[0.72 -0.70 0.06]	65	7.8	8.4	0.87	1.6	53	20.4	2.1	-2.6×10^7
MMS3	[0.67 -0.72 -0.18]	94	5.2	4.5	0.81	0.9	60	17.3	1.4	-1.8×10^7

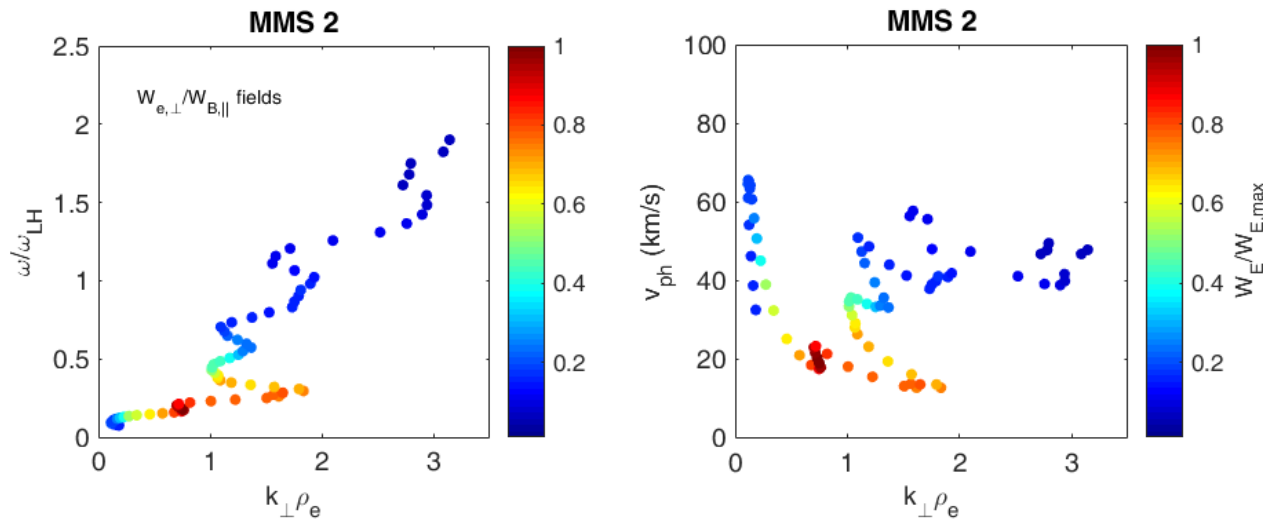


A new method is developed by Graham et al., 2019 to get the wave dispersion relation.

$$k_\perp(\omega) = \frac{1}{d_e} \sqrt{\frac{W_e(\omega)}{W_B(\omega)}}, \quad \text{Graham et al., 2019}$$

N_e is estimated from SC potential, and $\mathbf{V_e}$ is from EDP E field

	k (LMN)	V_{ph} (km/s)	Φ_B (V)	Φ_E (V)	C_Φ	$k_\perp \rho_e$	V_{ph} (km/s)	f (Hz)	$k_\perp \rho_e$	$\langle D_\perp \rangle$ (m ² s ⁻¹)
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In general, the result from different methods are consistent, but the V_{ph} estimated from Daniel's method seems smaller ($k_\perp$ is larger).

Again, no clear LHWs at MMS 2

Instability analysis (a finite plasma beta, ion frame)

$$0 = 1 - \frac{\omega_{pi}^2}{k^2 v_i^2} Z' \left(\frac{\omega}{k v_i} \right) + \frac{\omega_{pe}^2}{\Omega_{ce}^2} \left(1 + \frac{\omega_{pe}^2}{c^2 k^2} \right) + \frac{2\omega_{pe}^2}{k^2 v_e^2} \left(1 + \frac{\beta_i}{2} \right) \frac{k V_{de}}{\omega - k V_E},$$

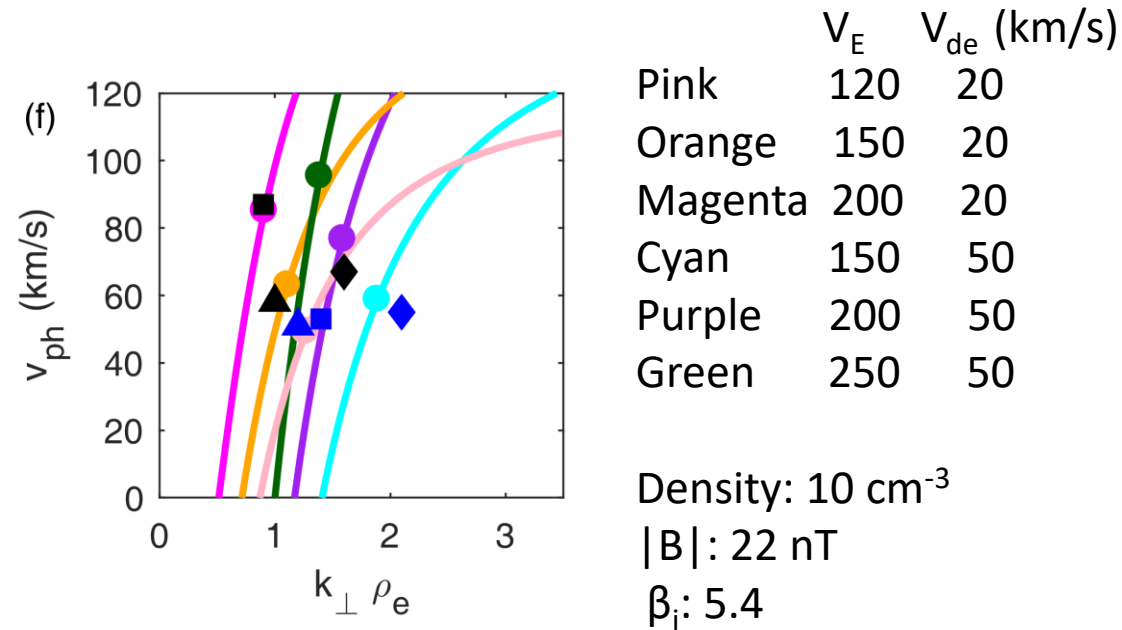
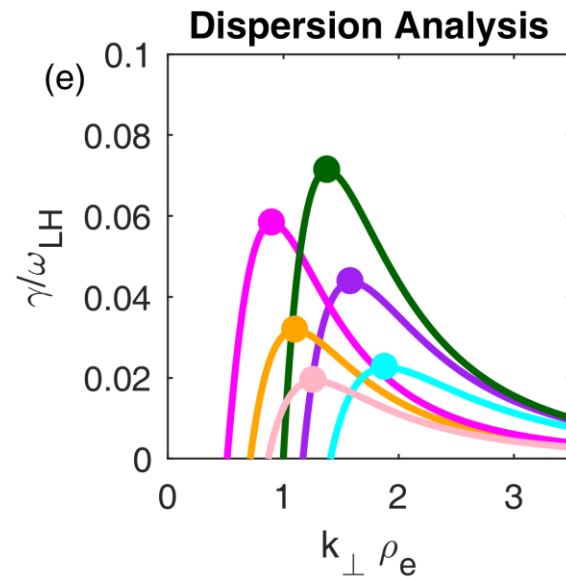
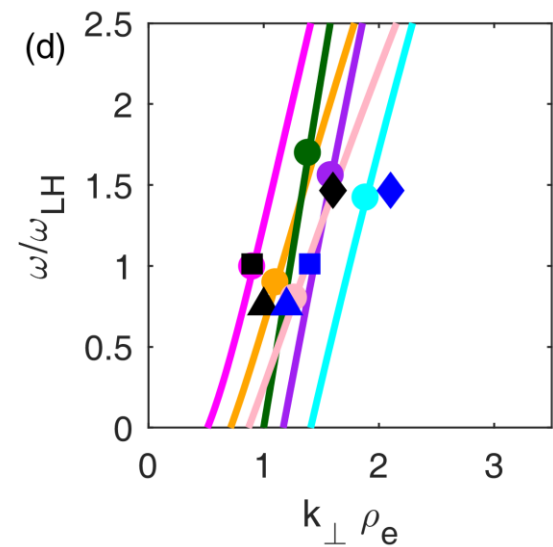
$$\mathbf{V}_{e\perp} \approx \mathbf{V}_E + \mathbf{V}_{de}$$

Davidson et al., 1977

	k (LMN)	V_{ph} (km/s)	Φ_B (V)	Φ_E (V)	C_Φ	k_⊥ρ_e	V_{ph} (km/s)	f (Hz)	k_⊥ρ_e	<D_⊥> (m ² s ⁻¹)
	Norgren et al. [2012]						Graham et al. [2019]			
MMS4	[0.62 -0.71 -0.33]	50(58)	5.7	5.8	0.83	1.0	43(51)	8.6(9.8)	1.2	~ 10 ⁵
MMS1	[0.72 -0.70 0.06]	65(67)	7.8	8.4	0.87	1.6	53(55)	20.4(20.9)	2.1	-2.6 × 10 ⁷
MMS3	[0.67 -0.72 -0.18]	94(87)	5.2	4.5	0.81	0.9	60(53)	17.3(15.7)	1.4	-1.8 × 10 ⁷

The numbers outside/inside the parentheses are estimated in the spacecraft/ion frame.

Instability analysis

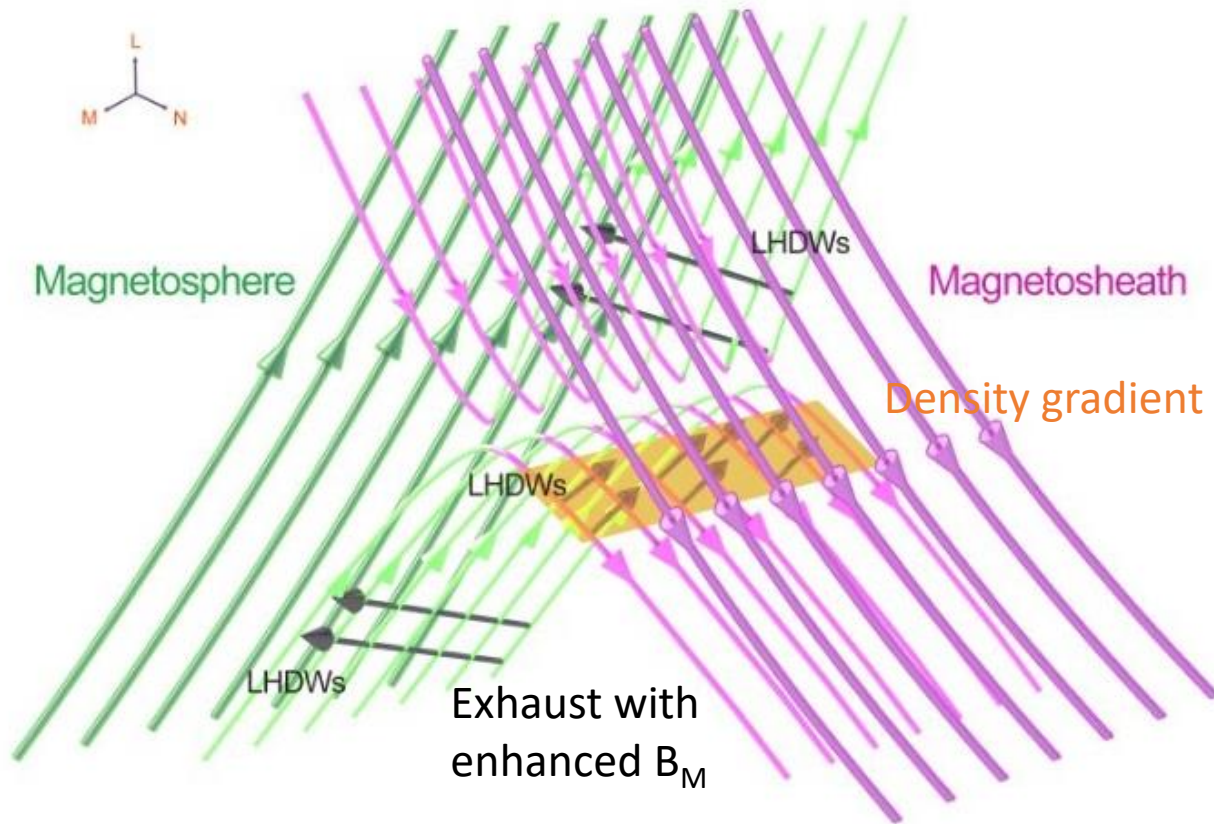


Black symbols: Norgren et al. [2012]

Blue ones: Graham et al. [2019]

	$\mathbf{k}$ (LMN)	V_{ph} (km/s)	Φ_B (V)	Φ_E (V)	C_Φ	$k_\perp \rho_e$	V_{ph} (km/s)	f (Hz)	$k_\perp \rho_e$	$\langle D_\perp \rangle$ ($\text{m}^2 \text{s}^{-1}$)
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MMS3	[0.67 -0.72 -0.18]	94(87)	5.2	4.5	0.81	0.9	60(53)	17.3(15.7)	1.4	-1.8×10^7

Summary



- LHDWs are observed at the magnetosheath separatrix in asymmetric guide-field reconnection;
- LHDWs are generated by lower hybrid drift instability;
- Properties of these LHDWs:
 - Active in limited regions;
 - Propagating toward the x-line, in the same direction the diamagnetic drift of the x-line;
 - $e\Phi/k_B T_e \sim 15\% - 25\% \rightarrow$ effectively transport particles from the sheath to the exhaust.

Thank you!

If you have any comments/suggestions, please contact
bbtang@spaceweather.ac.cn