

## Splitting in the static $k_2$ of the Neptune-Triton system

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### Abstract

We present fluid Love number  $k_{nm}$  values for the Neptune-Triton system as a function of the uncertainty in Neptune's  $J_4$  value and rotation rate. We discuss the unexpected found splitting in  $k_2$  for assumed slow and fast rotation, respectively, in regard to differences in the contribution functions of  $J_2$  and  $k_2$ .

### 1. Dissimilar Ice Giants?

The Ice Giants Uranus and Neptune have similar mass, radius, gravity field ( $J_2, J_4$ ), and similar magnetic fields. On the other hand, they differ in luminosity and rotation period  $P_{\text{rot}}$ . This poses the question of perhaps dissimilar interiors of the ice giants. Here we investigate if an observed fluid number  $k_2$  value for the Neptune-Triton system could help to further constrain the interior structure of Neptune.

### 2. Splitting in $I$ and $k_2$

Table 1: Some observational constraints.

Parameter	Value	Ref.
Neptune		
$P_{\text{rot}, 1}$ (fast)	16h 6m 40s	[5]
$P_{\text{rot}, 2}$ (slow)	17h 27m 29s	[2]
$R_{\text{eq}, 1}$ at 1 bar	24,766 km	[5]
$R_{\text{eq}, 2}$ at 1 bar	24,787 km	[2]
$q_{\text{rot}}$	0.02608	
Triton		
$M_{\text{Tri}}/M_{\text{Nep}}$	0.0048	
$q_{\text{tid}}$	$-2.13 \cdot 10^{-7}$	

The rotation rates of Uranus and Neptune are not even well known, with estimates differing by 40 min for Uranus and 1 h 21 min for Neptune, see Table 1. For Neptune, the faster rotation rate known as the *Voyager* rotation rate was inferred from radio and magnetic field data while the slower value is a prediction

from interior models that fit the shape data and minimize the observed wind velocity and dynamical heights [2].

In Ref. [6] it was found that interior models for the different rotation rates yield different normalized moment of inertia values  $I_{\text{slow}} = 0.2410(8)$  and  $I_{\text{fast}} = 0.2555(2)$  although the models would fit the same  $J_2$  value. In Ref. [7] a similar splitting was obtained for  $k_2$ . This is surprising since  $k_2$  is extraordinarily insensitive to different internal density distributions once  $J_2$  and  $J_4$  are fit, as has been shown for Jupiter [8] and Saturn [9].

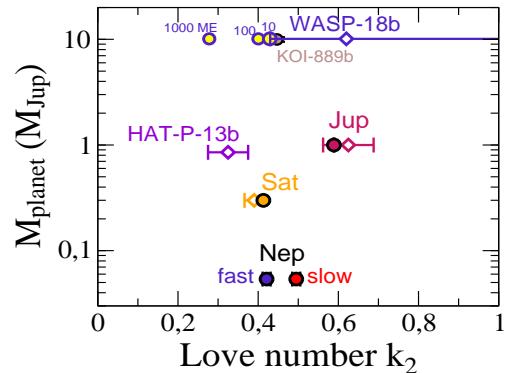


Figure 1: Observed (diamonds) and calculated  $k_2$  values of individual planets and brown dwarfs, plotted over their mass.

In Figure 1 we plot observed  $k_2$  values for Jupiter [3], Saturn [4], the hot Jupiter HAT-P-13b, and the massive hot Jupiter WASP-18b [1] in comparison to model predictions. For Neptune, the splitting in  $k_2$  between slow and fast rotating models is large; it is of same size as for changing the core mass of a  $10 M_{\text{Jup}}$  planet from 0 to  $\sim 300 M_{\text{E}}$ .

In contrast, models for Jupiter and Saturn that are constrained by  $J_2$  and  $J_4$  yield tight (0.02%)  $k_2$  ranges [8, 9], currently in agreement with the Juno and Cassini

data based measurements.

### 3. Outlook

In order to understand the behavior of  $k_2$  as a function of rotation rate we calculate the contribution functions of  $k_2$  and  $J_2$  in this work. We also explore whether different scalings could be responsible for the splitting. Gaining a better understanding of the behavior of  $k_2$  and the moment of inertia is important because at present, our results suggest that both parameters could be useful for inferring the solid rotation rate independently on radio or magnetic field data.

However, further effects may influence the planetary tidal response, such as internal oscillations or atmospheric dynamics. The tiny overlap in  $k_2$  between interior model predictions and the observed value for WASP-18b may indicate that our static approach is too simplistic.

Observing  $k_2$  at Neptune might require a polar orbiter that covers different longitudes.

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