

## N<sub>2</sub> states population and airglow in Titan's atmosphere

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### Abstract

Molecular nitrogen in Titan's atmosphere is excited to different vibrational and electronic states by direct photon absorption and electron impact. Here we present detailed calculations for the vibrational population of different electronic states and the emission rates from the de-excitation of these states.

### 1. Introduction

The interaction of N<sub>2</sub> with high energy photons and electrons excites the molecule to different vibrational and electronic levels (N<sub>2</sub><sup>\*</sup>). De-excitation from these states can result by spontaneous emission and by collisions of the N<sub>2</sub><sup>\*</sup> with other molecules. Competition among these processes will define the thermal structure of the atmosphere. Therefore a detailed calculation for the population of the different excited states is necessary in order to calculate an accurate temperature profile, as well as emission rates that can be compared with spacecraft observations such as those by the Cassini/UVIS instrument for Titan's case.

### 2. Model Description

The starting point of our calculations are the excitation and ionization rates based on high resolution cross sections of N<sub>2</sub> that we have performed in the past [1]. We expand on these calculations by calculating the population for the first 21 vibrational levels of the major singlet and triplet electronic states of N<sub>2</sub> and doublets of N<sub>2</sub><sup>+</sup>. These include the states: X <sup>1</sup>Σ<sub>g</sub><sup>+</sup>, A <sup>3</sup>Σ<sub>u</sub><sup>+</sup>, B <sup>3</sup>Π<sub>g</sub>, W <sup>3</sup>Δ<sub>u</sub>, B' <sup>3</sup>Σ<sub>u</sub><sup>-</sup>, α' <sup>1</sup>Σ<sub>u</sub><sup>-</sup>, α <sup>1</sup>Π<sub>g</sub>, w <sup>1</sup>Δ<sub>u</sub>, b <sup>1</sup>Π<sub>u</sub>, c' <sup>1</sup>Σ<sub>u</sub><sup>+</sup>, C <sup>3</sup>Π<sub>u</sub>, E <sup>3</sup>Σ<sub>g</sub><sup>+</sup> for N<sub>2</sub>, and X <sup>2</sup>Σ<sub>g</sub><sup>+</sup>, A <sup>2</sup>Π<sub>u</sub>, B <sup>2</sup>Σ<sub>u</sub><sup>+</sup> and C <sup>2</sup>Σ<sub>u</sub><sup>+</sup> for N<sub>2</sub><sup>+</sup>.

The population of each vibrational/electronic level depends on the rate this level is populated by direct excitation by photons/electrons and by cascade from

higher lying states (see e.g. Fig. 1), and the rate at which it is depopulated by radiative de-excitation and de-excitation by collisions with other atmospheric species. Therefore, in order to calculate the population of the different levels we need a large number of parameters that define the transition rates among different states. These include electron impact cross sections and Frank-Condon factors for the calculation of the production rates of different vibrational levels within each electronic state, transition rates among vibrational levels within each electronic state, transition rates among vibrational levels of different electronic states, as well as de-excitation rates of different states by molecular collisions and losses in state population by pre-dissociation. We obtain these parameters from the latest compilations and studies available [2, 3, 4, 5].

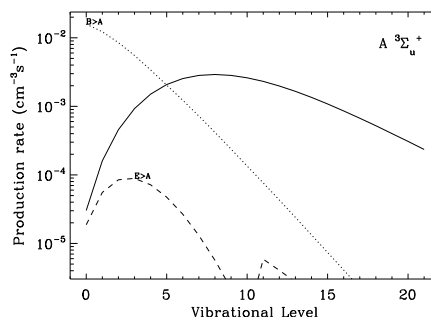


Figure 1: Production rates for the A <sup>3</sup>Σ<sub>u</sub><sup>+</sup> state at different vibrational levels. Contributions from direct electron impact excitation (solid line) and cascade from higher electronic states (broken lines) are shown.

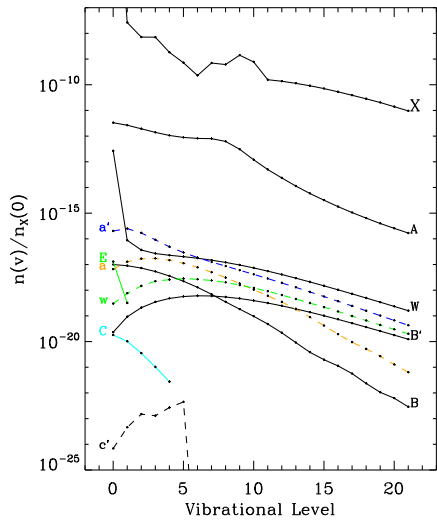


Figure 2: Vibrational populations for different  $N_2$  states scaled to the density of the  $N_2 X^1\Sigma_g^+(\nu = 0)$  level.

### 3. Results

Our calculations for the vibrational population for different electronic states close to 1100 km in Titan's atmosphere are shown in Fig. 2. The longest lived metastable state is the  $N_2 A^3\Sigma_u^+$  state followed by the ground vibrational level of the  $W^3\Delta_u$  state. All other states/levels have small populations due to their rapid de-excitation through radiative emission and/or collisions, the relative role of each of these processes depending on the local atmospheric density. Close to 1100 km radiative transitions dominate providing an emission spectrum that covers a wide range of wavelengths (Fig. 3). We will further discuss what constraints on the atmospheric properties we can derive by comparison with the available observations.

### References

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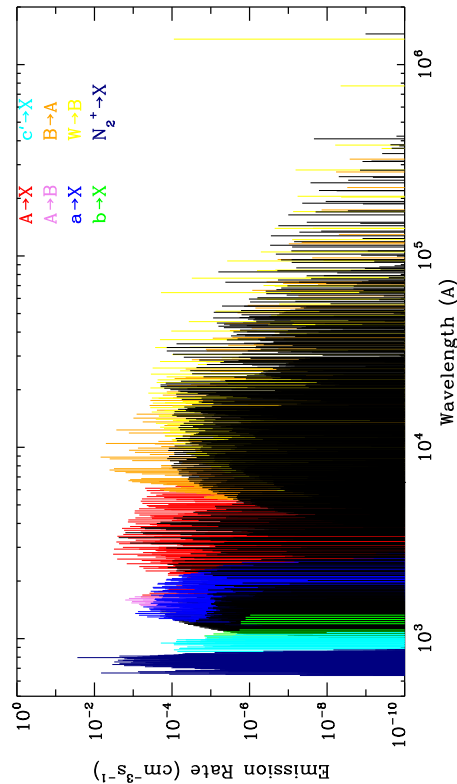


Figure 3: Emission rates for all transitions included in the model. Strongest emissions are color coded.