

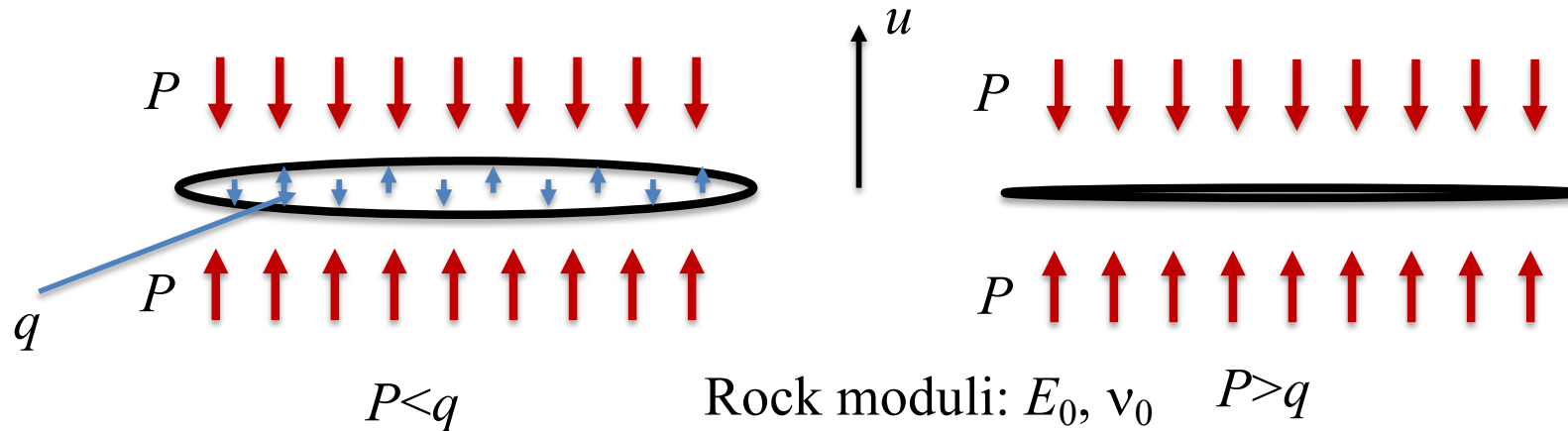
# Hydraulic fracture oscillations in response to strong impulse

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- Large hydraulic fracture with constrained opening as a bilinear oscillator
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- Spectrum of bilinear oscillator
- Excitation by a strong impulse and recovery of the stiffness ratio

# Large hydraulic fracture as a bilinear oscillator



Average displacement discontinuity,  $u$ , of a disc-like crack of radius  $R$

$$\bar{u} = \frac{16}{3\pi} \frac{1 - \nu_0^2}{E_0} R (q - P)$$

$$\bar{u} = -4\pi \frac{1 - \nu_0^2}{k_c} R^2 (q - P)$$

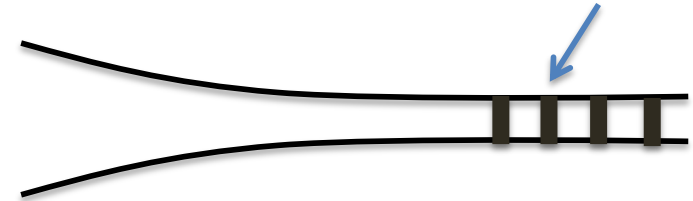
Stiffness of the contact

Stiffness ratio (compression to tension)

$$K = \frac{k_-}{k_+} = \frac{3}{4} \pi^2 \frac{E_0 R}{k_c} \rightarrow \infty \text{ as } R \rightarrow \infty \quad - \text{Strong bilinearity}$$

# Fractures with constraint opening

- Parts of unbroken rock connecting the opposite faces of the fracture
- Traditionally considered as obstacles for fluid flow
- Traditionally considered as a part of the process zone

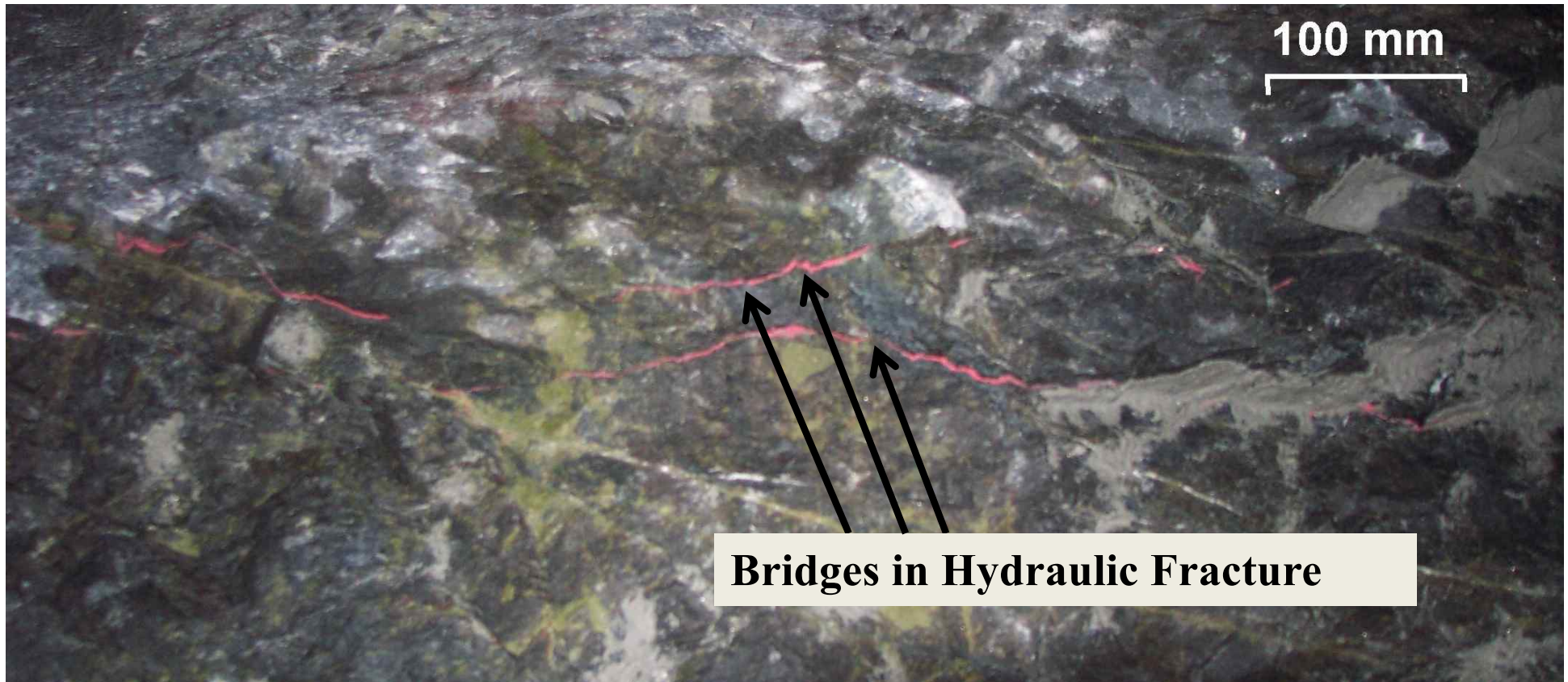


- What was overlooked before:

Bridges are all over the crack and can constrict the opening



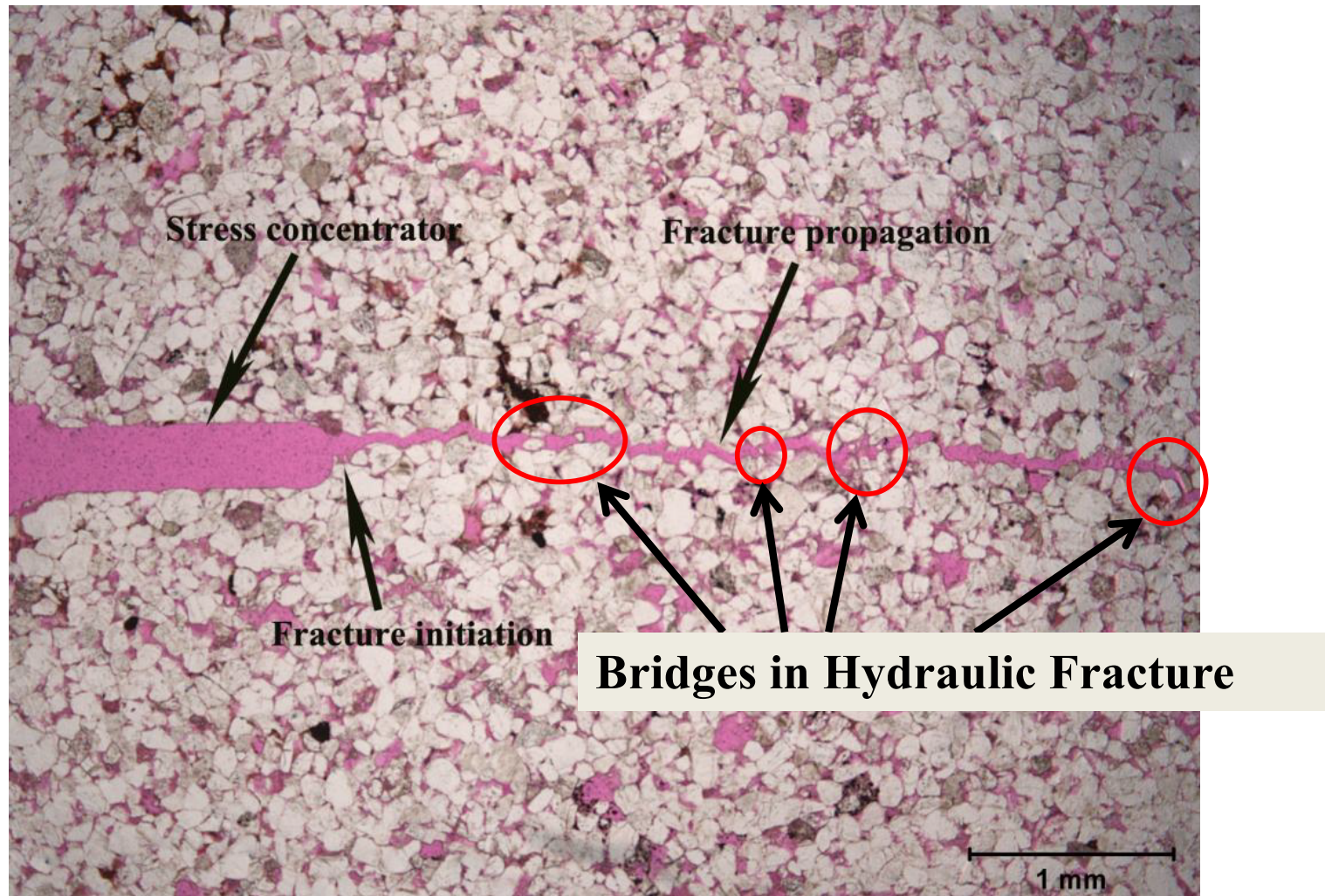
# Hydraulic fracture in mining



Northparkes mine, courtesy of Rob Jeffrey

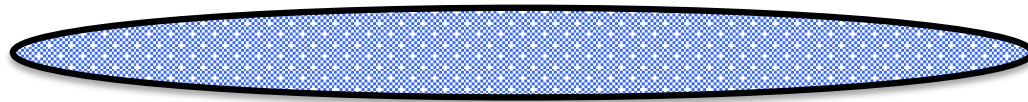


# Bridge Cracks



(Surdi 2010)

# Large hydraulic fracture with constrained opening (HFCO) as a bilinear oscillator



Conventional fracture



HF with constrained opening, effective when the crack radius,  $R$ , exceeds a constriction length,  $\lambda$ :

$$R \gg \lambda = E_0/k_b$$

**Fracture with constrained opening - bridges** distributed all over the fracture. Effective stiffness of bridges is low (He et al., 2020), hence  $k_b \ll k_c$ .

$$K = k_c/k_b \gg 1 \quad \text{- Strong bilinearity}$$

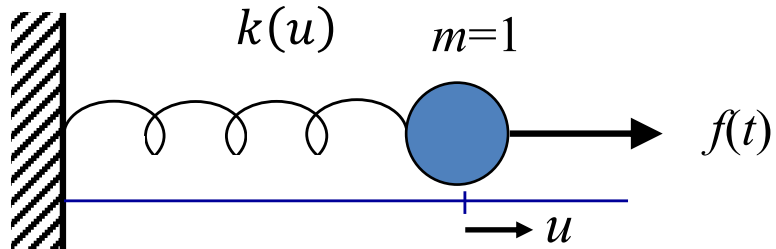
He, J., E. Pasternak and A.V. Dyskin, 2020. Bridges outside fracture process zone: Their existence and effect. *Engineering Fracture Mechanics*, 225, 106453

# Bilinear oscillator and multiple resonances

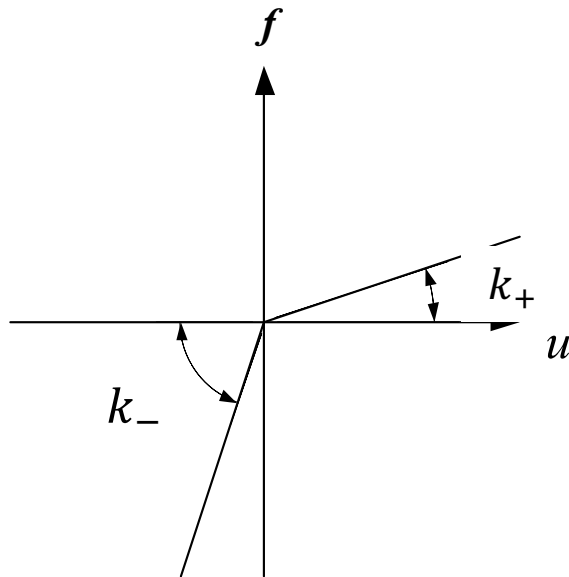
- Bilinear oscillator as a simplification of strong non-linearity
- Resonance frequency
- Multiple resonances
- Half resonance



# Bilinear oscillators



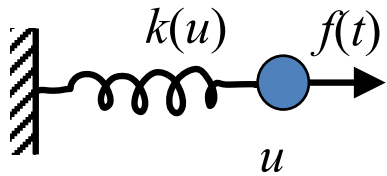
$$\ddot{u} + k(u)u = f(t)$$



$$k(u) = \begin{cases} \omega_+^2, & u > 0 \\ \omega_-^2, & u < 0 \end{cases}$$

Dyskin, A.V., E. Pasternak and E. Pelinovsky, 2012. Periodic motions and resonances of impact oscillators. *Journal of Sound and Vibration* 331(12) 2856-2873

# Bilinear oscillator

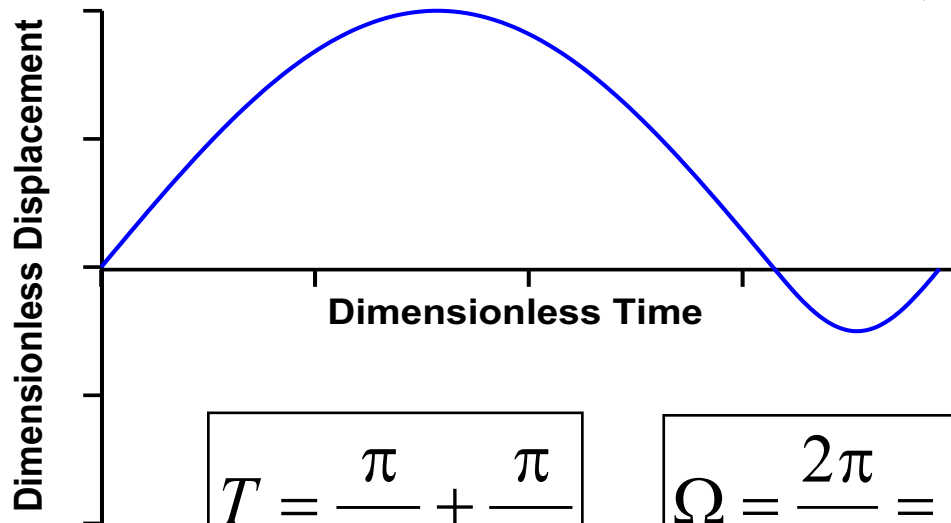


$$\ddot{u} + k(u)u = f(t)$$

$$k(u) = \begin{cases} \omega_+^2, & u > 0 \\ \omega_-^2, & u < 0 \end{cases}$$

Solution, one cycle

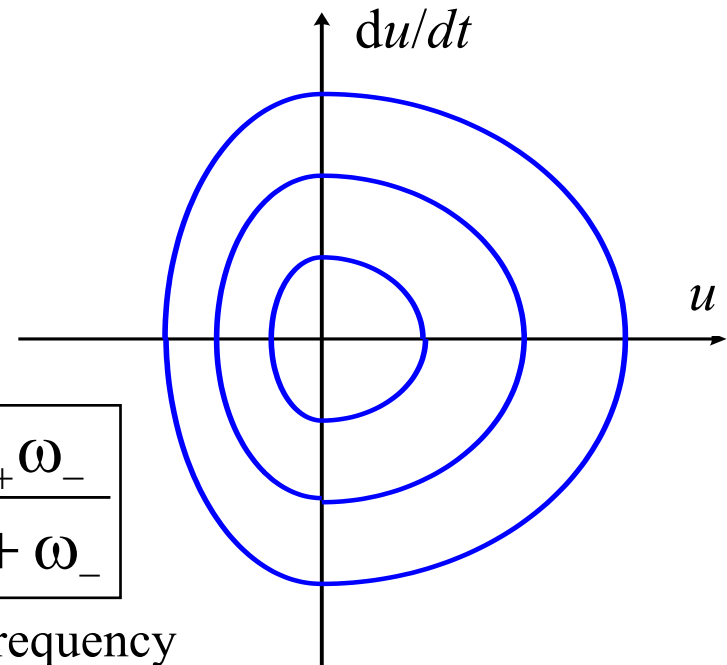
$$u(t) = A \begin{cases} \sin(\omega_+ t), & 0 < \omega_+ t < \pi, \\ \frac{\omega_+}{\omega_-} \sin[\omega_- t + \varphi], & \pi < \omega_- t + \varphi < 2\pi, \end{cases}$$



$$T = \frac{\pi}{\omega_+} + \frac{\pi}{\omega_-}$$

$$\Omega = \frac{2\pi}{T} = \frac{2\omega_+\omega_-}{\omega_+ + \omega_-}$$

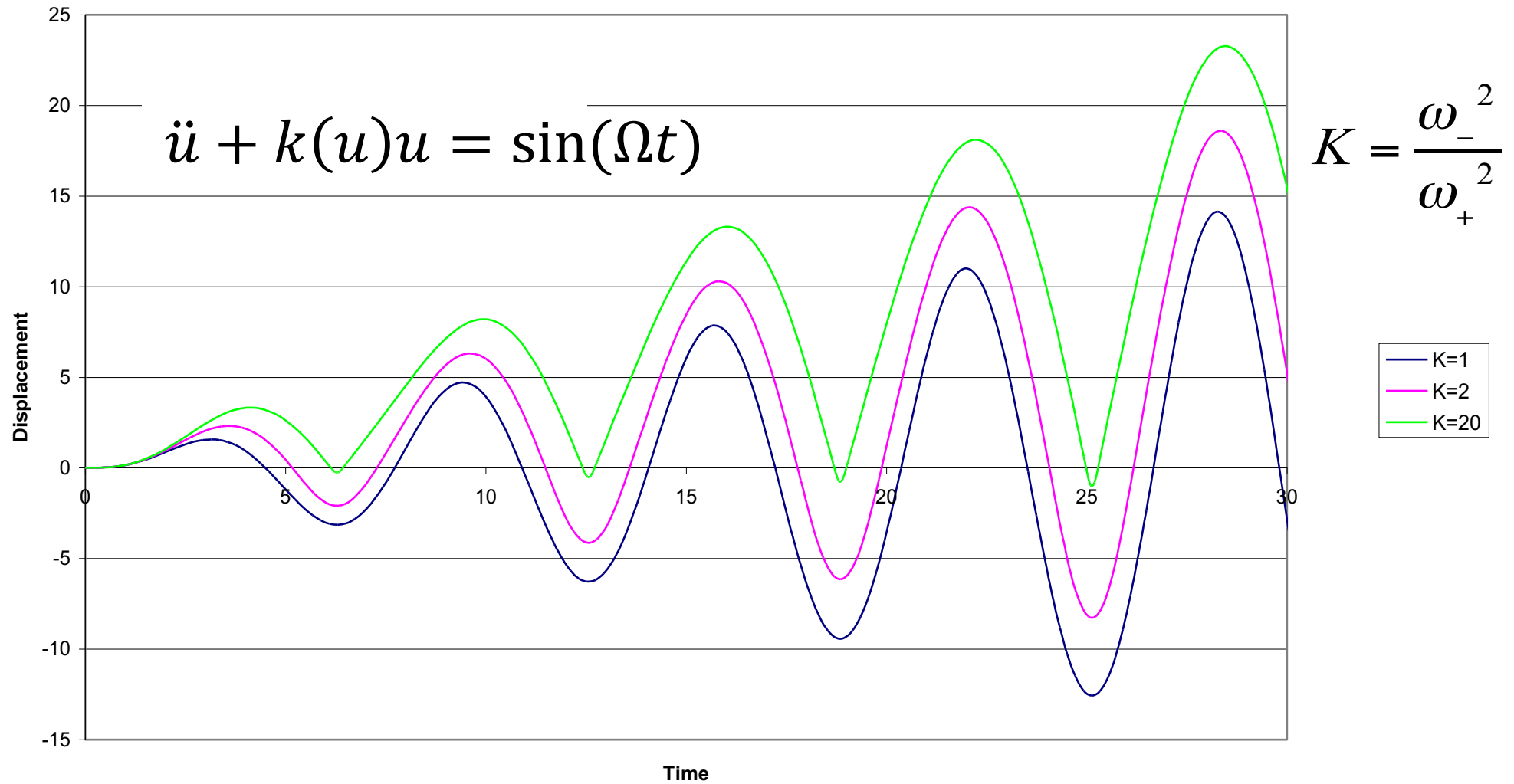
Resonance frequency



Thomson, J.M.T. (1983)

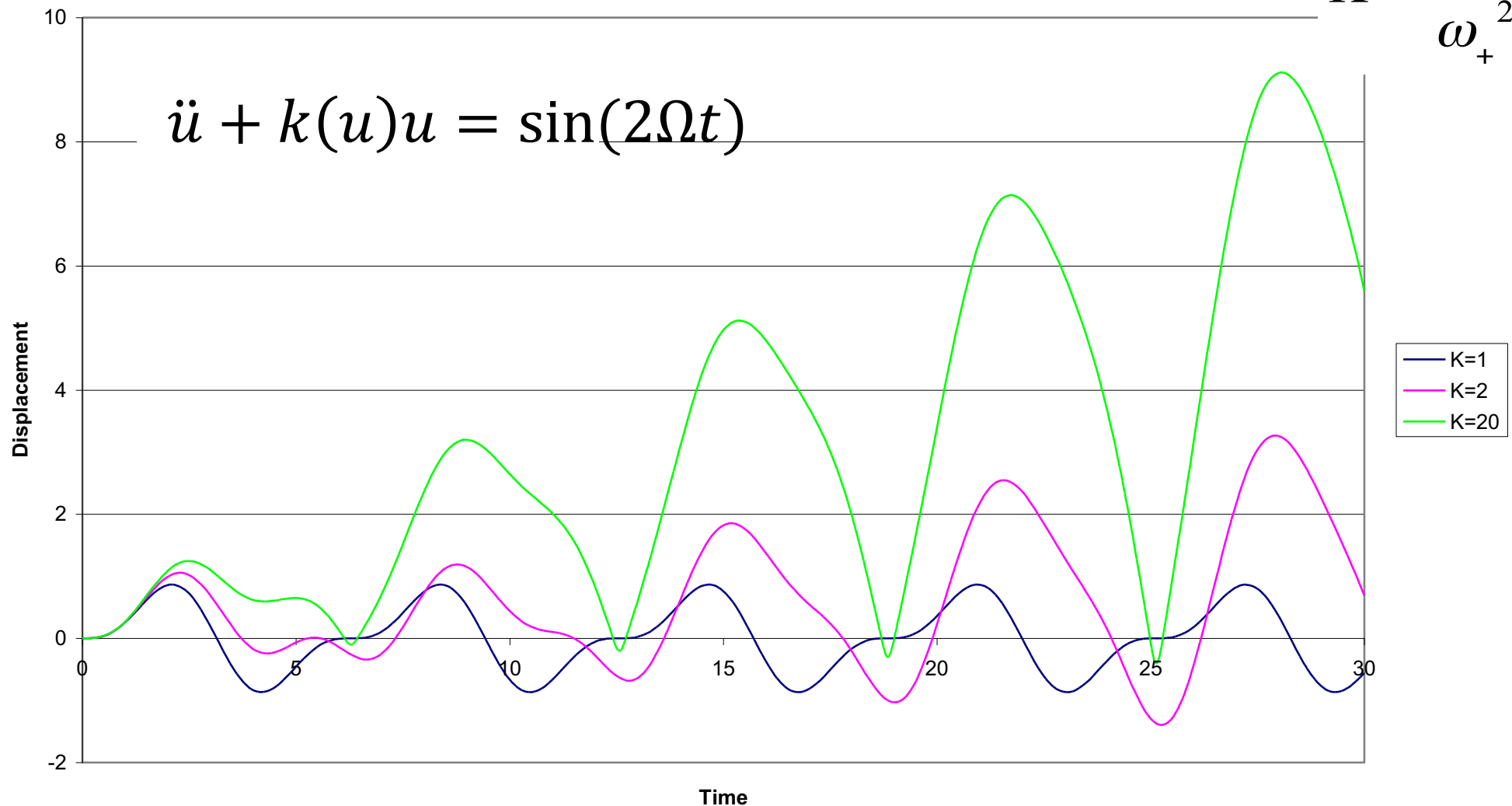
Dyskin, A.V., E. Pasternak and E. Pelinovsky, 2012. Periodic motions and resonances of impact oscillators. *Journal of Sound and Vibration* 331(12) 2856-2873

# Basic resonance



# Second resonance (excitation at double resonance frequency)

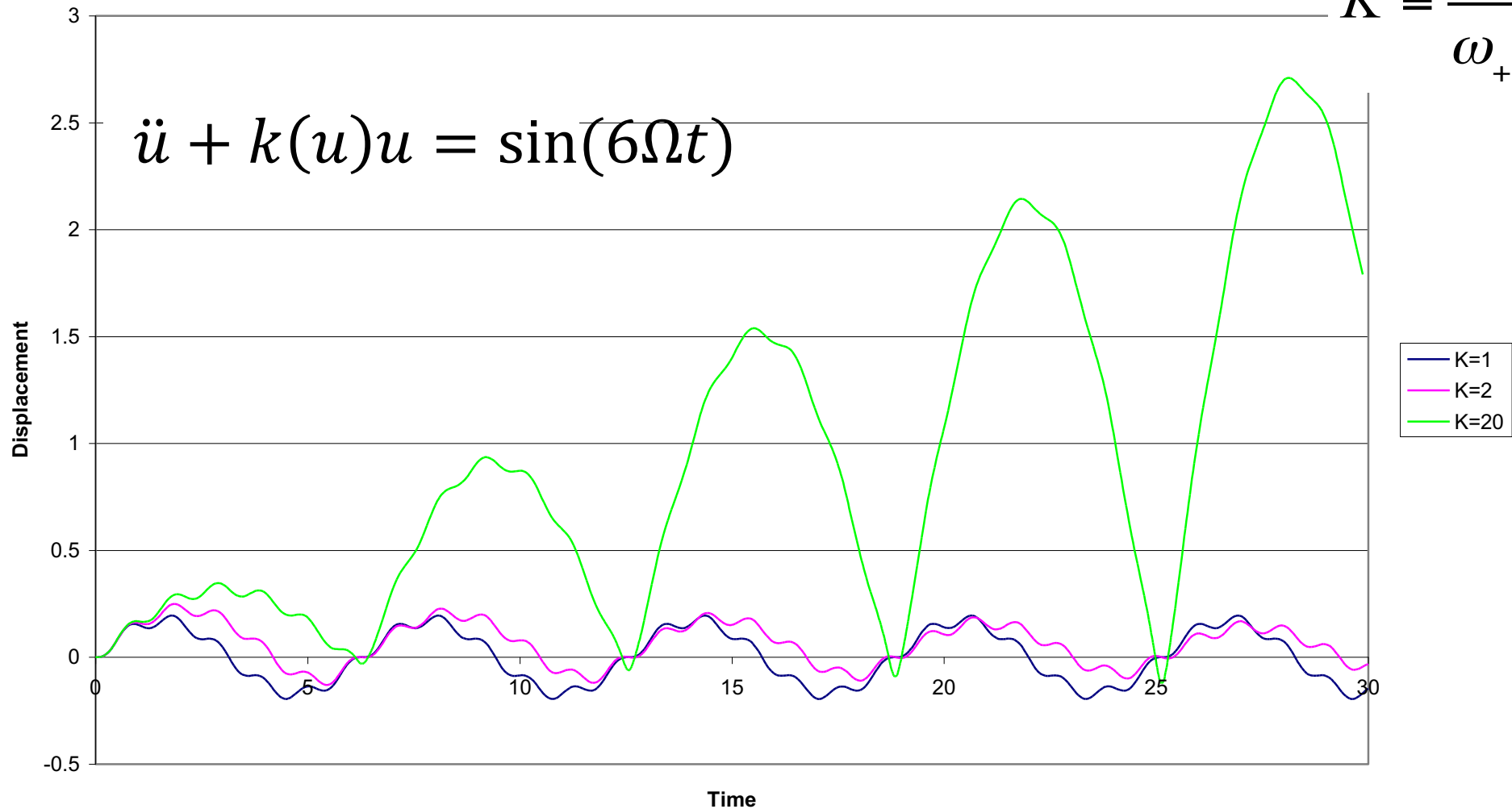
$$K = \frac{\omega_-^2}{\omega_+^2}$$



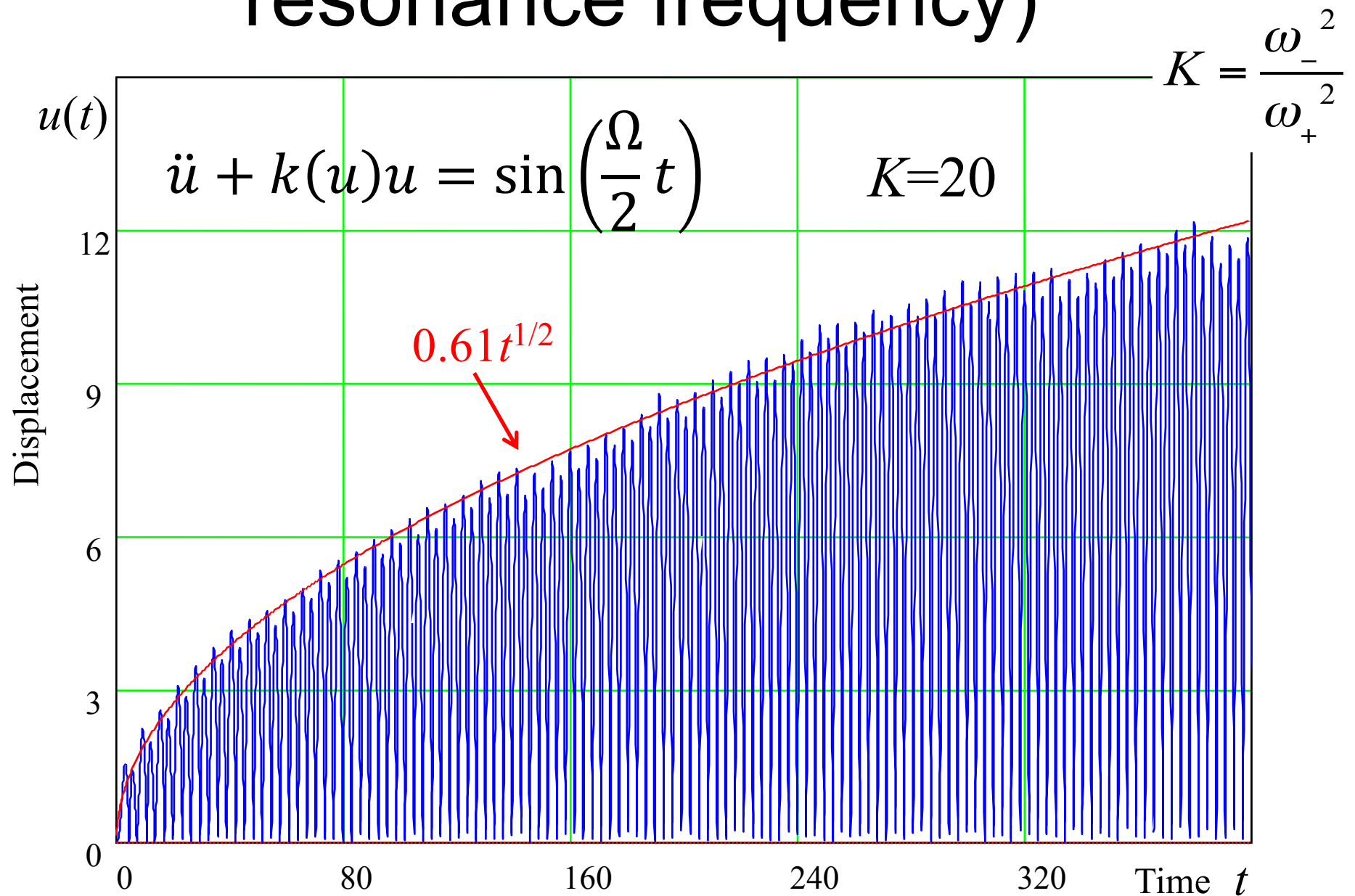
# Sixth resonance (excitation at 6 x resonance frequency)

$$K = \frac{\omega_-^2}{\omega_+^2}$$

$$\ddot{u} + k(u)u = \sin(6\Omega t)$$



# Half resonance (excitation at **half** resonance frequency)





# Spectrum of bilinear oscillator

- Period  $T = \frac{\pi}{\omega_+} + \frac{\pi}{\omega_-}$
- Fourier series for displacement discontinuity

$$u(t) = \frac{a_0}{2} + \sum_{k=1}^{\infty} \left[ a_k \cos\left(\frac{2\pi kt}{T}\right) + b_k \sin\left(\frac{2\pi kt}{T}\right) \right]$$

- Spectral amplitude

$$c_k = \sqrt{a_k^2 + b_k^2}$$

# First two harmonics

- Introduce  $\tau = \frac{\pi}{T\omega_+}$
- Ratio of the first two harmonics

$$\frac{c_1}{c_2} = \frac{\cos\left(\frac{\pi\tau}{2}\right)}{\cos(\pi\tau)} \frac{|4\tau(2-\tau) - 3||4\tau(2-\tau) + 5|}{|\tau(2-\tau) - 1||\tau(2-\tau) + 3|}, \quad \frac{1}{2} < \tau \leq 1$$

# Recovery of the stiffness ratio

- Find period  $T$
- Find ratio of two first harmonics  $\frac{c_1}{c_2}$
- Find  $\tau$
- Find stiffness ratio  $K = \frac{k_-}{k_+} = \frac{\omega_-^2}{\omega_+^2} = \frac{\tau^2}{(1 - \tau)^2}$

# Interpretation of the stiffness ratio

- Conventional HF

$$K = \frac{3}{4} \pi^2 \frac{E_0 R}{k_c}$$

- If rock modulus  $E_0$  and contact stiffness  $k_c$  are known, the fracture radius  $R$  is determinable

- HF with constrained opening (large crack radius),  $R \gg \lambda = \frac{E_0}{k_b}$

$$K = \frac{k_c}{k_b}$$

- Effective stiffness  $k_b$  of bridges is determinable

# Conclusions

- Large hydraulic fractures can behave as bilinear oscillators with stiffness in compression being considerably higher than stiffness in tension
- Bilinear oscillators exhibit multiple resonances
- Spectrum of bilinear oscillator is controlled by the period and stiffness ratio  $K = \frac{k_-}{k_+} = \frac{\omega_-^2}{\omega_+^2}$
- Finding spectrum of oscillations in response to external excitation allows the estimation of hydraulic fracture dimensions

# Acknowledgements

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# Publications

1. Dyskin A.V., E. Pasternak and E. Pelinovsky, 2007. Modelling resonances in topological interlocking structures. *ACAM2007*, Proc. 5th Australasian Congress on Applied Mechanics, 10-12 December 2007, Brisbane, Australia (F. Albermani, B. Daniel, J. Griffiths, D. Hargreaves, P. Meehan, A. Tan and M. Veidt, eds), 2, 408-413. Paper 1266.
2. Dyskin, A.V., Pasternak, E. and Pelinovsky, E. 2010. Coupled bilinear oscillators, their resonances and controlling parameters, *Proc. 6<sup>th</sup> Australasian Congress on Applied Mechanics*, ACAM 6, Kian Teh, Ian Davies and Ian Howard (eds.), 12-15 December 2010, Perth, Paper 1170
3. Dyskin, A.V., E. Pasternak and E. Pelinovsky, 2012. Periodic motions and resonances of impact oscillators. *Journal of Sound and Vibration* 331(12) 2856-2873. ISBN/ISSN 0022-460X, 04/06/2012.
4. Dyskin, A.V., E. Pasternak and I. Shufrin, 2014. Structure of resonances and formation of stationary points in symmetrical chains of bilinear oscillators. *Journal of Sound and Vibration* 333, 6590–6606.
5. He, J., A.V. Dyskin, E. Pasternak, M. Lebedev and B. Gurevich, 2017. The effect of constriction in hydraulic fracturing. *Bifurcation and Degradation of Geomaterials with Engineering Applications - Proceedings of the 11th International Workshop on Bifurcation and Degradation*. Limassol, Cyprus. Papamichos, E., P. Papanastasiou, E. Pasternak, A.V. Dyskin (eds.), Springer, Cham, Switzerland, 613-619.
6. He, J., E. Pasternak, A.V. Dyskin, M. Lebedev and B. Gurevich, 2017. The constricting effect of bridges in hydraulic fracturing [online]. In: 9th Australasian Congress on Applied Mechanics (ACAM9). Sydney: Engineers Australia, 318-323.
7. He, J., Pasternak, E. and A.V. Dyskin, 2020. Bridges outside fracture process zone: Their existence and effect. *Engineering Fracture Mechanics*, 225, 106453.
8. Pasternak, E., A. Dyskin and Ch. Qi, 2020. Impact oscillator with non-zero bouncing point. *International Journal of Engineering Science*, 103203.